

Light Hadron Spectroscopy in Two-flavor Lattice QCD with Small Sea Quark Masses

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1. Introduction
2. Simulations
3. Results(Chiral extrapolations)
4. Conclusion

Introduction

purposes for hadron spectrum calculations

- Verification of QCD at low energies
- Basis of the other lattice calculations
(matrix elements, finite temperature/density)
- Quark masses

Features of lattice QCD

Advantage]

- Non-perturbative calculations

Disadvantage]

- Statistical and systematic errors
 - Continuum extrapolation($a \rightarrow 0$)
→ improvement of actions and operators
 - Renormalization factor
→ non-perturbative renormalization method
 - Finite size effect(studied well)
 - **Chiral extrapolation**

Chiral extrapolation]

in lattice QCD, sea quark masses

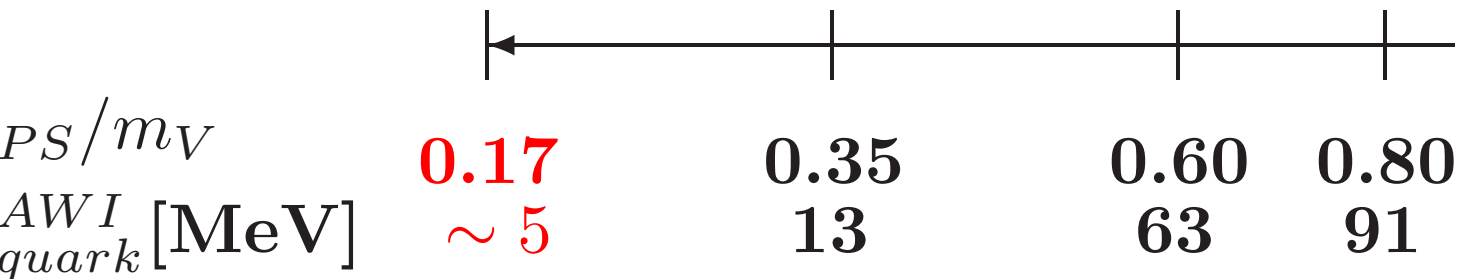
in two-flavor cases, $m_{ud} \equiv (m_u + m_d)/2$

are heavier than those in our real world

because of computational costs

lattice data must be extrapolated
to the physical point.

lattice $m_{PS}/m_V \rightarrow$ real $m_\pi/m_\rho = 0.17$



- Long extrapolations involve sizable systematic errors(as we see later)
→ more realistic simulations are desirable

3.1 [Computational cost]

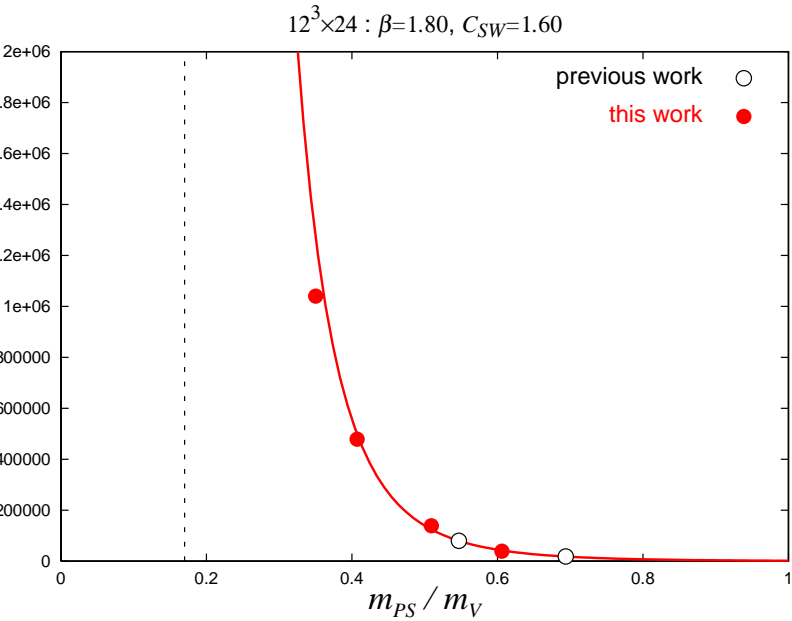
cost increases extremely

as sea quark masses become small

- CPU time ($\sim N_{inv}/dt$) $\propto (m_{PS}/m_V)^{-6}$

N_{inv} : # of inversions for the quark matrix

dt : step size in molecular dynamics



Simulations

$f = 2$ full QCD simulations

Action and run parameters]

- Small sea quark masses
($m_{PS}/m_V = 0.60 - 0.35$, $m_{quark} = 63 - 13$ MeV)
- Coarse lattices ($a = 0.2[\text{fm}]$, $L = 2.4 - 3.2[\text{fm}]$)
- RG gauge + tadpole-improved Clover quark
- HMC with the simple leapfrog scheme
- Measurements : every 5 trajectory (1 conf.)

Statistics]

combined with our previous work, **CP-PACS,2002**

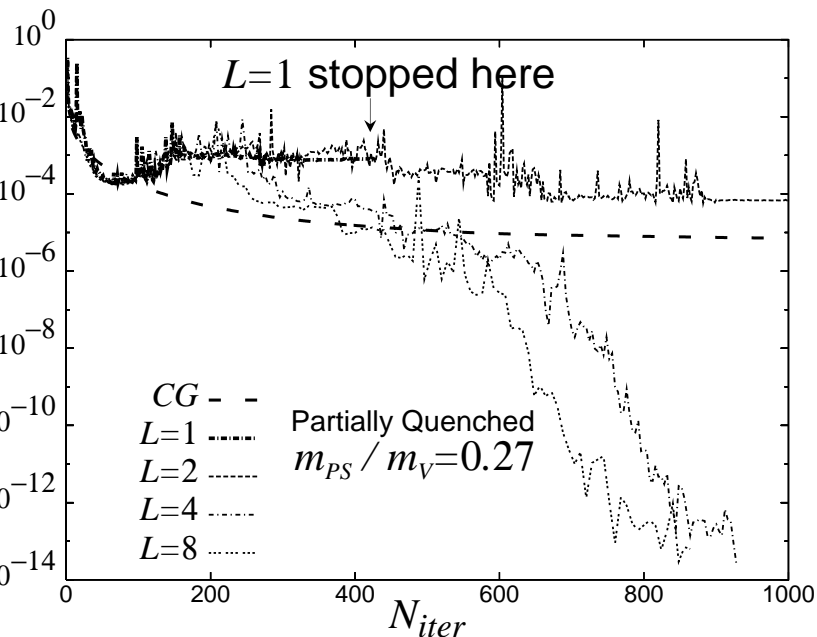
[$12^3 \times 24$]

κ_{sea}	m_{PS}/m_V	# traj.
0.14090	0.80	6250
0.14300	0.75	5000
0.14450	0.70	7000
0.14585	0.60	4000
0.14640	0.55	5250
0.14660	0.50	4000
0.14705	0.40	4000
0.14720	0.35	1400

3.2 [Issues on algorithms]

Two problems in small sea quark mass regions

- BiCGStab for quark inversions often does not work : $Ax = b \rightarrow x = A^{-1}b$
 \rightarrow Solved by BiCGStab(L)
 L : the order of the minimal residual polynomial
S.Itoh and YN,2003



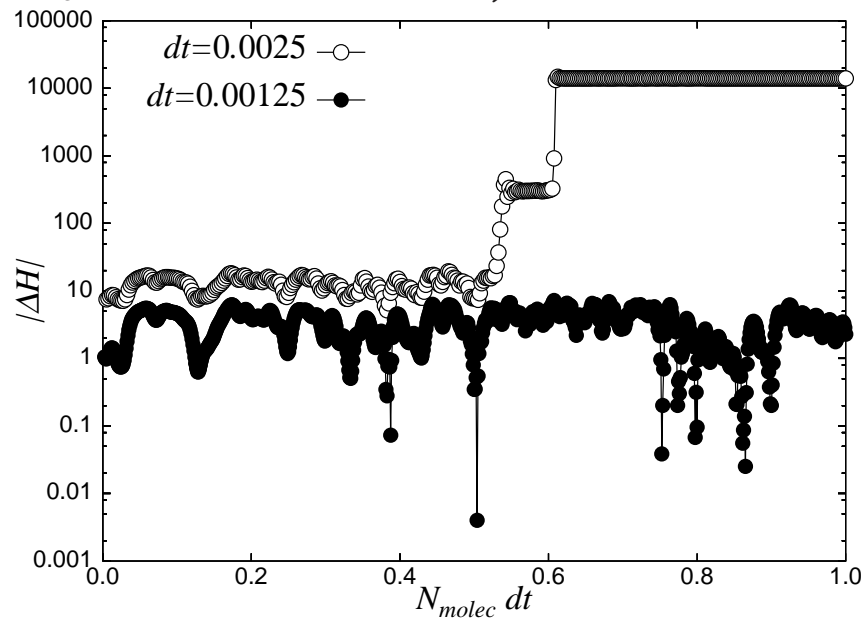
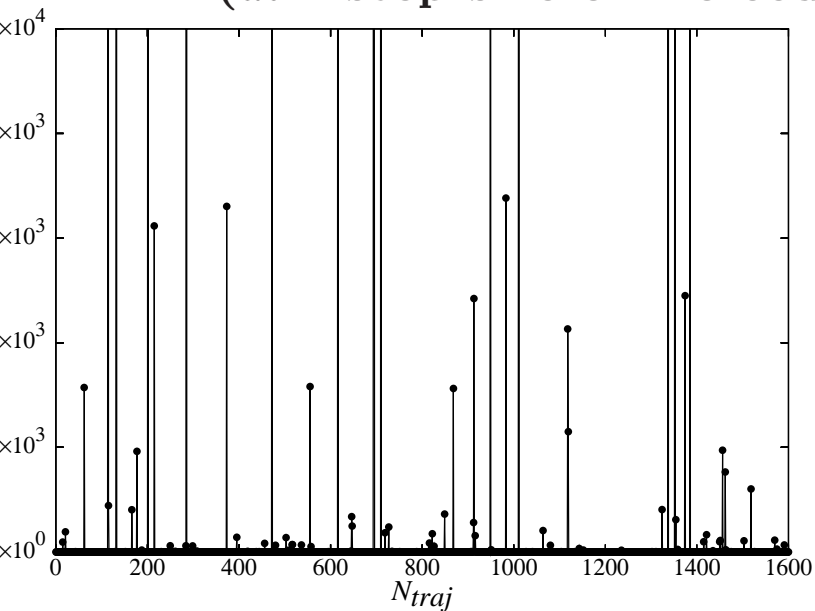
- **Efficiency decrease by huge $\Delta H \equiv H_{trial} - H_{start}$**
Jansen and Sommer, 1998

- As sea quark masses become smaller, more frequently huge ΔH appears

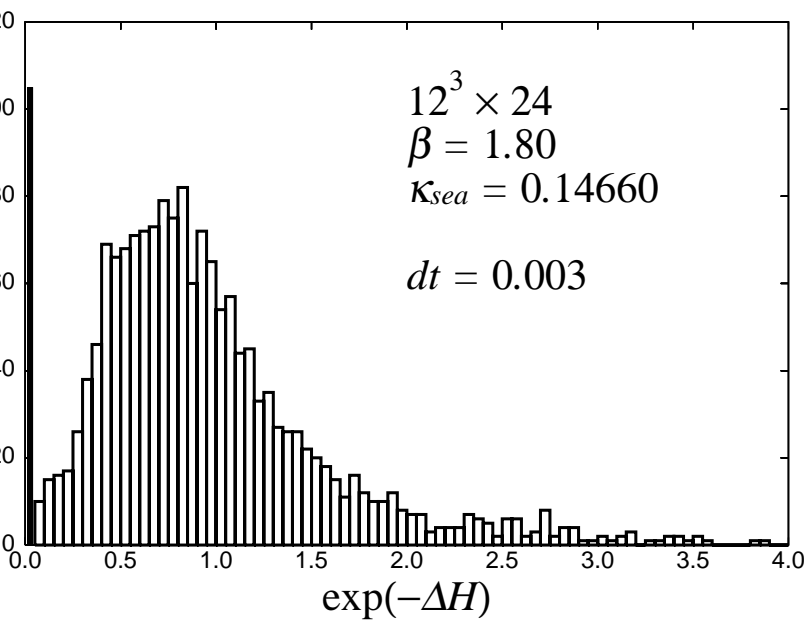
$$P_{acc} = \text{erfc} \left(\frac{1}{2} \sqrt{\langle \Delta H \rangle} \right) \sim \exp \left(-\sqrt{\frac{\langle \Delta H \rangle}{\pi}} - \frac{\langle \Delta H \rangle}{2\pi} \right)$$

$$\langle \Delta H \rangle \sim O((dt)^4 V) \sim O(1)$$

→ The analysis has not been finished yet.
 The prescription is employing small dt .
 (dt : step size of molecular dynamics in HMC)



the histogram of spikes



Results

[Chiral extrapolation]

parameterize lattice data O with m_{quark}
or chiral extrapolations

$$O = f(m_{quark})$$

choice of the extrapolation function f

- polynomial (conventional method)
 - Chiral perturbation (ChPT) formulae
 - Wilson ChPT (WChPT) formulae
- The physical quark mass $m_{quark}^{physical}$ is obtained by

$$m_{quark}^{physical} = f^{-1}(O^{physical})$$

Polynomial (conventional) extrapolation

- Check the validity of the previous quadratic fit from $m_{PS}/m_V = 0.80 - 0.55$ by comparing with new small sea quark mass data $m_{PS}/m_V = 0.50 - 0.35$

[Vector meson mass]

We parameterize vector meson masses m_V as a function of pseudoscalar meson masses m_{PS}

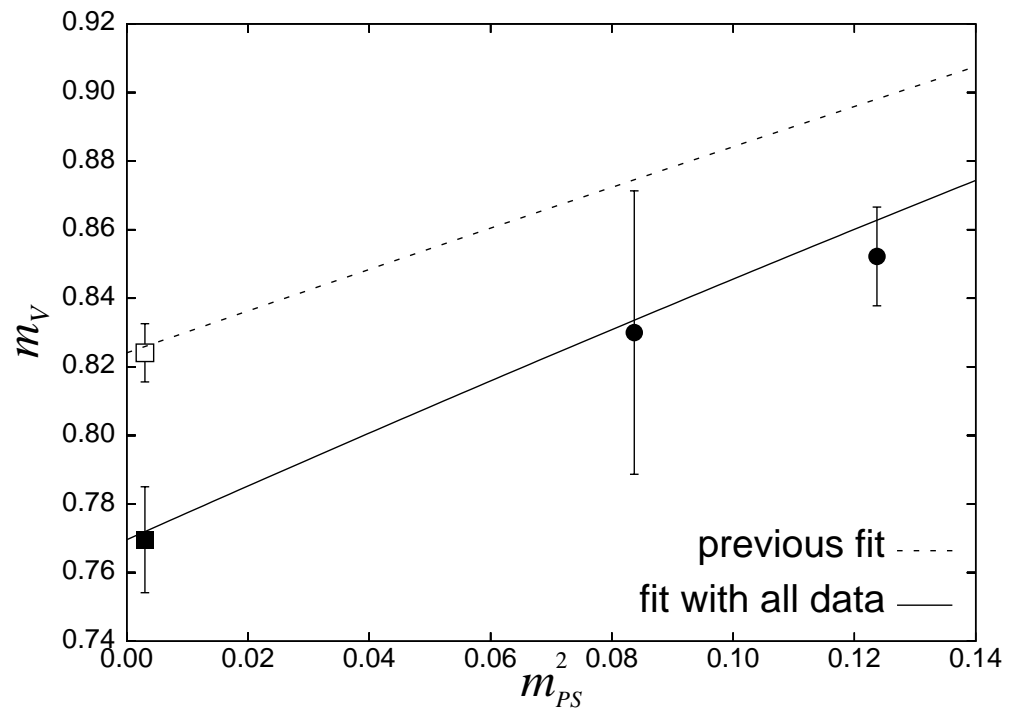
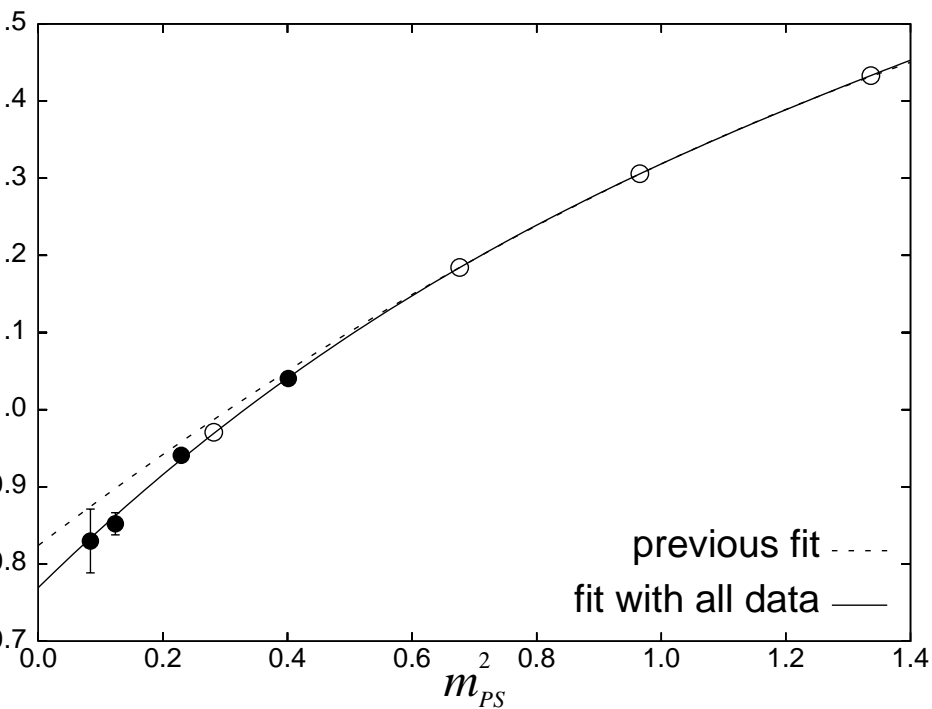
$$m_V = A + Bm_{PS}^2 + Dm_{PS}^4 + Fm_{PS}^6 \quad \text{for } m_{PS}/m_V = 0.80 - 0.35$$

$$m_V = A + Bm_{PS}^2 + Dm_{PS}^4 \quad \text{for } m_{PS}/m_V = 0.80 - 0.55$$

(previous fit) **CP-PACS,2002**

- We use m_{PS}^2 instead of m_{quark} for comparison with the previous work (We parameterize m_{PS}^2 as a function of m_{quark} , as explained later)

- Systematic deviations from the previous fit are observed in small sea quark mass region
 → Sea quark mass dependence was underestimated by the previous quadratic extrapolation from $m_{PS}/m_V = 0.80 - 0.55$
 The deviation in the chiral limit is $7\%(3.5\sigma)$



Pseudoscalar meson mass]

We parameterize pseudoscalar meson masses m_{PS}
as a function of VWI quark masses m_{quark}^{VWI}

$$m_{PS}^2 = Bm_{quark}^{VWI} + C(m_{quark}^{VWI})^2 + D(m_{quark}^{VWI})^3 + E(m_{quark}^{VWI})^4$$

for $m_{PS}/m_V = 0.80 - 0.35$

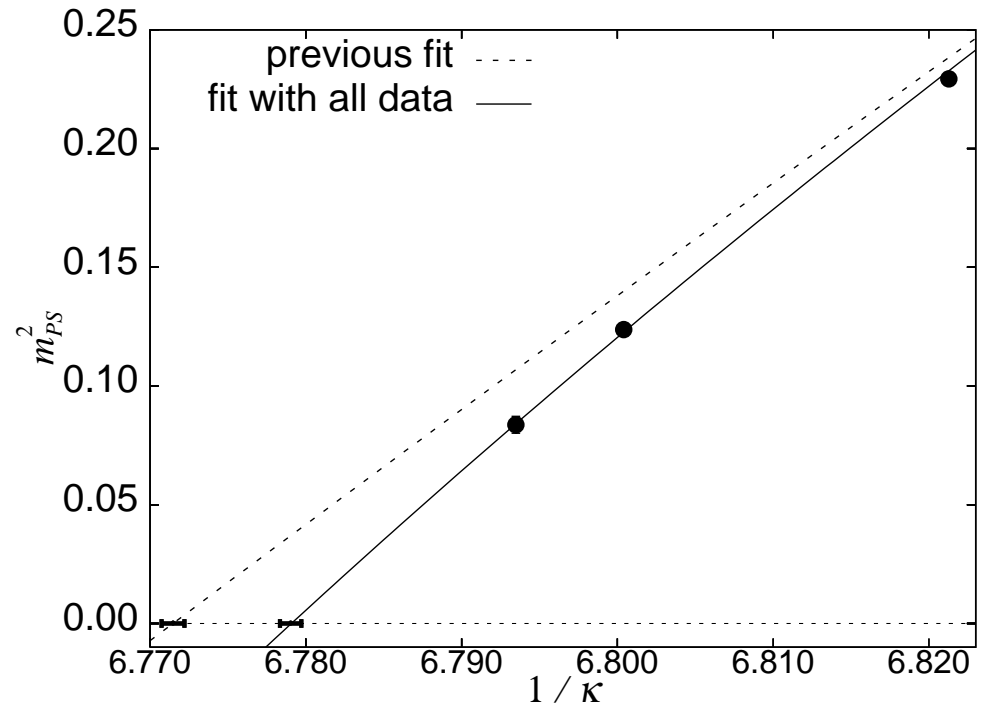
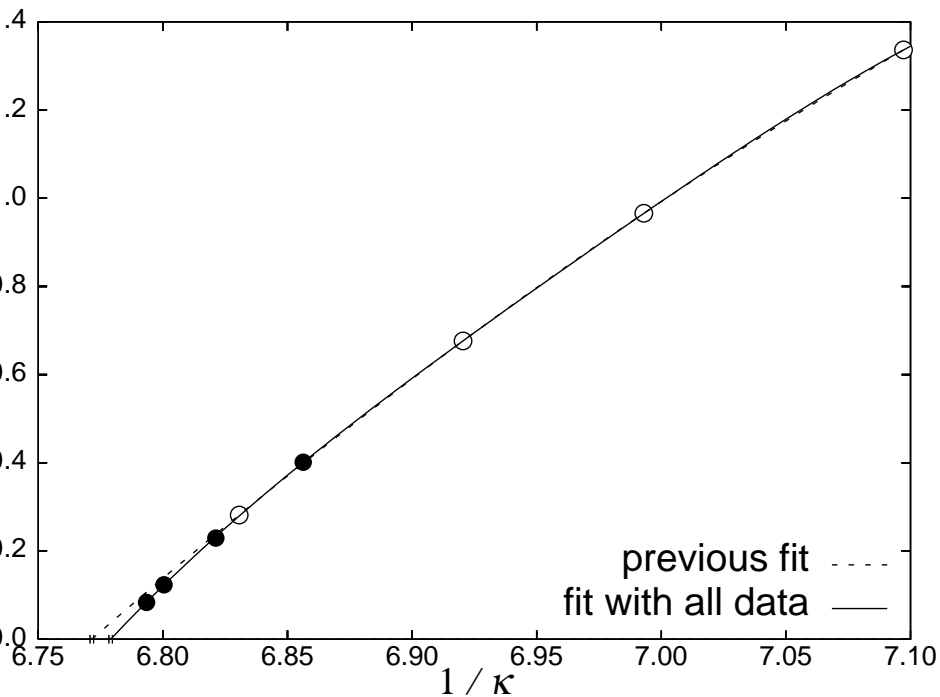
$$m_{PS}^2 = Bm_{quark}^{VWI} + C(m_{quark}^{VWI})^2$$

for $m_{PS}/m_V = 0.80 - 0.55$ (previous fit) **CP-PACS,2002**

where

$$m_{quark}^{VWI} = \frac{1}{2} \left(\frac{1}{\kappa} - \frac{1}{\kappa_c} \right)$$

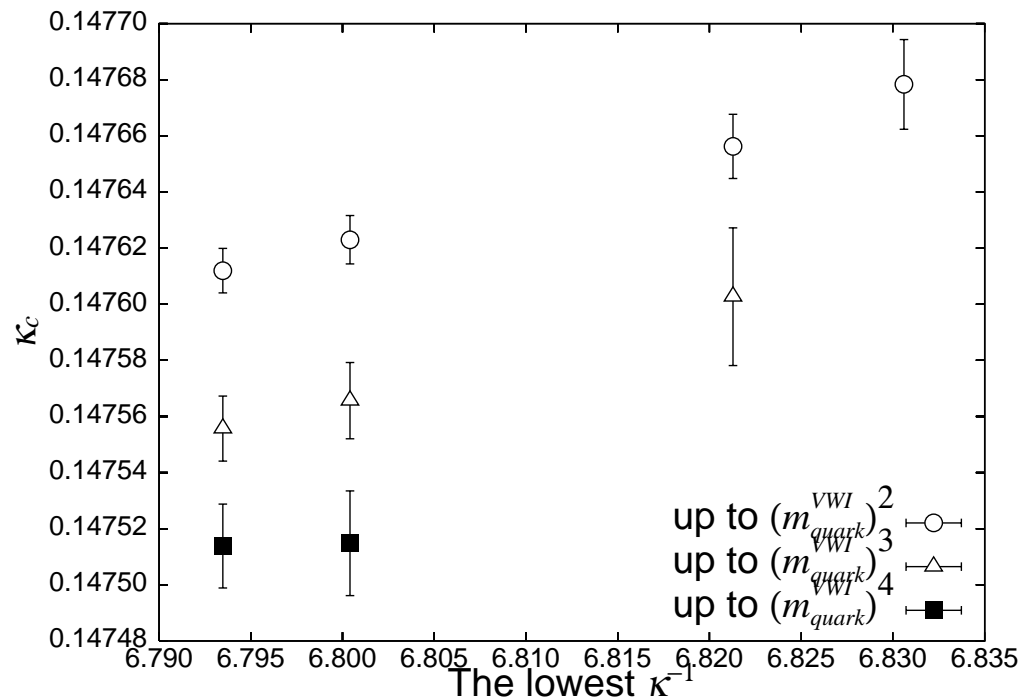
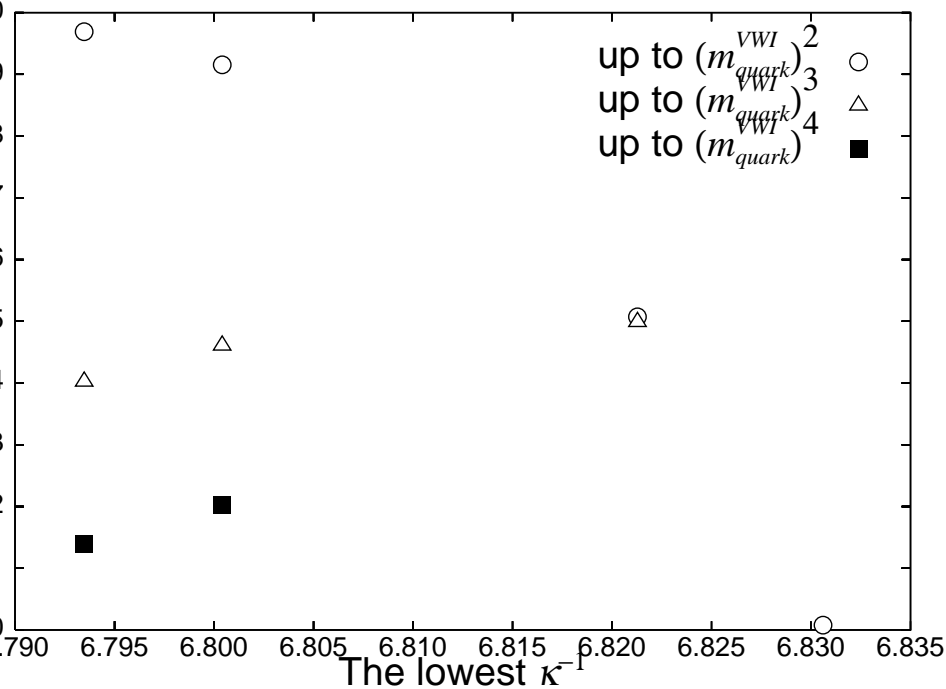
- Large deviations from the previous fit are also observed in small sea quark mass region (10σ difference in κ_c)
- The quadratic fit form can not describe our data of $m_{PS}/m_V = 0.80 - 0.35(\chi^2/dof = 8)$
cf. $m_{PS}/m_V = 0.80 - 0.55(\chi^2/dof = 0.1)$



3 How do the sea quark mass dependence

in m_{PS}^2 increase?

- With $m_{PS}/m_V = 0.55 - 0.80$ data, the quadratic fit works ($\chi^2/dof = 0.1$)
- As the smaller sea quark mass data is added to the fit, χ^2/dof increases and κ_c changes.

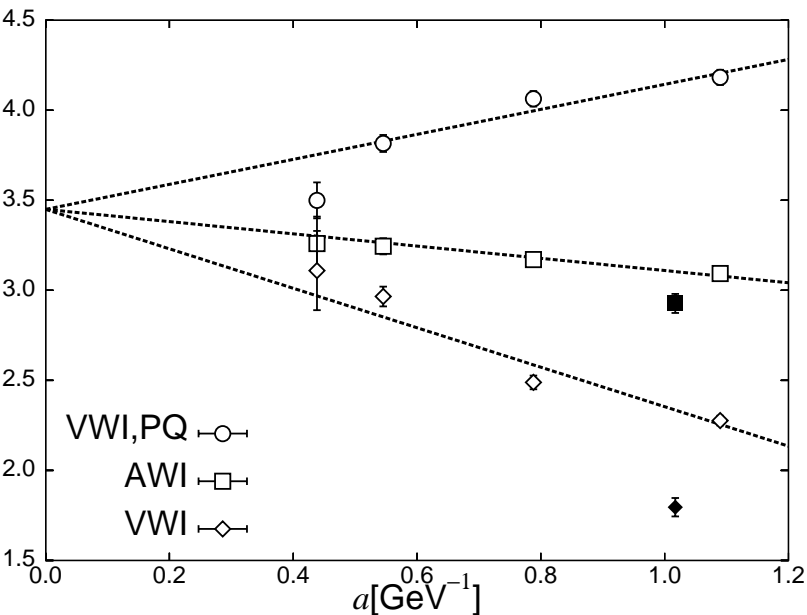


Results of polynomial extrapolations]

- Smaller sea quark mass data require higher order polynomials towards the chiral limit in addition to the quadratic term

quark masses in physical units from $m_{PS}/m_V = 0.80 - 0.35]$

- $m_{ud,R}^{AWI,\overline{\text{MS}}}(\mu = 2\text{GeV})$ decreases by **8%**(4σ)
- $m_{ud,R}^{VWI,\overline{\text{MS}}}(\mu = 2\text{GeV})$ decreases by **23%**(10σ)

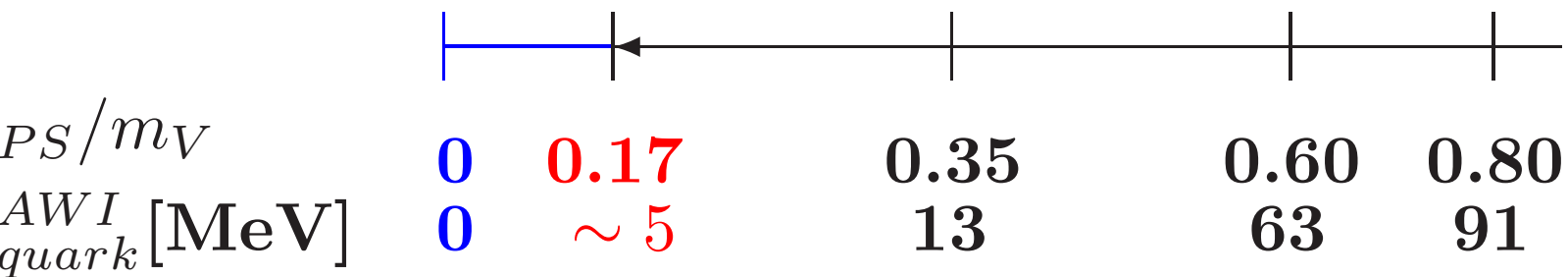


hiral extrapolations based on ChPT

- In the case of polynomial extrapolations, more higher order terms may be needed if smaller sea quark mass data are available
→ ChPT may give a guide for extrapolations, which is expected to describe the sea quark mass dependence around the physical point.

hPT : a low energy effective theory of QCD

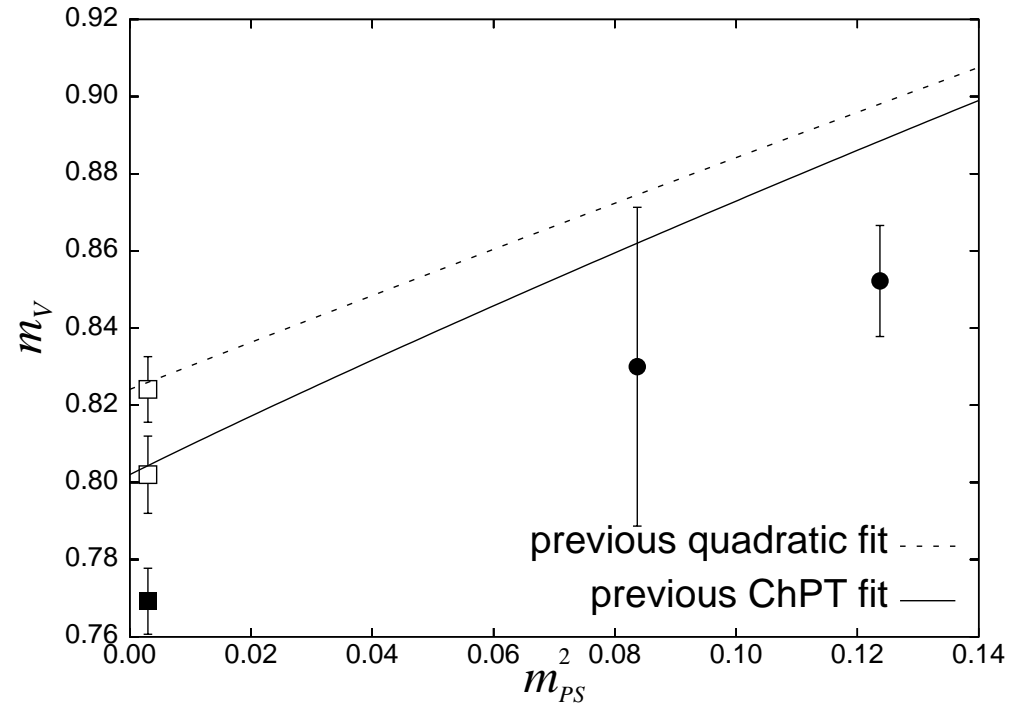
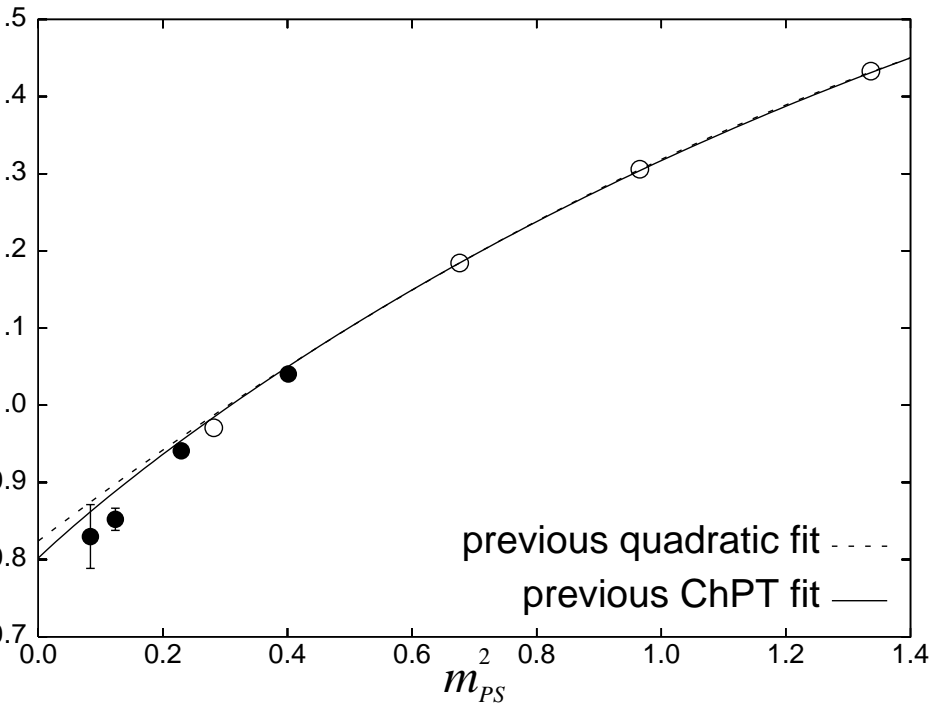
einberg,1979;Gasser and Leutwyler,1984,1985



vector meson mass with ChPT in the static limit]

$$m_V = A + Bm_{PS}^2 + Cm_{PS}^3 \quad \text{Jenkins et.al.,1995}$$

- ChPT fit with the previous data $m_{PS}/m_V = 0.80 - 0.55$ predicted the new $m_{PS}/m_V = 0.50 - 0.35$ data better than the quadratic fit



One-loop ChPT fit to m_{PS}^2 and f_{PS}

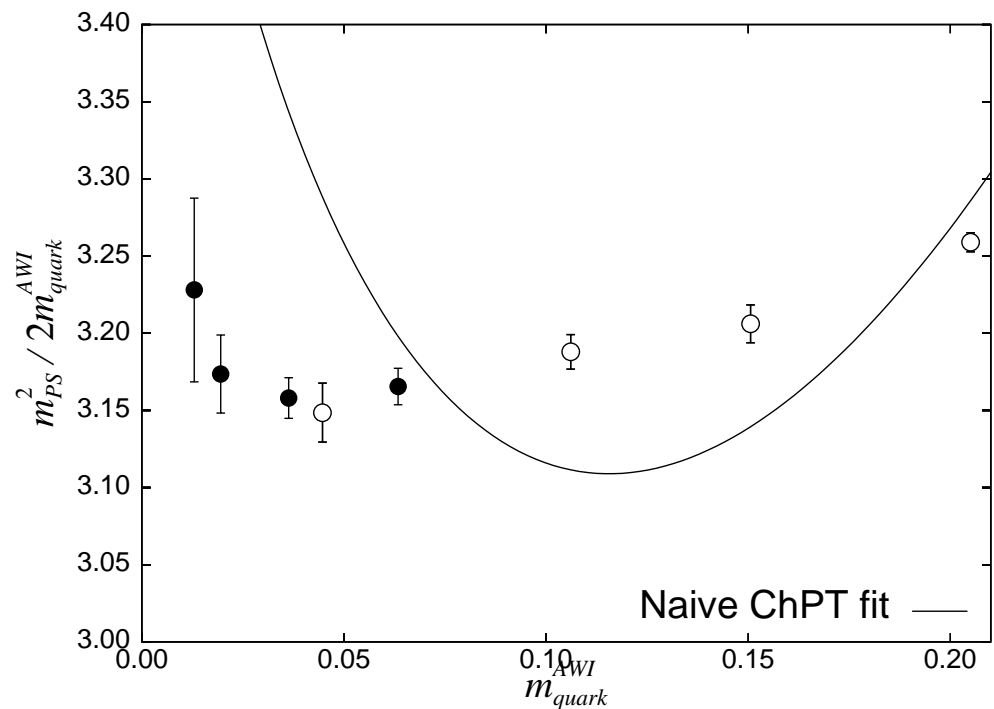
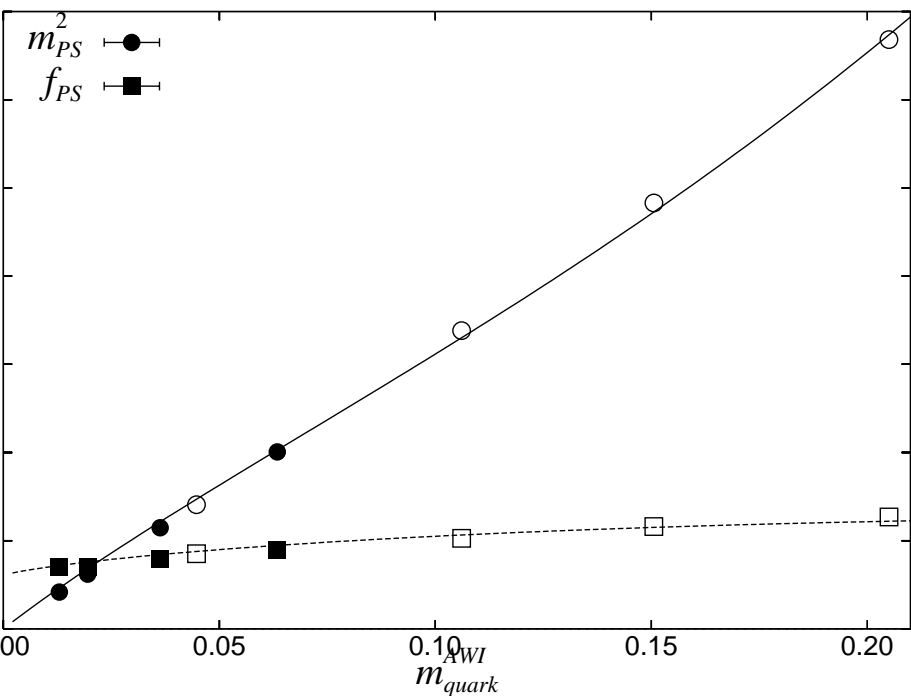
$$m_{PS}^2 = 2m_{quark}B_0 \left(1 + \frac{B_0m_{quark}}{(4\pi f)^2} \log \frac{2B_0m_{quark}}{\Lambda_3} \right)$$

$$f_{PS} = f \left(\frac{2B_0m_{quark}}{(4\pi f)^2} \log \frac{2B_0m_{quark}}{\Lambda_4} \right)$$

- Fitting parameters : $B_0, f, \Lambda_3, \Lambda_4$

ChPT fails to describe our data of $m_{PS}/m_V = 0.80 - 0.35$

[$m_{quark}^{AWI} : \chi^2/dof = 71$]



Reasons of failure of the naive ChPT formulae]

- Sea quark masses are still too heavy for the one-loop formulae (a fit dropping the heavy data $m_{PS}/m_V = 0.60 - 0.35$ gives $\chi^2/dof \sim 4$)
- **Chiral symmetry is explicitly broken for the Wilson-type fermions**
→ **Modify ChPT for Wilson-type fermions(WChPT)**
Sharpe and Singleton,1998; Rupak and Shoresh,2002;
O.Baer *et.al.*,2003; S.Aoki,2003

$$\mathcal{L}_{WChPT} = \frac{f^2}{4} \left(1 + c_0(S^0 - 1)\right) \mathbf{Tr} (\partial_\mu U \partial^\mu U) + c_1 S^0 + c_2 (S^0)^2$$

$$S_0 \equiv \mathbf{Tr}(U + U^\dagger), \quad U \equiv \exp \left(i \sum_{a=1}^3 \frac{\pi_a \sigma^a}{f} \right) \in SU(N_f = 2)$$

$$c_0 = W_0 a,$$

$$c_1 = W_1 a + B_1 m_{quark},$$

$$c_2 = W_2 a^2 + V_2 (m_{quark} a) + O(m_{quark}^2).$$

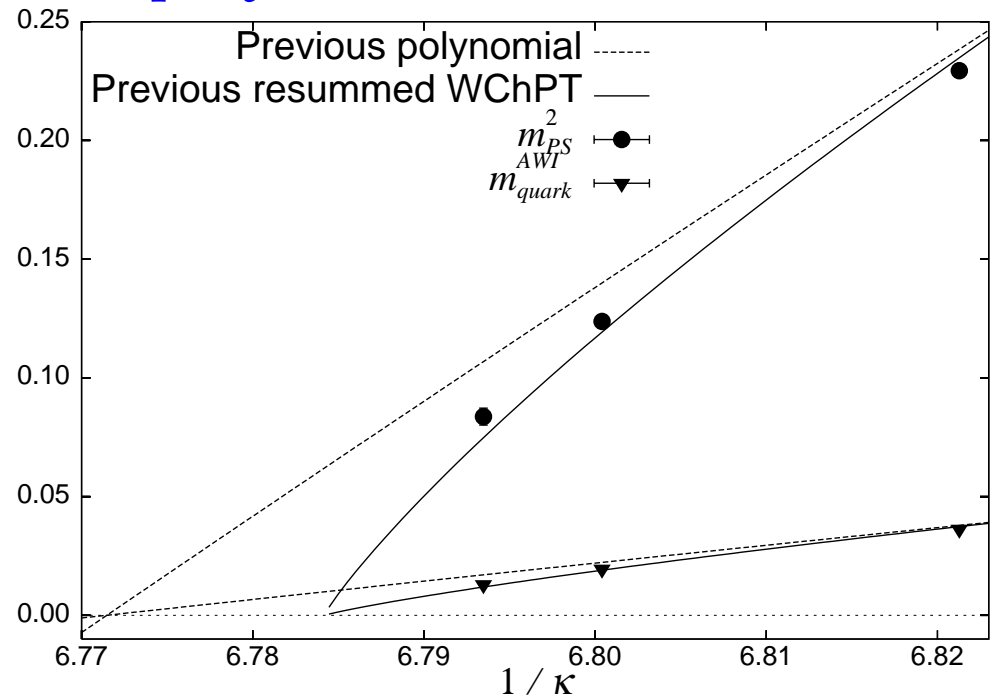
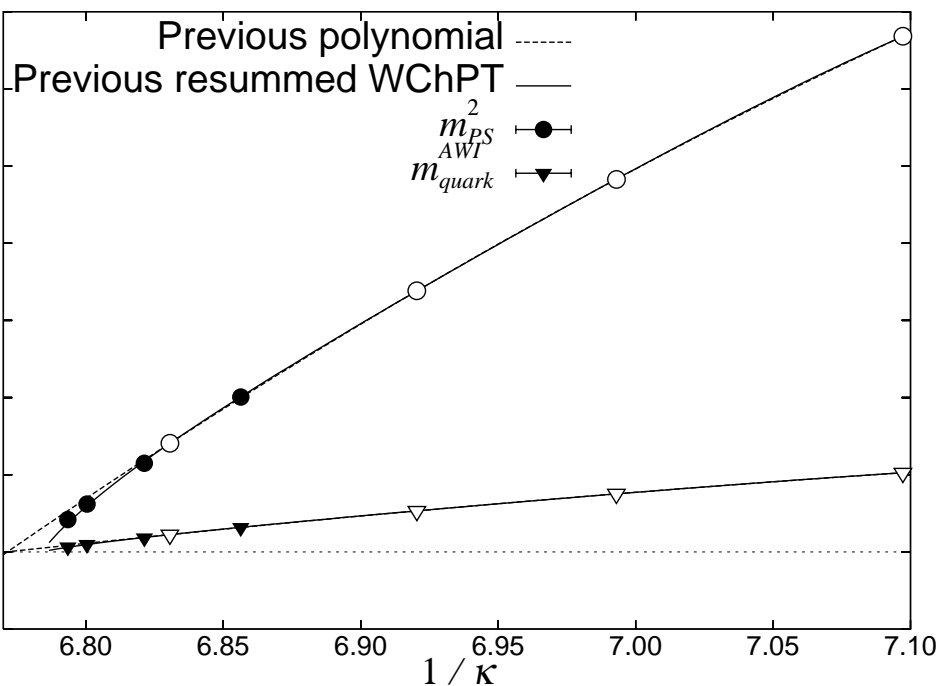
here $W = V = 0$ recovers the naive ChPT Lagrangian

Resummed WChPT formulae at one-loop] S.Aoki,2003

$$m_{PS}^2 = Am_{quark}^{VWI} \left(-\log \left(\frac{m_{quark}^{VWI}}{\Lambda_0} \right) \right)^{\omega_0} \left(1 + \omega_1^{PS} m_{quark}^{VWI} \log \left(\frac{Am_{quark}^{VWI}}{\Lambda_3^2} \right) \right)$$

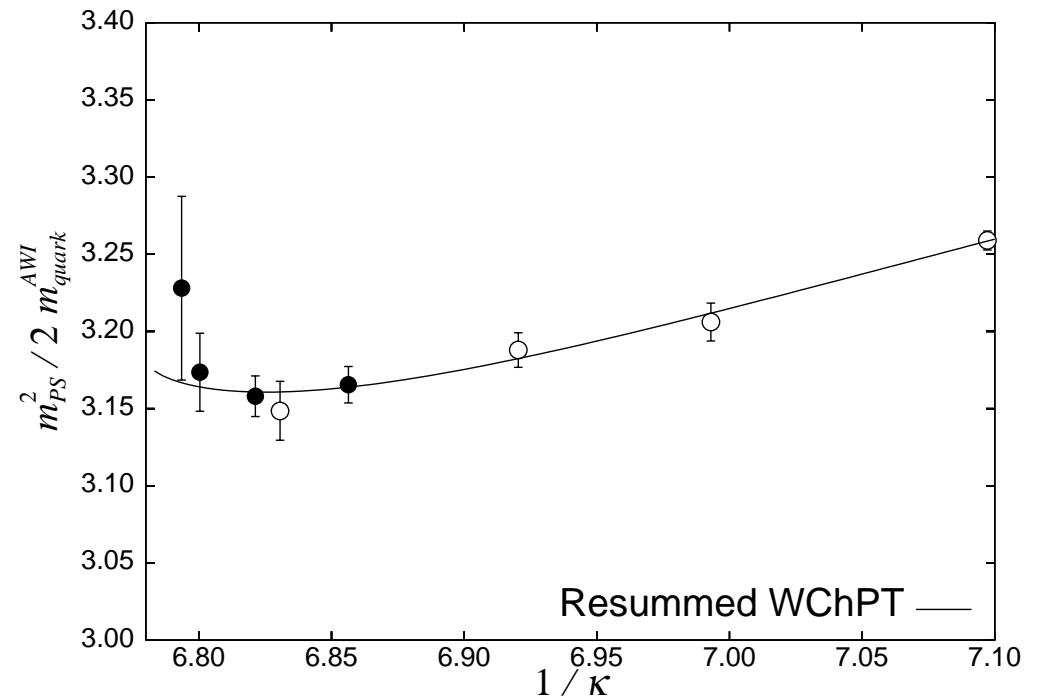
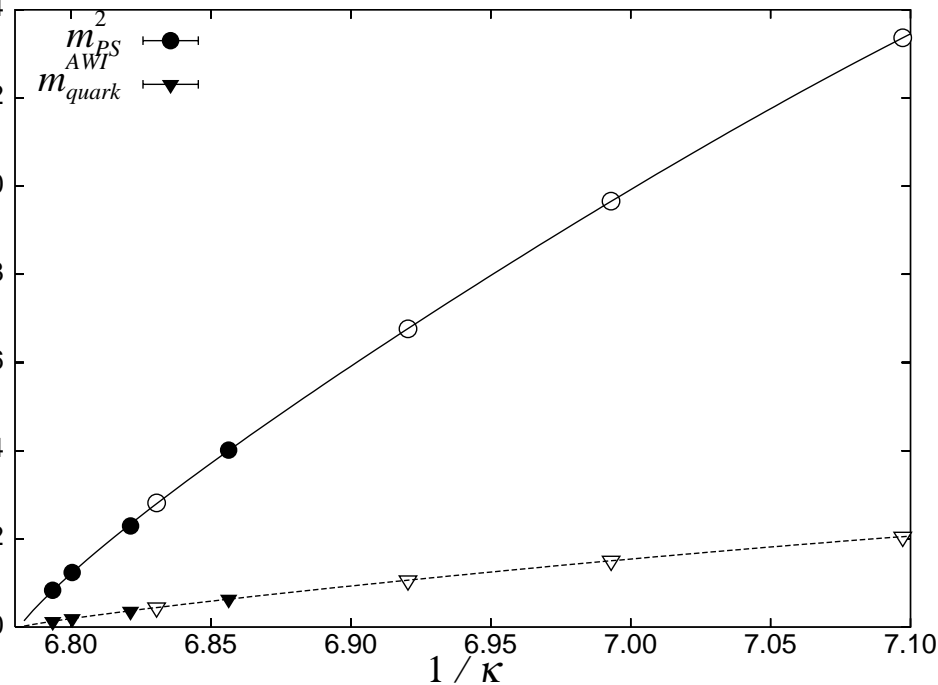
$$m_{quark}^{AWI} = m_{quark}^{VWI} \left(-\log \left(\frac{m_{quark}^{VWI}}{\Lambda_0} \right) \right)^{\omega_0} \left(1 + \omega_1^{AWI} m_{quark}^{VWI} \log \left(\frac{Am_{quark}^{VWI}}{\Lambda_{3,AWI}^2} \right) \right)$$

- The WChPT formulae with the previous data
 $m_{PS}/m_V = 0.80 - 0.55$ predicted the data of
 $m_{PS}/m_V = 0.50 - 0.35$ better than polynomials



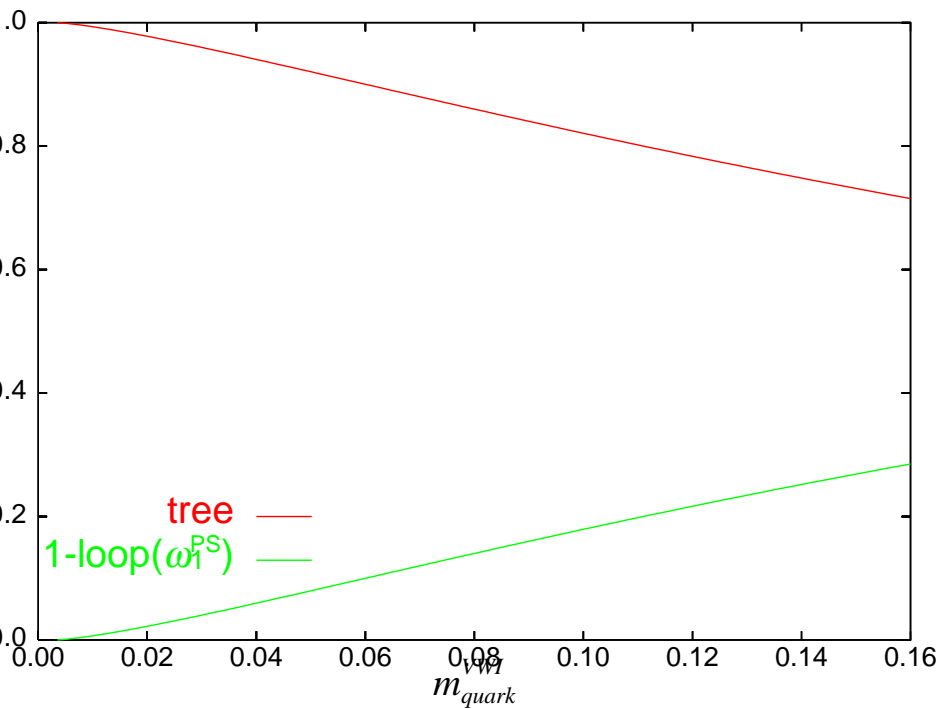
WChPT fit with our all data of $m_{PS}/m_V = 0.80 - 0.35]$

- The WChPT formulae describe our data well ($\chi^2/dof = 1.6$), in contrast to the naive ChPT case.
- $\Lambda_3 = 0.15[\text{GeV}]$ is comparable with the phenomenological estimate ($0.2 < \Lambda_3 < 2[\text{GeV}]$ **Gasser and Leutwyler, 1984**)
- $O(a)$ contribution ($\omega_1^{PS} - \omega_1^{AWI}$) suppresses the curvature in $m_{PS}^2/2m_{quark}^{AWI}$ to approximately 10% of the naive ChPT



Convergence check of the resummed WChPT]
at the heaviest point,

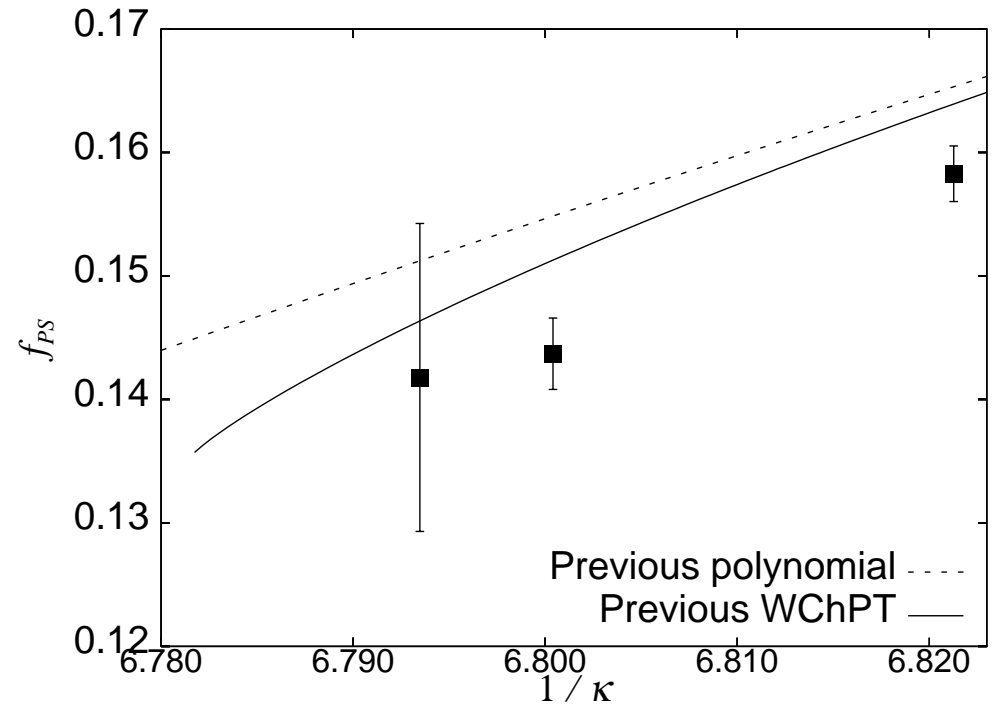
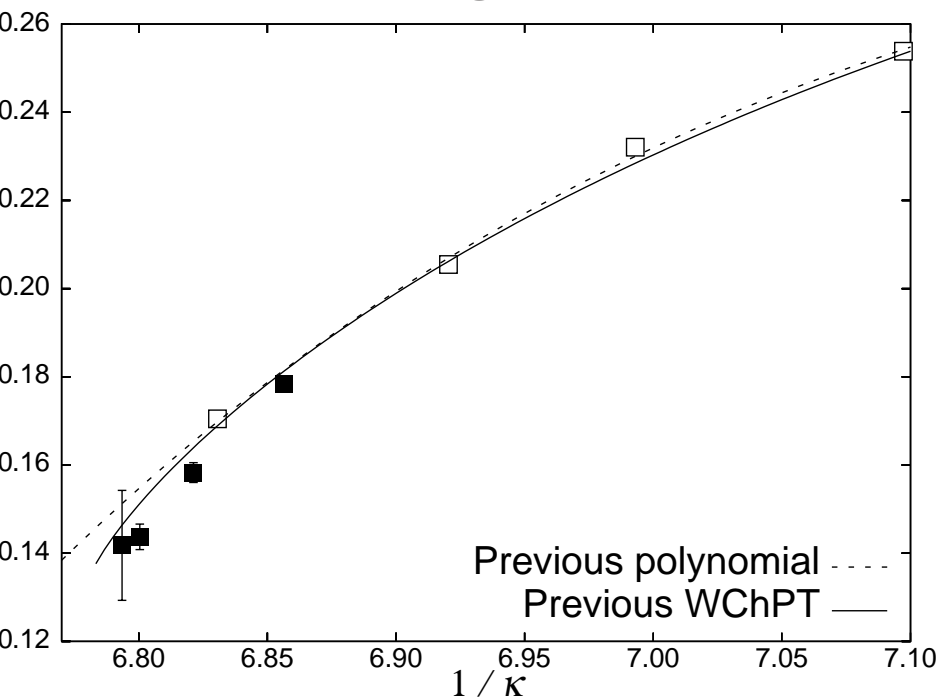
- Tree : 71%
- one-loop : 29%



WChPT formula for the pseudoscalar decay constant]

$$f_{PS} = f \left(1 - \omega_1^{f_{PS}} m_{quark}^{VWI} \log \left(\frac{A m_{quark}^{VWI}}{\Lambda_4^2} \right) \right)$$

- **WChPT formula with the previous data $m_{PS}/m_V = 0.80 - 0.55$ predicted $m_{PS}/m_V = 0.50 - 0.35$ data better than polynomial**
- **The WChPT describes our data of $m_{PS}/m_V = 0.80 - 0.35$ ($\chi^2/dof = 3.6$)**
- **$\Lambda_4 = 2.44(13)$ is comparable with (a little larger than) the phenomenological estimate ($\Lambda_4 = 1.26 \pm 0.14$ [GeV] **Colangelo *et al.*, 2001**)**
 → **The scaling violation leads to the deviation!?**



Conclusion

We studied light hadron spectrum using several chiral extrapolations with our data of

$m_{PS}/m_V = 0.80 - 0.35$. Our results imply that

it is better to employ WChPT formulae instead of conventional polynomials for chiral extrapolations.

[Reasons]

- WChPT has a concrete theoretical background
- WChPT formulae from $m_{PS}/m_V = 0.80 - 0.55$ predicted $m_{PS}/m_V = 0.50 - 0.35$ data better than polynomials
- WChPT formulae describe our data of $m_{PS}/m_V = 0.80 - 0.35$ in contrast to the naive ChPT formulae
- Convergence of one-loop WChPT formulae is reasonable

[Prediction by WChPT formulae]

$$m_{ud,R}^{VWI,AWI}(\overline{\text{MS}}, \mu = 2\text{GeV}) = 1.230(94), 2.725(38) [\text{MeV}]$$

$$m_{ud,R}^{VWI,AWI}(\overline{\text{MS}}, \mu = 2\text{GeV}) = 1.749(53), 2.851(58) [\text{MeV}] \quad \text{cf. polynomial}$$

future works]

- Check the validity of WChPT formulae (smaller sea quark masses, scaling)
- WChPT formulae for vector mesons and baryons

3.4 [Finite size effect]

$L = 2.4\text{fm}$ vs $L = 3.2\text{fm}$

- No finite size effects in the meson quantities
($m_{PS}, m_V, m_{quark}^{AWI}$)
- $1 - 3\%$ ($0.8 - 2.3\sigma$) decreases
in the baryon quantities (m_N, m_Δ)

