

Toward $O(a)$ -improvement of axial current in lattice QCD

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for the ALPHA/CP-PACS/JLQCD collaborations

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I. Introduction

CP-PACS/JLQCD project in three-flavor QCD

(T.Ishikawa's talk for details)

- RG-improved gauge + NP $O(a)$ -improved Wilson quark
- use all machines available : 2.3 TFLOPS in total

CP-PACS(0.3), SR8K/G1(0.2), VPP-5000(0.3) @ Univ.of Tsukuba

SR8K/F1(0.8) @ KEK

Earth Simulator (0.8) @ JAMSTEC

- one-loop value for c_A , $Z_{\{A,P,m\}}$, $b_{\{A,P,m\}}$

↓

- $m_q \Rightarrow$ fundamental parameter : $O(ag^4)$ scaling violation

↑

NP determination of c_A , $Z_{\{A,P,m\}}$, $b_{\{A,P,m\}}$

ALPHA

- $O(a)$ -improvement
- plaquette gauge + NP $O(a)$ -improved Wilson quarks
- non-perturbative determination of improvement and renormalization constants in quenched and two-flavor QCD



collaboration in computation of c_A , ... in full QCD

ALPHA : c_A in two-flavor QCD → talk by Della Morte

CP-PACS/JLQCD : a/L dependence of c_A in quenched QCD
with plaquette gauge action → this talk

contents :

1. Introduction
2. Topics on determination of c_A
3. Old method
4. New method
5. Summary

II. Topics on determination of c_A

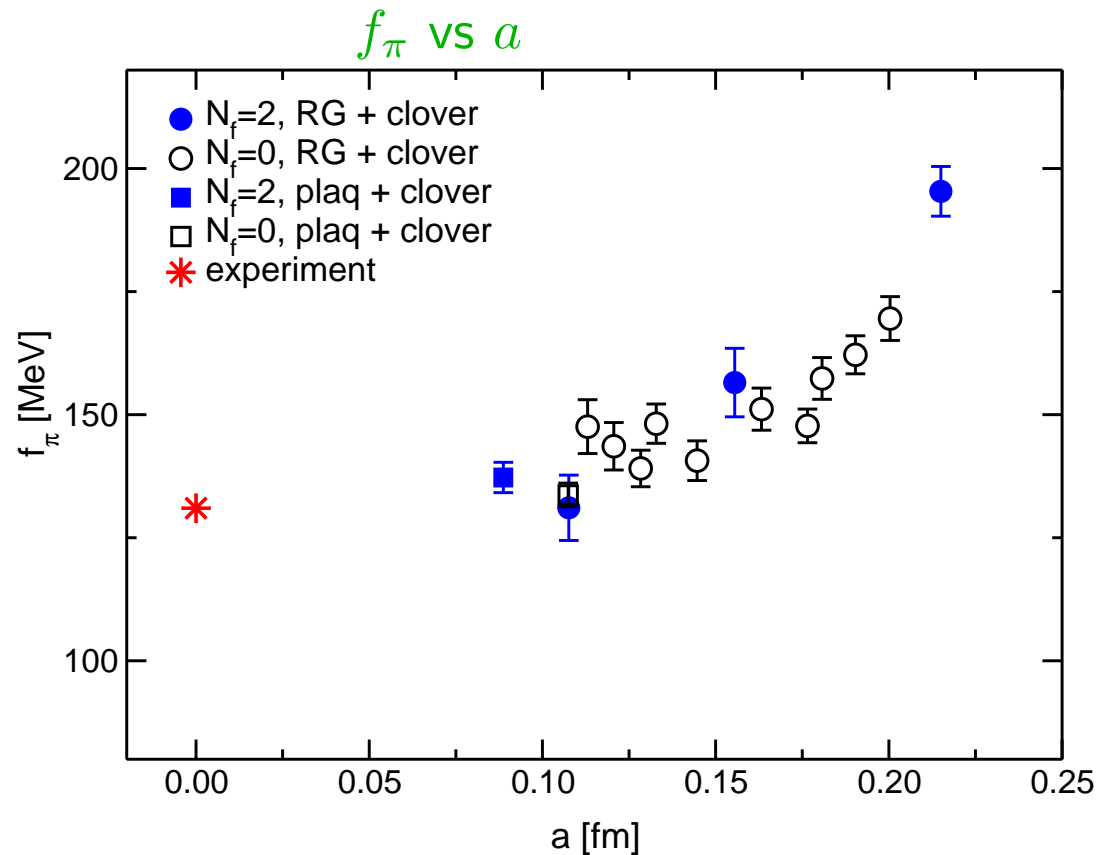
contents :

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II-1 motivation

CP-PACS study in two-flavor QCD

f_{PS} (using one-loop c_A, Z_A, b_A) \Rightarrow large scaling violation



- scaling violation :

$$\sim O(a)$$

\Rightarrow NP determination of c_A
(or $O(a^2)$?)

- renormalization factor Z_A

\Rightarrow talk by K.Ide

JLQCD data in quenched QCD

- with plaq. gauge + NP $O(a)$ -improved Wilson quarks
- AWI $m_q = \langle A_4 P \rangle / (2 \langle PP \rangle)$

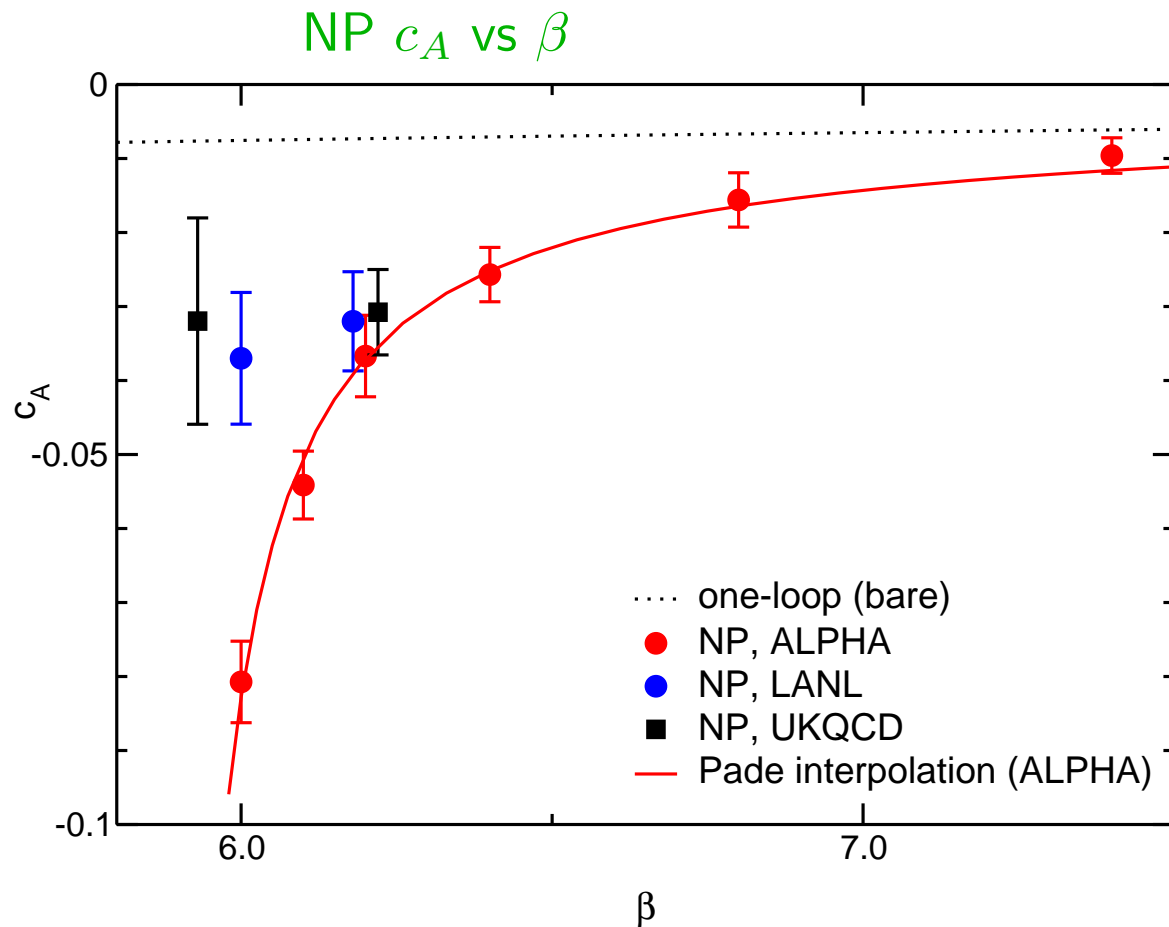
$$m_s^{\text{AWI}} = \begin{cases} 104.1(1.6) \text{ MeV} & \text{with one-loop } c, Z, b \\ 91.9(1.4) \text{ MeV} & \text{with NP } c, Z, b \text{ by ALPHA} \end{cases}$$

\Rightarrow 12% difference

- $(Z_A/Z_P)_{\text{one-loop}} = 1.19 \quad \Leftrightarrow (Z_A/Z_P)_{\text{NP}} = 1.22$
- $(b_A - b_P)_{\text{one-loop}} = -0.013 \quad \Leftrightarrow (b_A - b_P)_{\text{NP}} = 0.171$: but small $O(am_q)$
- $c_{A,\text{one-loop}} = -0.013 \quad \Leftrightarrow c_{A,\text{NP}} = -0.083$

II-2 $O(a)$ effects in c_A

previous results of NP c_A in quenched QCD



⇒ $O(a)$ effects in c_A is not small ?

⇒ large difference in $O(a^2)$ scaling property?

⇒ study of a/L dependence is important

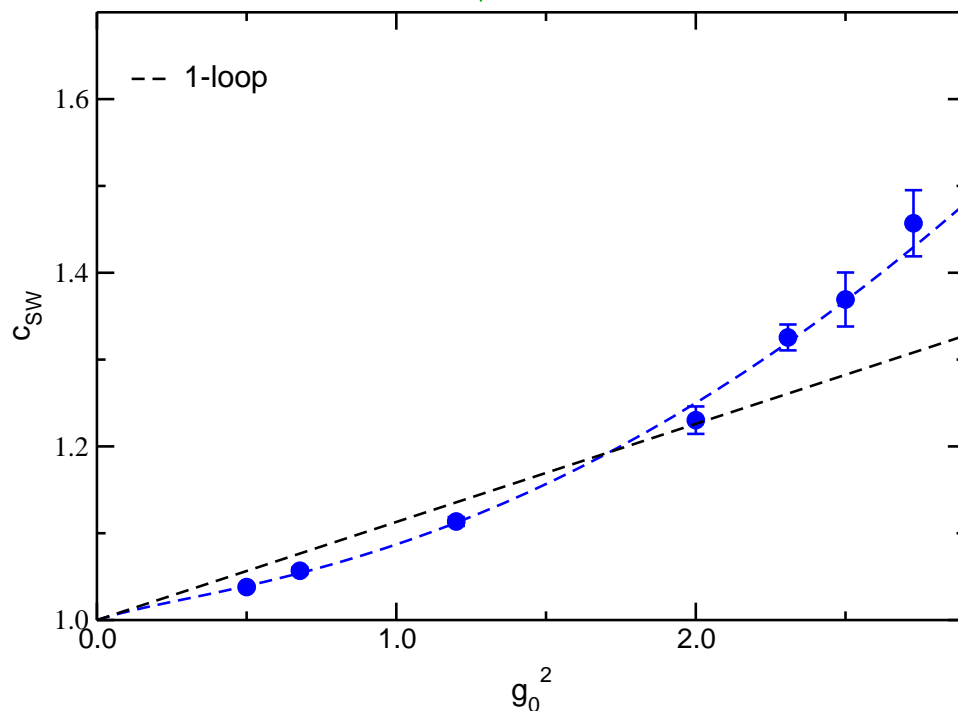
CP-PACS/JLQCD determination of c_{SW}

- in three-flavor QCD with RG improvd gauge

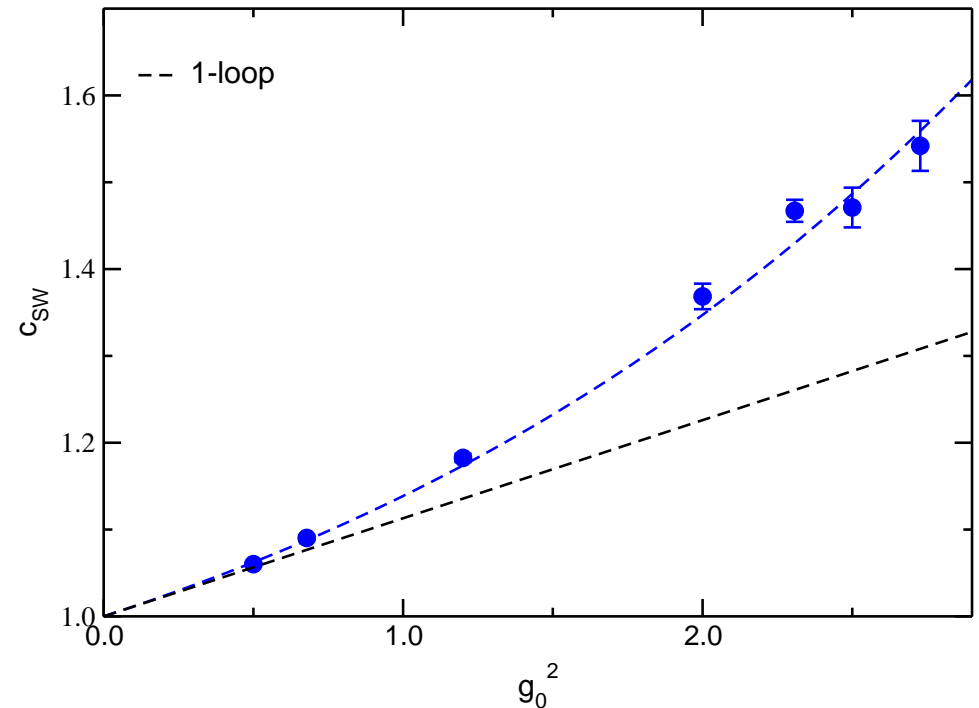
⇒ non-negligible a/L dependence of c_{SW}

⇒ correct NP c_{SW} to $L=L^*$ ($L^* = 6a(\beta = 1.9)$)

$L/a = 8$



interpolated to $L=L^*$



III. Old method

contents :

1. Introduction
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III-1 Method

method to determine c_A

1) calculate quark masses with two values of θ

$$\psi(x + L\hat{k}) = e^{i\theta_k} \psi(x), \quad \bar{\psi}(x + L\hat{k}) = \bar{\psi}(x) e^{-i\theta_k}$$

$$f_A(x_0) = -\frac{1}{3} \langle A_0^a \mathcal{O}^a \rangle, \quad f_P(x_0) = -\frac{1}{3} \langle P^a \mathcal{O}^a \rangle$$

$$m(\theta, x_0) = r(\theta, x_0) + c_A s(\theta, x_0),$$

$$r(x_0) = \frac{1}{4} (\partial_0 + \partial_0^*) f_A(x_0) / f_P(x_0), \quad s(x_0) = \frac{1}{2} a \partial_0 \partial_0^* f_P(x_0) / f_P(x_0),$$

2) determine c_A by imposing that quark masses coincide with each other

$$m(\theta', x_0) - m(\theta, x_0) = \Delta m^{(0)},$$

where we take $\theta = 0$, $\theta' = 1$, $x_0 = T/2$

simulation parameters

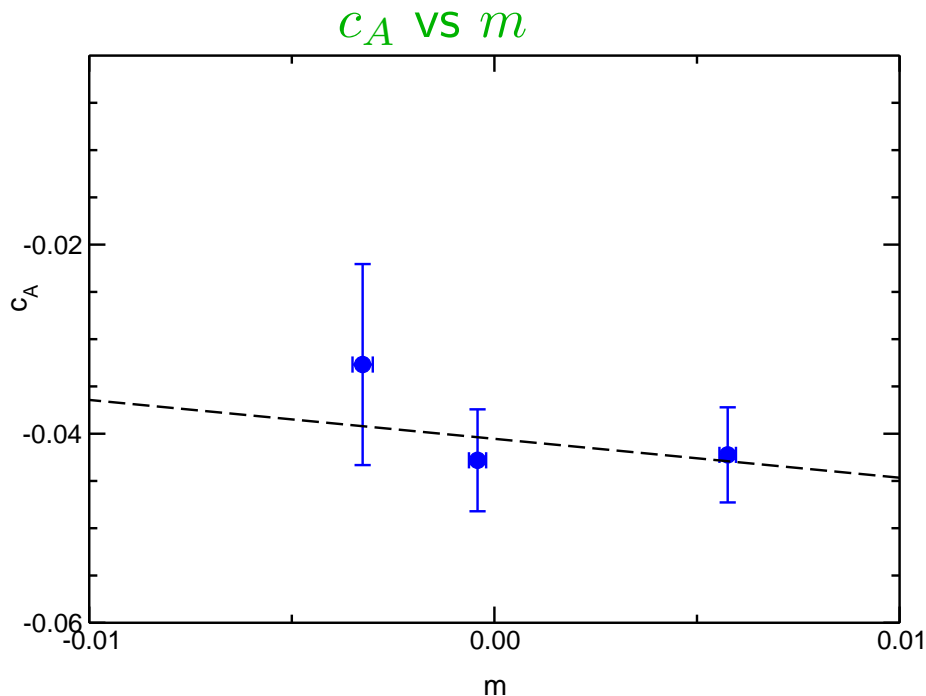
- in quenched QCD with plaq. gauge + $O(a)$ -improved Wilson quarks
- use HMC algorithm on SR8K at KEK
- take NP c_{SW} from ALPHA
- take K_c on $16^3 \times 32$ from ALPHA
 - exceptional conf \Rightarrow heavier mass at $\beta = 6.0$
 - mass dependence of c_A is checked at $\beta = 12.0$ and 6.2 on $8^3 \times 16$

statistics [traj]

	$6^3 \times 12$	$8^3 \times 16$	$10^3 \times 20$	$12^3 \times 24$	$14^3 \times 28$
$\beta = 12.0$	40000	40000	37000	40000	40000
8.0	60000	35000	40000	48000	40000
6.8	80000	45000	–	40000	40500
6.4	–	55000	–	60000	–
6.2	100000	80000	–	80000	–
6.0	–	40000	–	–	–

III-2 Mass dependence of c_A

at $\beta=6.2$ on $8^3 \times 16$

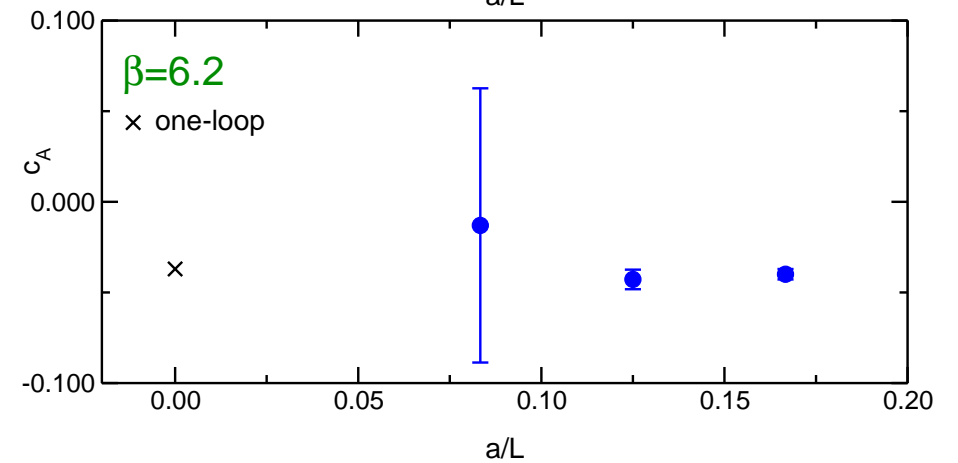
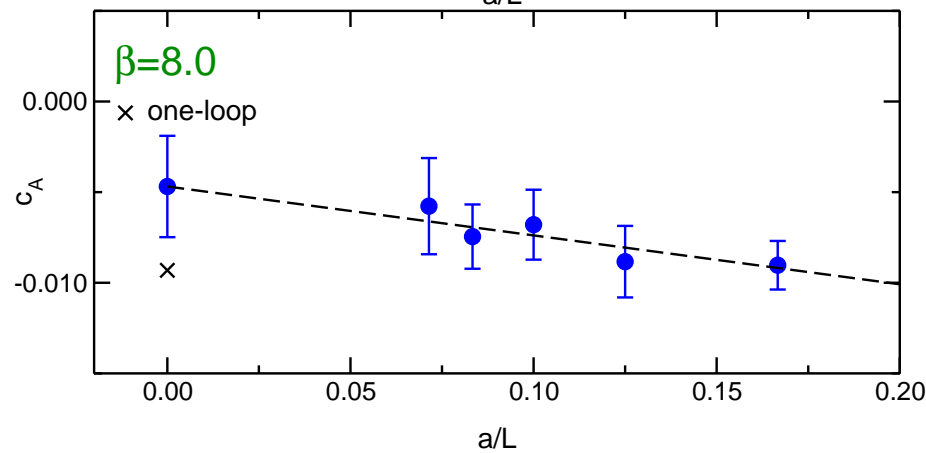
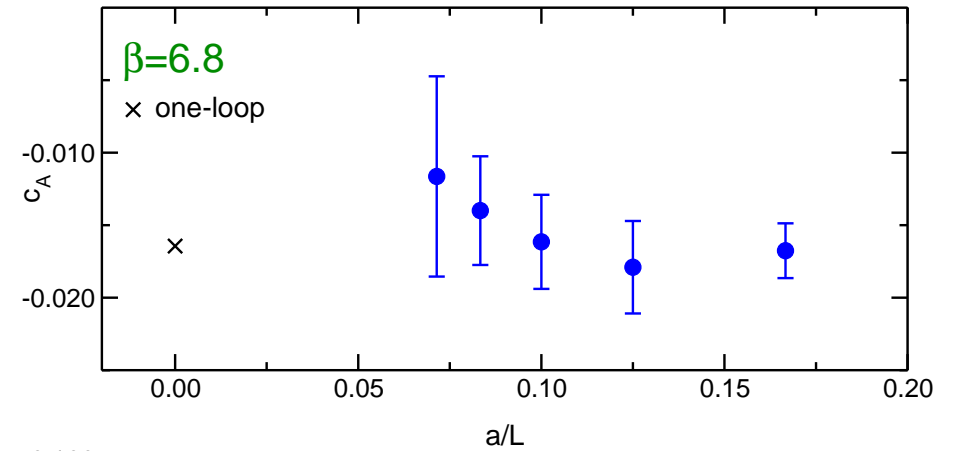
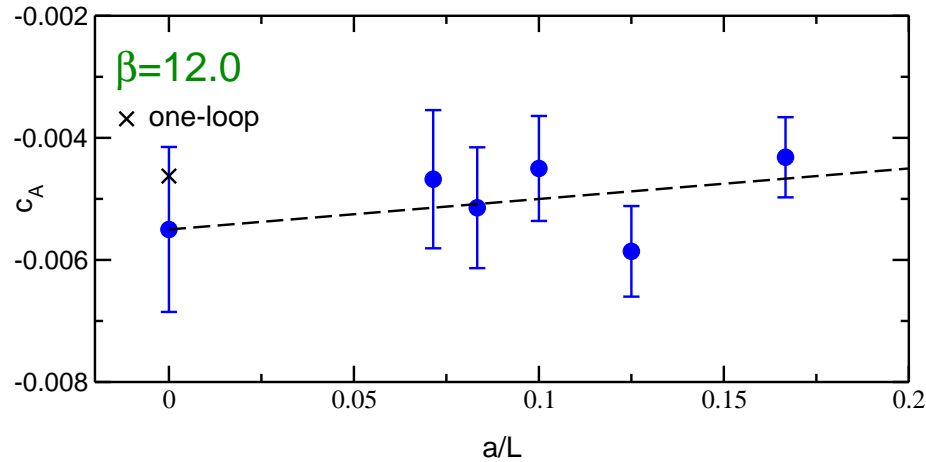


- small m dependence of c_A
 - ⇐ confirm ALPHA's observation
- $K_c(16^3 \times 32)$ is close to $K_c(8^3 \times 16)$
 - $m = -0.00042(21)$ at $\beta=6.2$ on $8^3 \times 16$
 - $m = -0.000050(34)$ at $\beta=12.0$ on $8^3 \times 16$

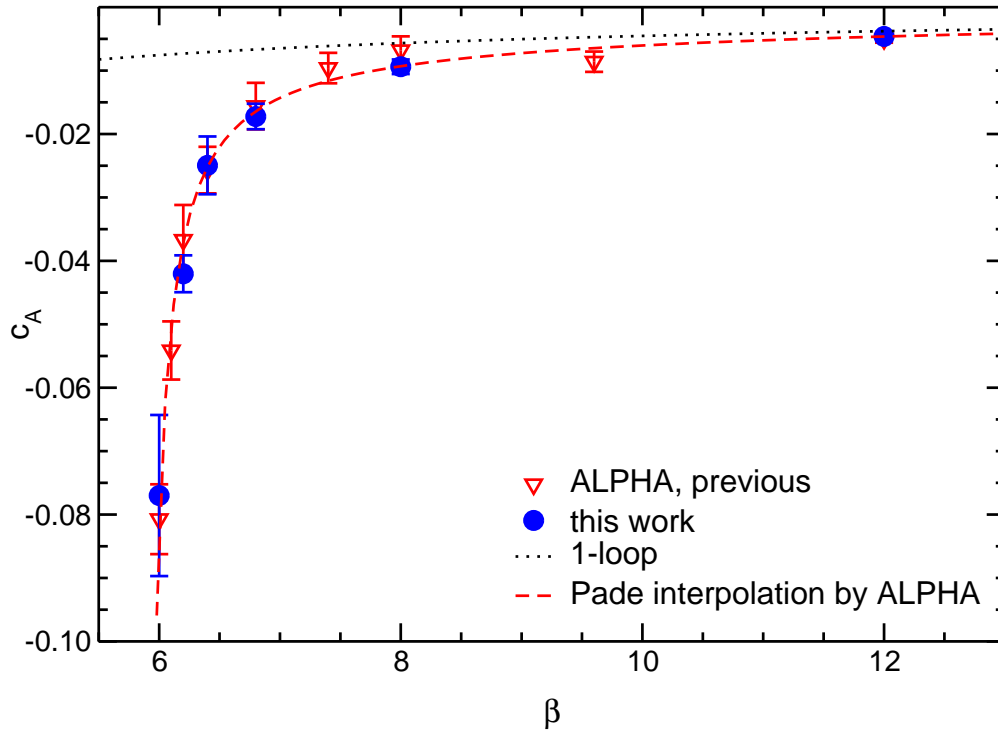
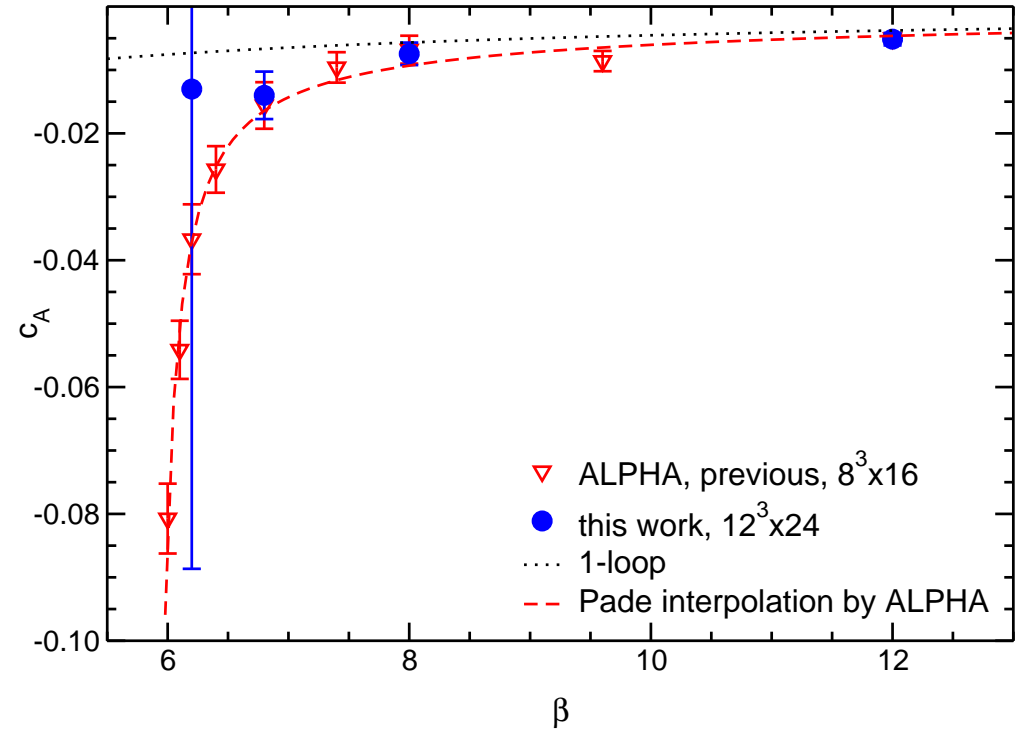
⇒ ignore m dependence of c_A in the following

III-3 volume dependence of c_A

c_A vs a/L



error of c_A increases for larger L/a

c_A vs β on $8^3 \times 16$ on $12^3 \times 24$ large error obscures β dependence

sensitivity

$$c_A = \frac{\Delta m^{(0)} - (r(\theta', x_0) - r(\theta, x_0))}{s(\theta', x_0) - s(\theta, x_0)}$$

larger $L \Rightarrow \theta$ has smaller effects ($p_k = \theta_k/L$)

\Rightarrow smaller difference between $s(\theta')$ and $s(\theta)$

$\Rightarrow c_A$ loses its sensitivity to the improvement condition

IV. New method

contents :

1. Introduction
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IV-1 Method

“slope method” to determine c_A

- improvement condition

$$m' - m = \Delta m^{(0)}$$

old method : difference between $\theta \Leftrightarrow \theta' = 1$ at fixed $x_0 (= T/2)$

new method : **difference between $x_0 \Leftrightarrow x_{0,\text{ref}}$ with fixed $\theta (= \pi)$**

\Rightarrow simulations with a single value of θ

- choice of x_0 and $x_{0,\text{ref}}$

x_0 : contribution of the ground state dominates ($= T/2$)

$x_{0,\text{ref}}$: contribution of the excited state dominates ($= T/4$)

simulation parameters

- in quenched QCD with plaq. gauge $+O(a)$ -improved Wilson quarks
- use HMC algorithm on SR8K at KEK
- take NP c_{SW} from ALPHA
- take K_c on $16^3 \times 32$ from ALPHA
 - exceptional conf \Rightarrow heavier mass at $\beta = 6.0$ and 6.2
 - mass dependence of c_A is checked at $\beta = 12.0$ and 6.4
- $T/L=2$
 - $T/L=3$ is tested at $\beta = 12.0$

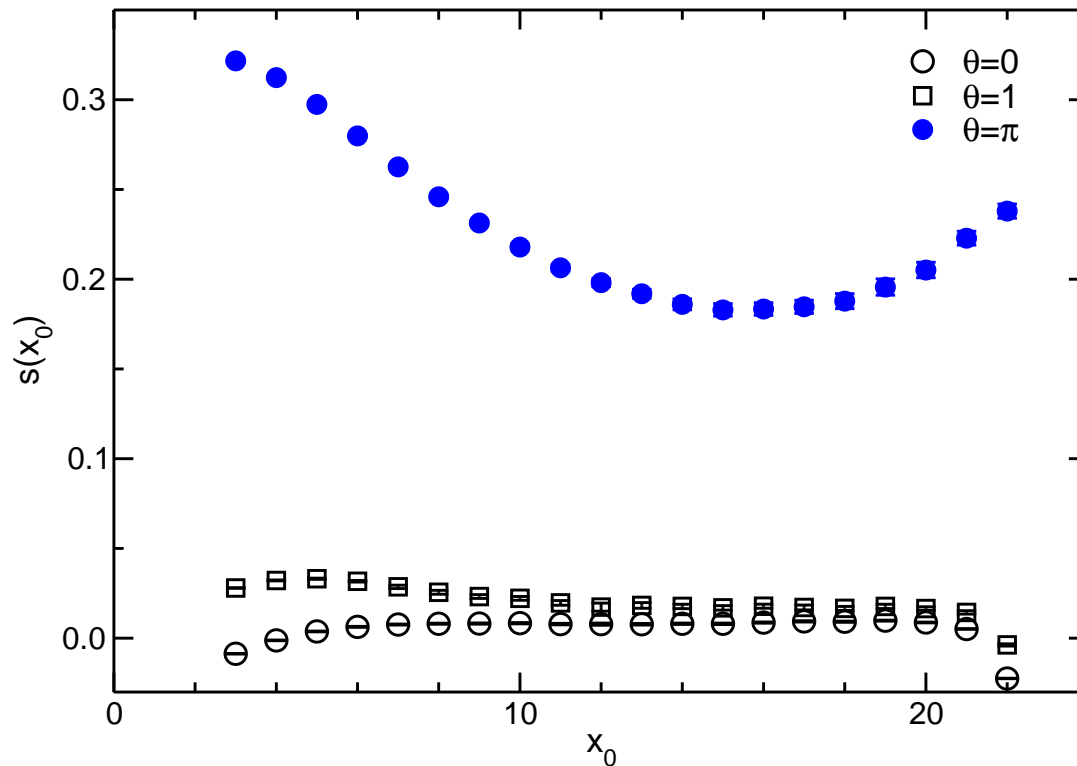
statistics [traj]

	$\beta = 12.0$	8.0	6.8	6.4	6.2	6.0
$8^3 \times 16$	40000	40000	60000	40000	30000	60000
$12^3 \times 24$	18000	80000	40000	40000	38000	46000
$16^3 \times 32$	30000	40000	40000	40000	40000	40000

IV-2 Results of s and c_A

$s(x_0)$ on $12^3 \times 24$

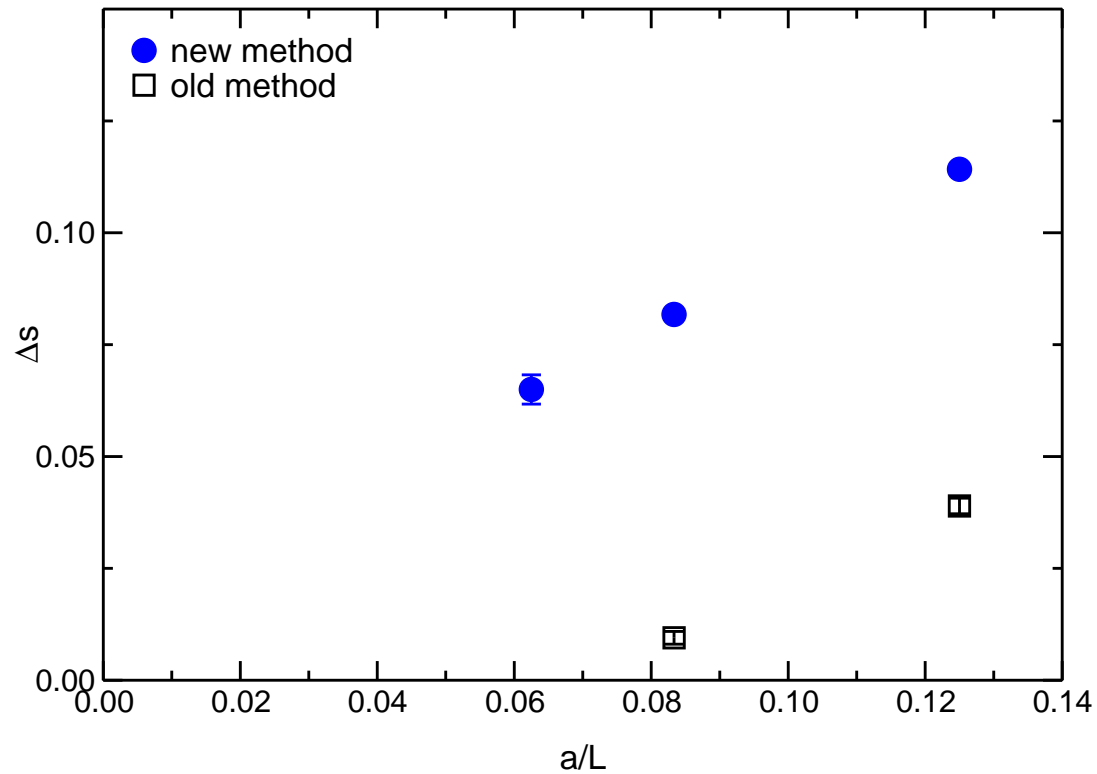
$\beta=6.4, c_{SW}=1.52562, 12^3 \times 24, K=0.13572$



- small difference between $s(T/2)$ with $\theta = 0$ and 1
- much larger difference between $s(T/2)$ and $s(T/4)$ with $\theta = \pi$

$$c_A = \frac{\Delta m^{(0)} - (r' - r)}{s' - s}$$

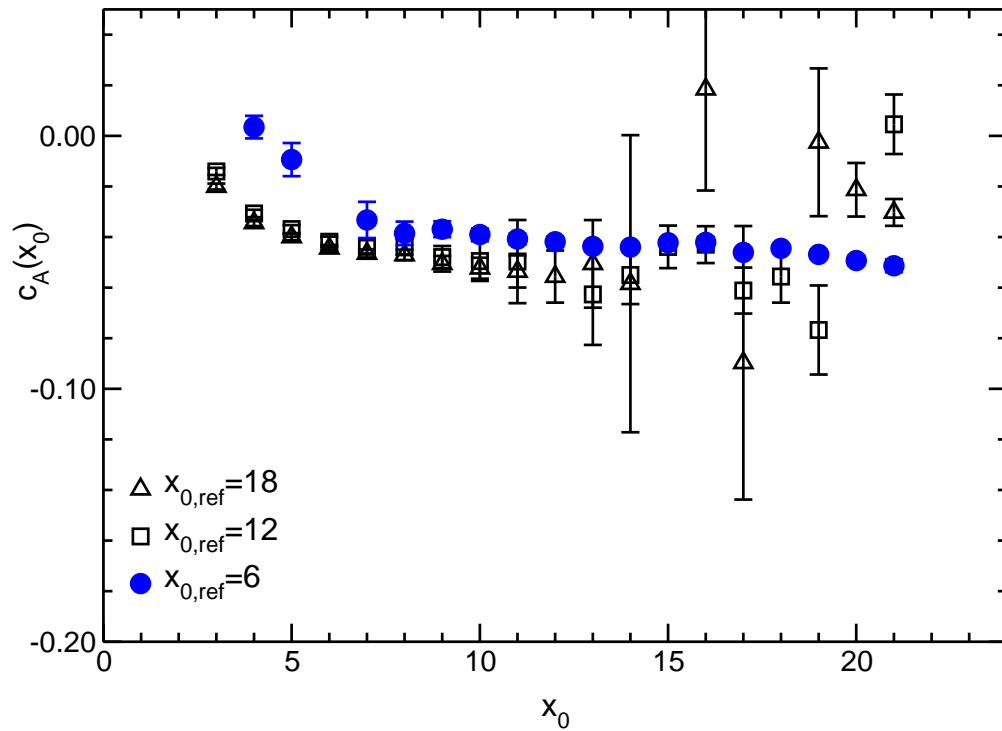
Δs at $\beta=6.4$



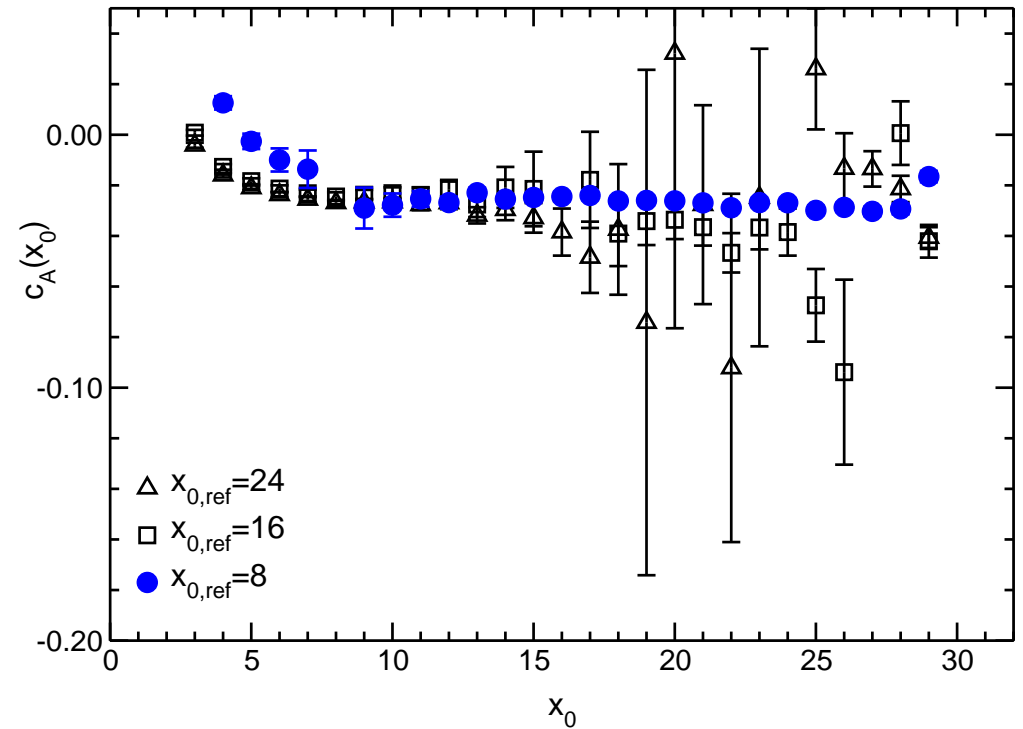
- old method : Δs almost vanishes at $L/a \sim 12$
- new method : much larger Δs

$c_A(x_0)$ at $\beta = 12.0$

$N_f=0$, $12^3 \times 24$, $\beta=12.0$, $c_{SW}=1.16373$, $K=0.12991$



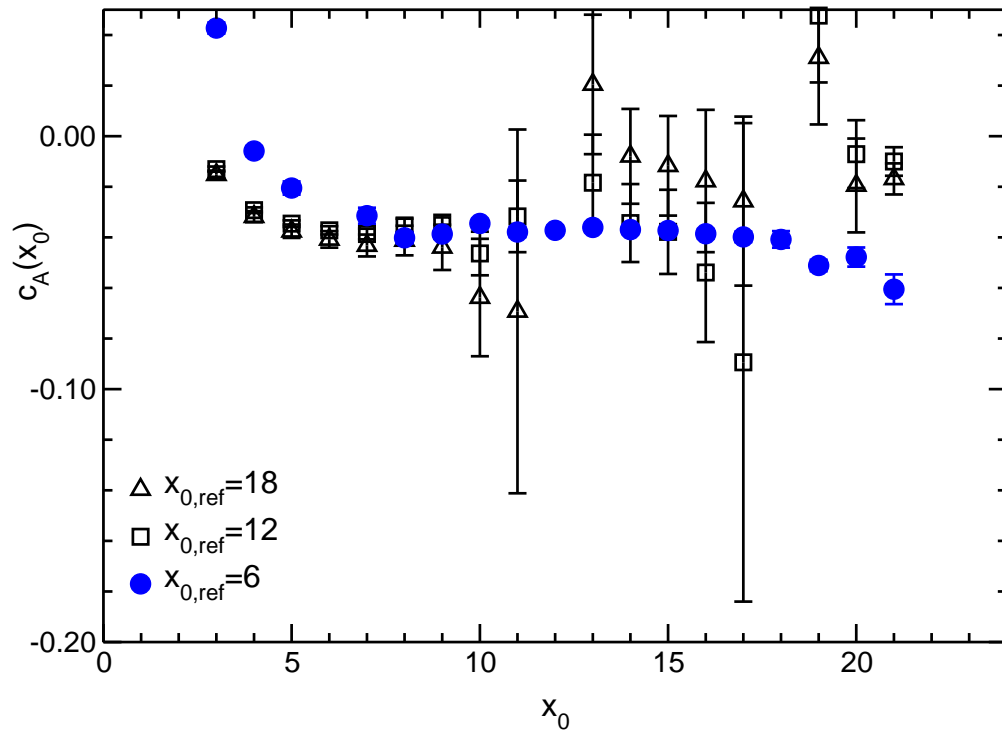
$N_f=0$, $16^3 \times 32$, $\beta=12.0$, $c_{SW}=1.16373$, $K=0.12991$



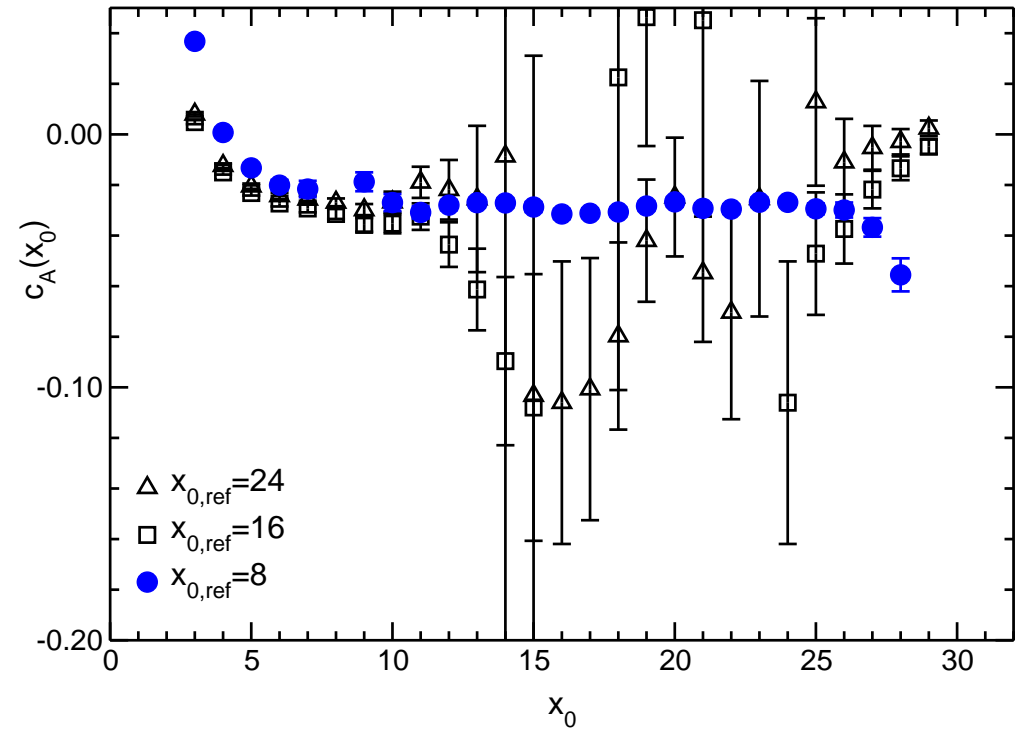
- x_0, x_0, ref dependence of c_A is not large
- take $x_0 = T/2$, $x_0, \text{ref} = T/4$

$c_A(x_0)$ at $\beta = 6.2$

$N_f=0$, $12^3 \times 24$, $\beta=6.2$, $c_{SW}=1.61375$, $K=0.13485$



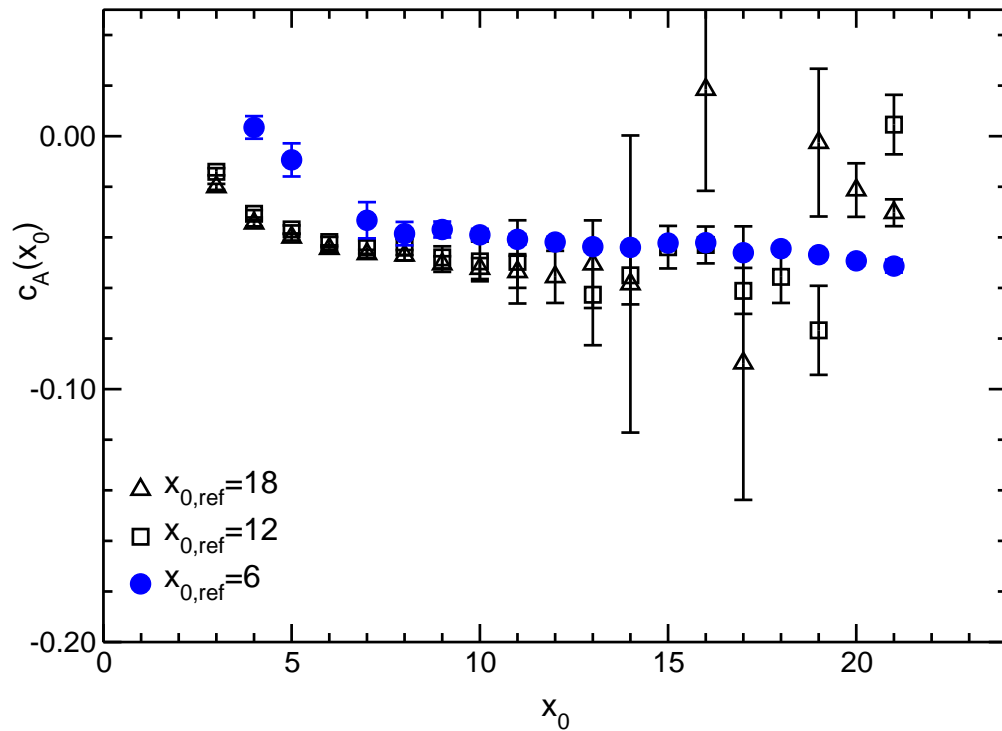
$N_f=0$, $16^3 \times 32$, $\beta=6.2$, $c_{SW}=1.61375$, $K=0.13485$



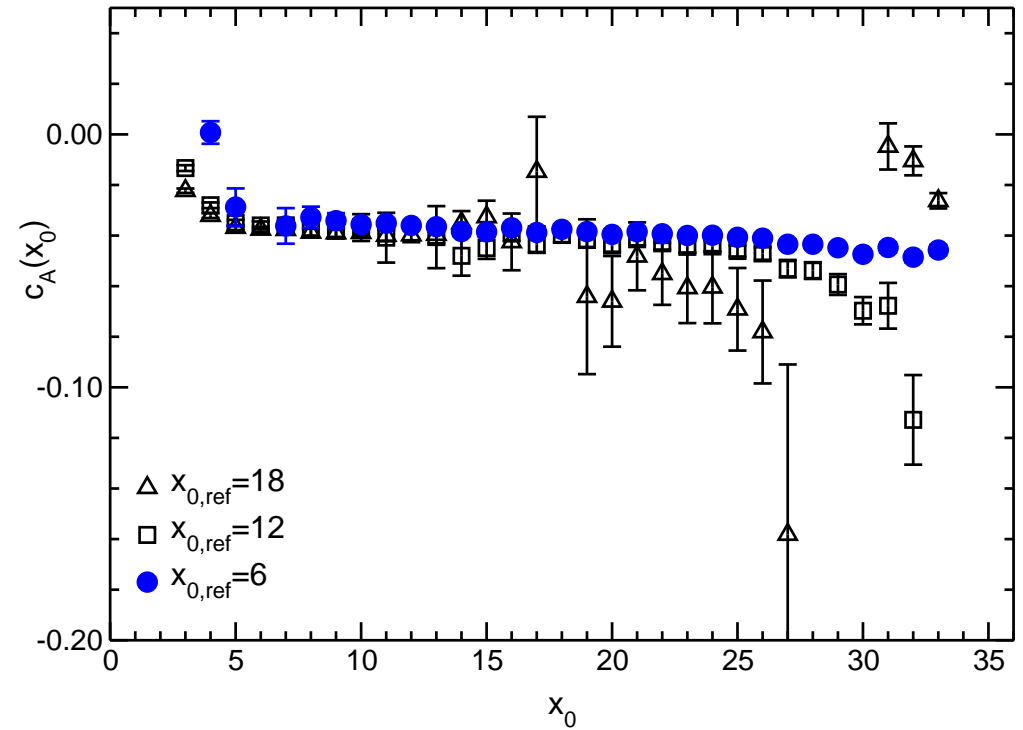
- exceptional conf appears at $\beta = 6.4$ and $6.2 \Rightarrow$ filtered out

$T/L = 2$ or $T/L = 3$

$N_f=0$, $12^3 \times 24$, $\beta=12.0$, $c_{sw}=1.16373$, $K=0.12991$



$N_f=0$, $12^3 \times 36$, $\beta=12.0$, $c_{sw}=1.16373$, $K=0.12991$

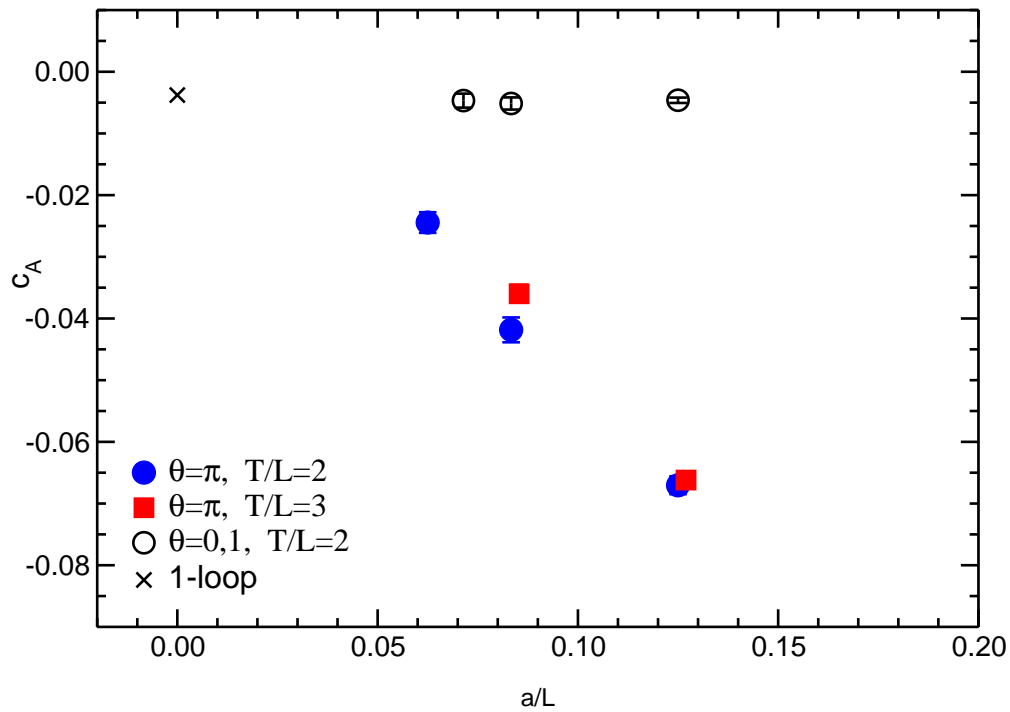


- give consistent c_A with each other
 \Rightarrow take $T/L = 2$ to reduce CPU time

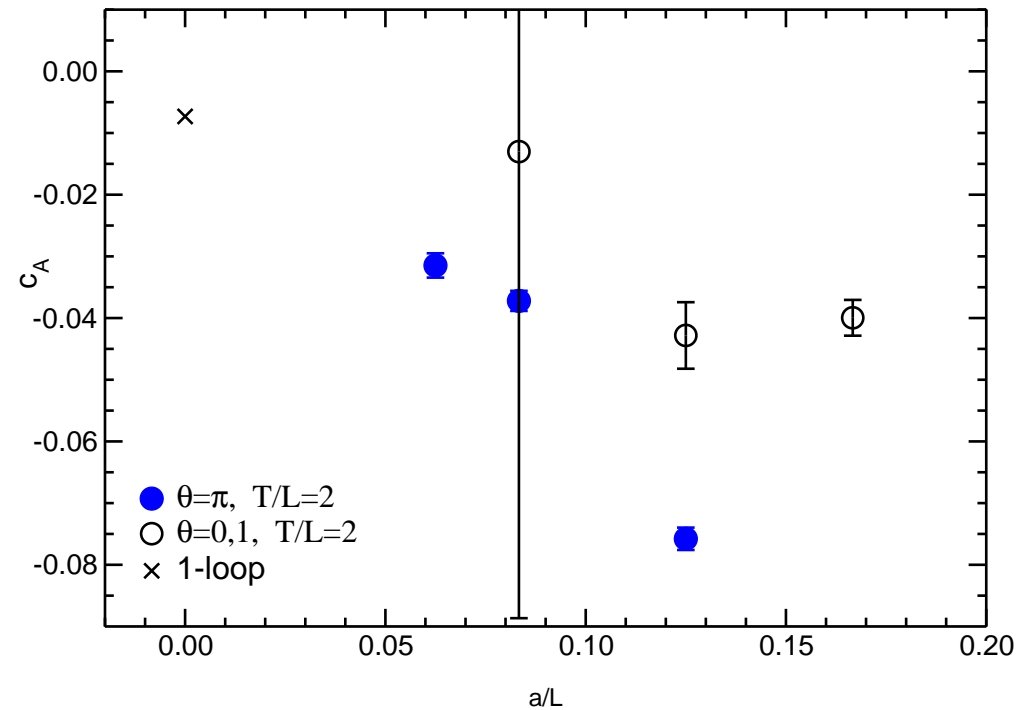
IV-3 a/L dependence of c_A

c_A vs a/L

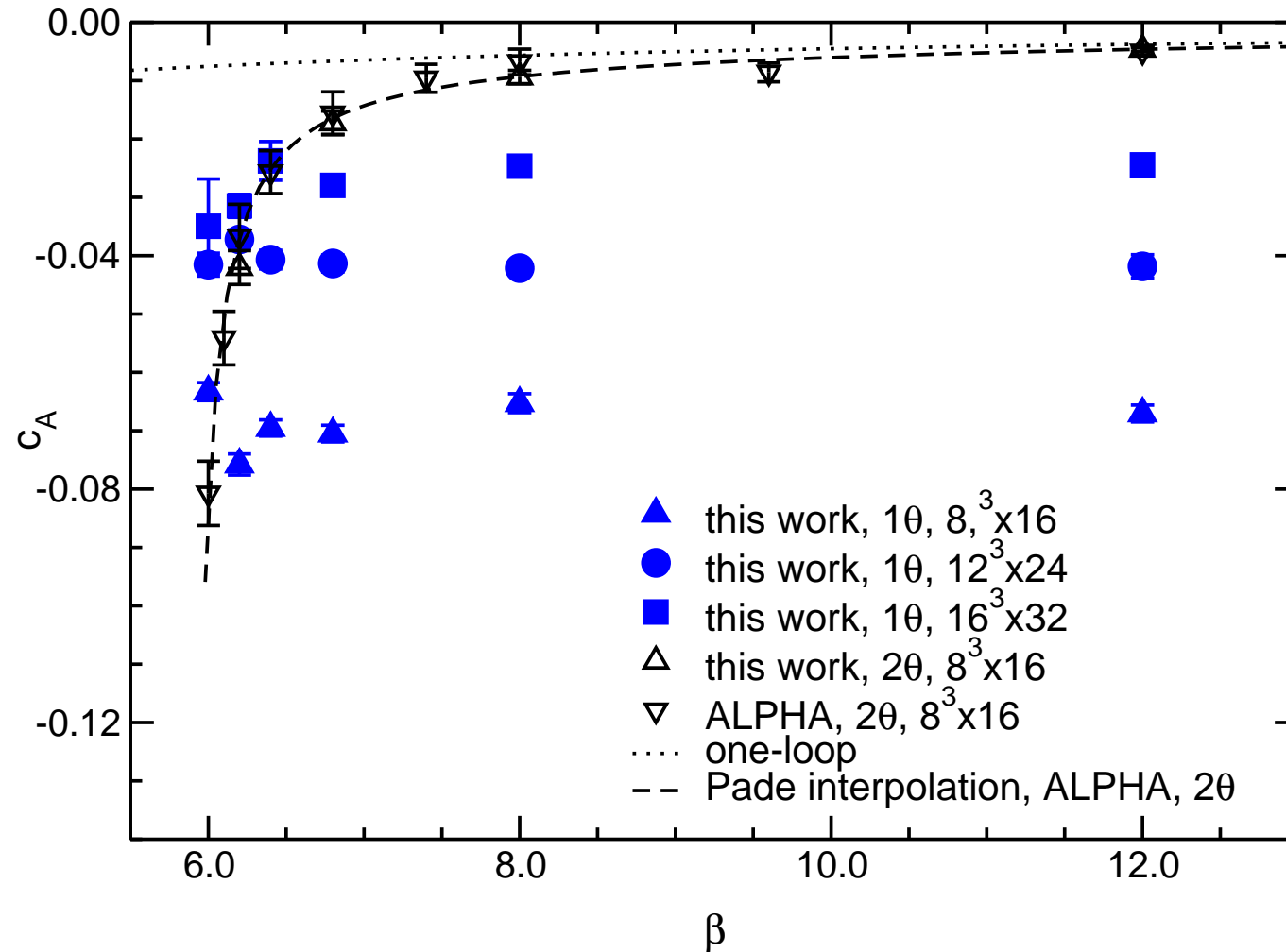
$N_f=0, \beta=12.0, c_{sw}=1.16373, K=0.12991$



$N_f=0, \beta=6.2, c_{sw}=1.61375, K=0.13485$



- c_A has large a/L dependence both at weak and strong couplings

c_A vs β 

- small g^2 dependence of c_A with L/a fixed

perturbation

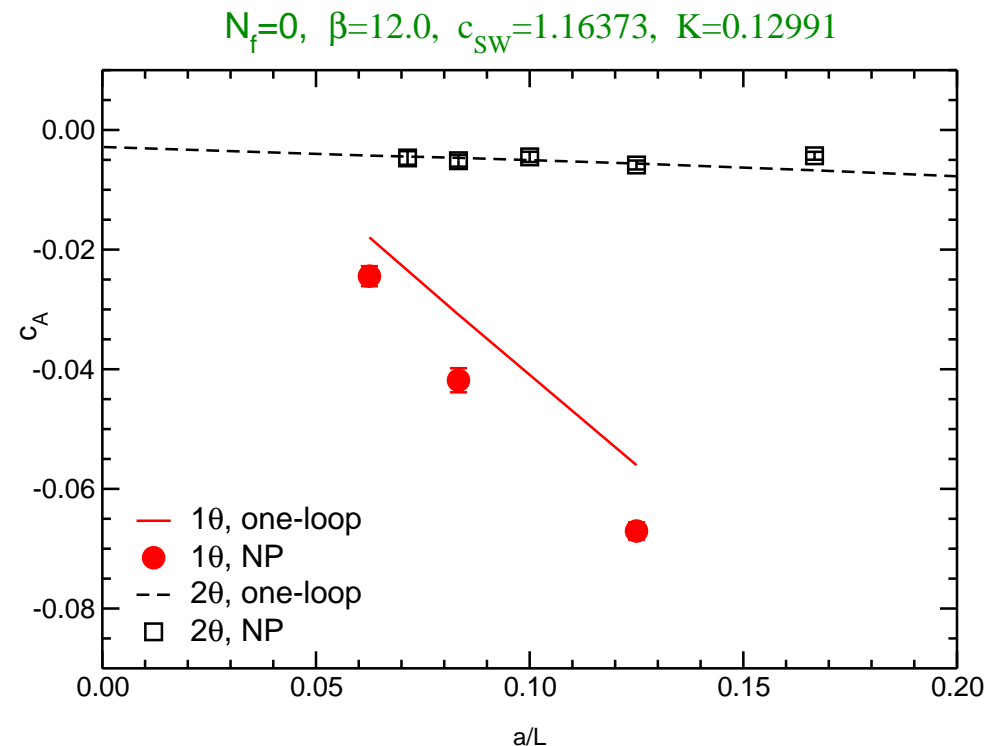
- at tree-level
 - $r(x_0)$, $s(x_0)$ does not depend on x_0
 - $\Rightarrow c_A$ is ill-determined at tree-level

$$c_A = \frac{\Delta m^{(0)} - (r(\theta, x_0) - r(\theta, x_{0,\text{ref}}))}{s(\theta, x_0) - s(\theta, x_{0,\text{ref}})}$$

- at one-loop

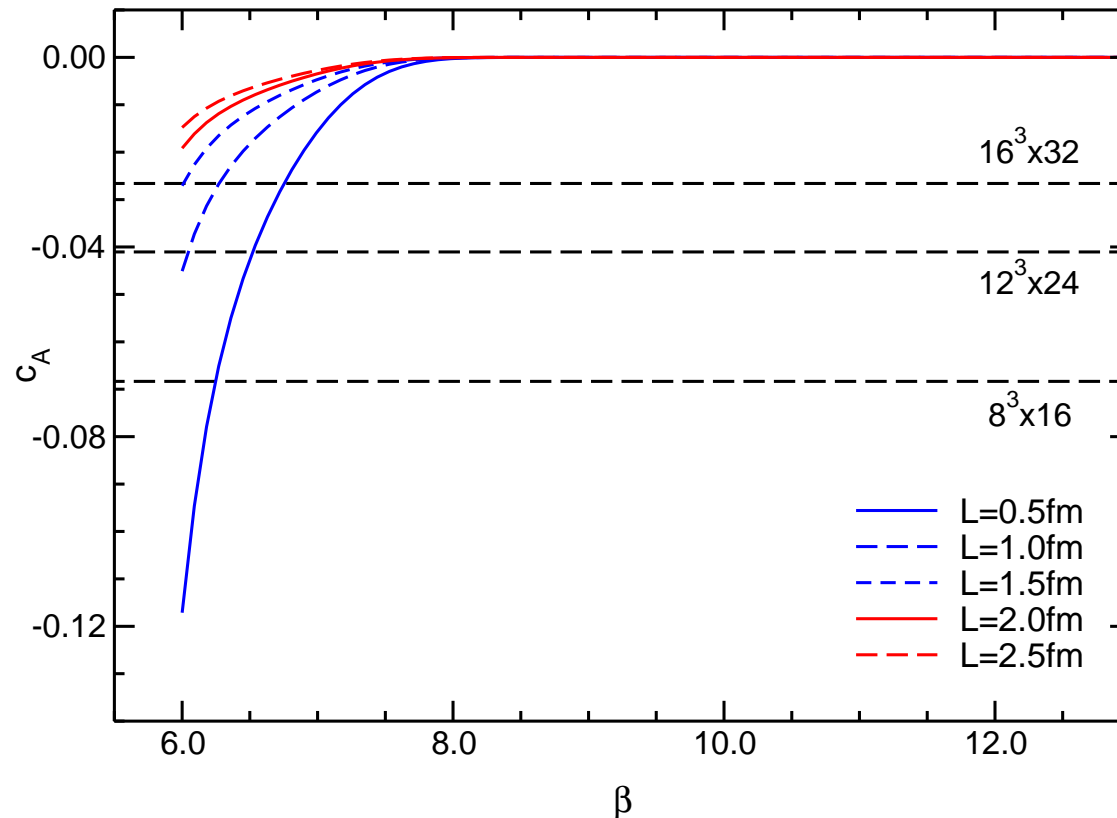
- $c_A \sim O(1)$

\Leftarrow one-loop calc. at $g_0^2 = 1$
and 0.5 give the same c_A
(Aoki, 2003)



on a constant physics line

- assume c_A doesn't depend on g^2
- fit c_A as a function of a/L
- use ALPHA's interpolation formula for r_0 (Guagnelli *et al.*, 1998)



- $c_A \rightarrow 0$ as $g^2 \rightarrow 0$
- strongly depends on the choice of L at $\beta \sim 6$.

$$L = 0.5 - 2.0 \text{ fm}$$

$$\Rightarrow \Delta m_q^{\text{AWI}} \sim 15\%$$

V. Summary

ALPHA/CP-PACS/JLQCD

- collaboration on NP computation of c_A
this talk : a/L dependence in quenched QCD with plaq. gauge

old method

- loses sensitivity to improvement condition
 \Rightarrow error of c_A rapidly blow up as $L \rightarrow \infty$

“slope criterion” method

- good sensitivity
- fixed a/L : independent of g^2 (?), strongly depends on a/L
- fixed L : large volume dependence at low β

“gap criterion” method

- a/L dependence should be studied