

Lattice Perturbation Theory for Heavy Quark Physics

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ILFT Workshop
Shuzenji, Izu, Japan
Sept 23, 2004

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Fermilab, MILC, HPQCD [hep-ph/0408306](https://arxiv.org/abs/hep-ph/0408306)

Outline

- The need for lattice perturbation theory
- The Fermilab action
- Techniques for lattice PT
- Current matching calculations

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- Monte Carlo simulation times scale as $\frac{1}{a^6}$
- Use improved actions to reduce cutoff effects

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- Conversion between α_V and $\alpha_{\overline{MS}}$
- Matching LEFT operators to QCD one

$$\mathcal{O}^{\text{Cont}} = Z_{\mathcal{O}} \left[\mathcal{O}_0^{\text{latt}} + \sum_n c_n \mathcal{O}_n \right]$$

- Short distance coefficients, can match perturbatively if $\alpha_V(\pi/a)$ is not too large.
- Important for determining decay constants, form factors, etc.

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$$\begin{aligned}
 S = & a^4 \sum_x \bar{\psi}(x) \left\{ m_0 + \frac{1}{2} [(1 + \gamma_0) D_0^- - (1 - \gamma_0) D_0^+] \right. \\
 & \left. + \zeta \vec{\gamma} \cdot \vec{D} - \frac{ar_s \zeta}{2} \Delta^{(3)} - \frac{iac_B \zeta}{2} \vec{\Sigma} \cdot \vec{B} - \frac{ac_E \zeta}{2} \vec{\alpha} \cdot \vec{E} \right\} \psi(x)
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- At tree level $\zeta = r_s = c_B = 1$, and $c_E = \frac{1 + \frac{m_0(2+m_0)}{4(1+m_0)}}{1+m_0} \approx 1$

Lattice Perturbation Theory Techniques

- The lattice cutoff means that Feynman rules are much more complicated
- Non-linear connection $U_\mu(x) = e^{igaA_\mu(x)}$ produces an infinite tower of extra vertices
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- Use of highly improved gauge and quark actions adds additional operators
- Automated perturbation theory addresses these issues
- Generate vertex functions (and diagrams) automatically
- Perform loop integrals using VEGAS

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- Correct $\mathcal{O}(a)$ errors by using rotated spinors

$$\Psi(x) = \left[1 + ad_1\vec{\gamma} \cdot \vec{D}\right] \psi(x), \quad d_1 = \frac{m_0}{2(1+m_0)(2+m_0)}$$

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- Matching condition: $\bar{\psi}_c\Gamma_\mu\psi_c - Z_\Gamma\bar{\Psi}_l\Gamma_\mu\Psi_l = 0 + \mathcal{O}(a^2)$

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- $Z_{V_4}^{qq}$ can be determined non-perturbatively
- Expect that the perturbative corrections to ρ will be small, for $m_0^h a$ not too large
- Verified (to one loop) by Harada *et.al.* for Wilson glue and clover light quarks
- Improved glue? Asqtad light quarks?

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- Is $\rho^{(1)} \approx 0$ over the mass range of interest?
- Fermilab approach means $Z_{\Gamma_4} \neq Z_{\Gamma_i}$
- To compute ρ for a given mass pair we need to evaluate six Z factors ($\Gamma = \gamma_4$, $\Gamma = \gamma_i$, $\Gamma = \gamma_5\gamma_4$ and $\Gamma = \gamma_5\gamma_i$)

Computing $Z^{(1)}$

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- Wilson-like fermions have a non-trivial wavefunction renormalization at tree-level

$$Z_{\mathcal{O}}^{(0)} = \sqrt{(1 + m_q)(1 + m_Q)}$$

Computing $Z^{(1)}$ (cont)

- Account for dominant mass dependence by computing $\mathcal{Z} = Z/Z^{(0)}$

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- Important to use the same infrared regulator for both lattice and continuum theories
- Results presented here use a gluon mass
- Could also use twisted periodic boundary conditions
- Lattice to lattice matching is possible for these currents

Calculation Details (continuum)

- Continuum calculation uses a gluon mass IR regulator
- Pauli-Villars regulator Λ for the UV
- Continuum result does not depend on Λ
- Three diagrams

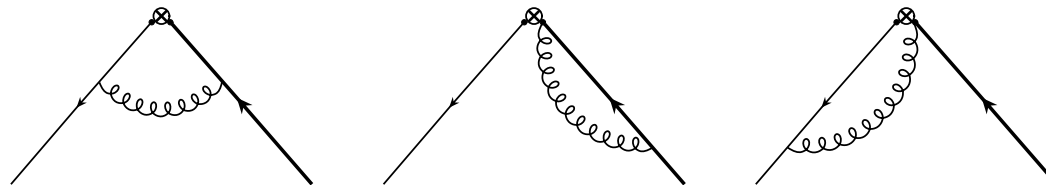
$$\begin{aligned}
 \langle q | \mathcal{O}_c | Q \rangle^{(1)} &= \frac{1}{2} \frac{\partial}{\partial p_0} \text{diagram}_1 \times \text{diagram}_2 + \frac{1}{2} \frac{\partial}{\partial p_0} \text{diagram}_3 \times \text{diagram}_4 \\
 &+ \text{diagram}_5
 \end{aligned}$$

The diagrams are:

- diagram_1 : A fermion line with momentum p and a gluon loop (wavy line) attached to it.
- diagram_2 : A fermion line with momentum p and a Pauli-Villars regulator line (double line with a dot) with momentum \bar{p} .
- diagram_3 : A Pauli-Villars regulator line with momentum \bar{p} and a gluon loop (wavy line) attached to it.
- diagram_4 : A fermion line with momentum p and a Pauli-Villars regulator line with momentum \bar{p} .
- diagram_5 : A fermion line with momentum p and a Pauli-Villars regulator line with momentum \bar{p} , with a gluon loop (wavy line) attached to the Pauli-Villars line.

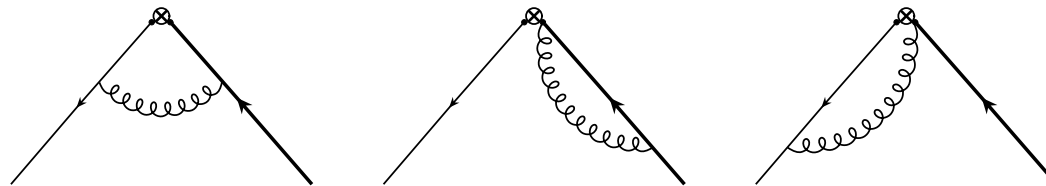
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- Lattice calculation uses a gluon mass IR regulator
- Five diagrams without the rotations
- For clover light, and Fermilab heavy, the rotations add 12 further diagrams
- For Asqtad light the rotations add 3 diagrams



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- evaluate diagrams by brute force VEGAS integration

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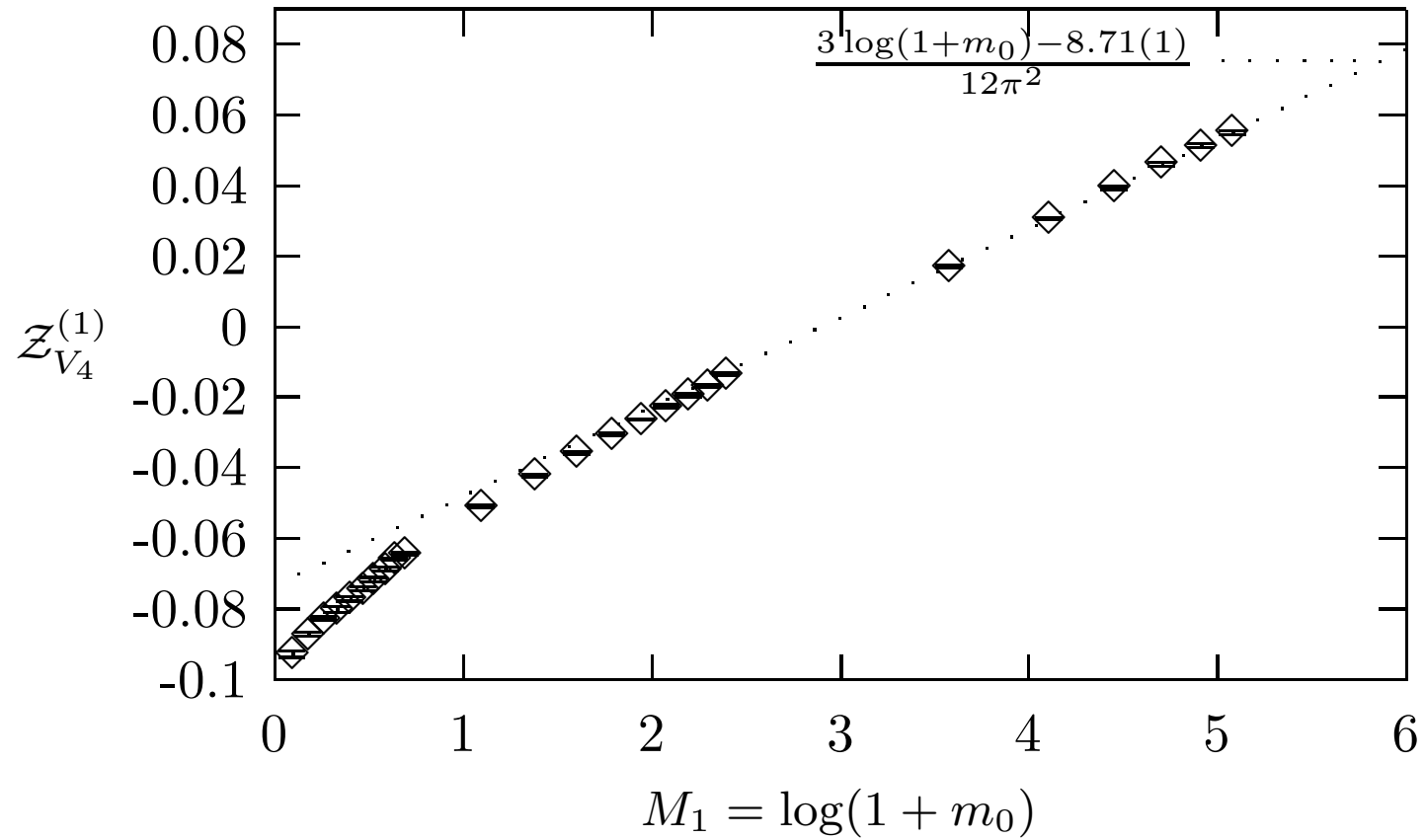
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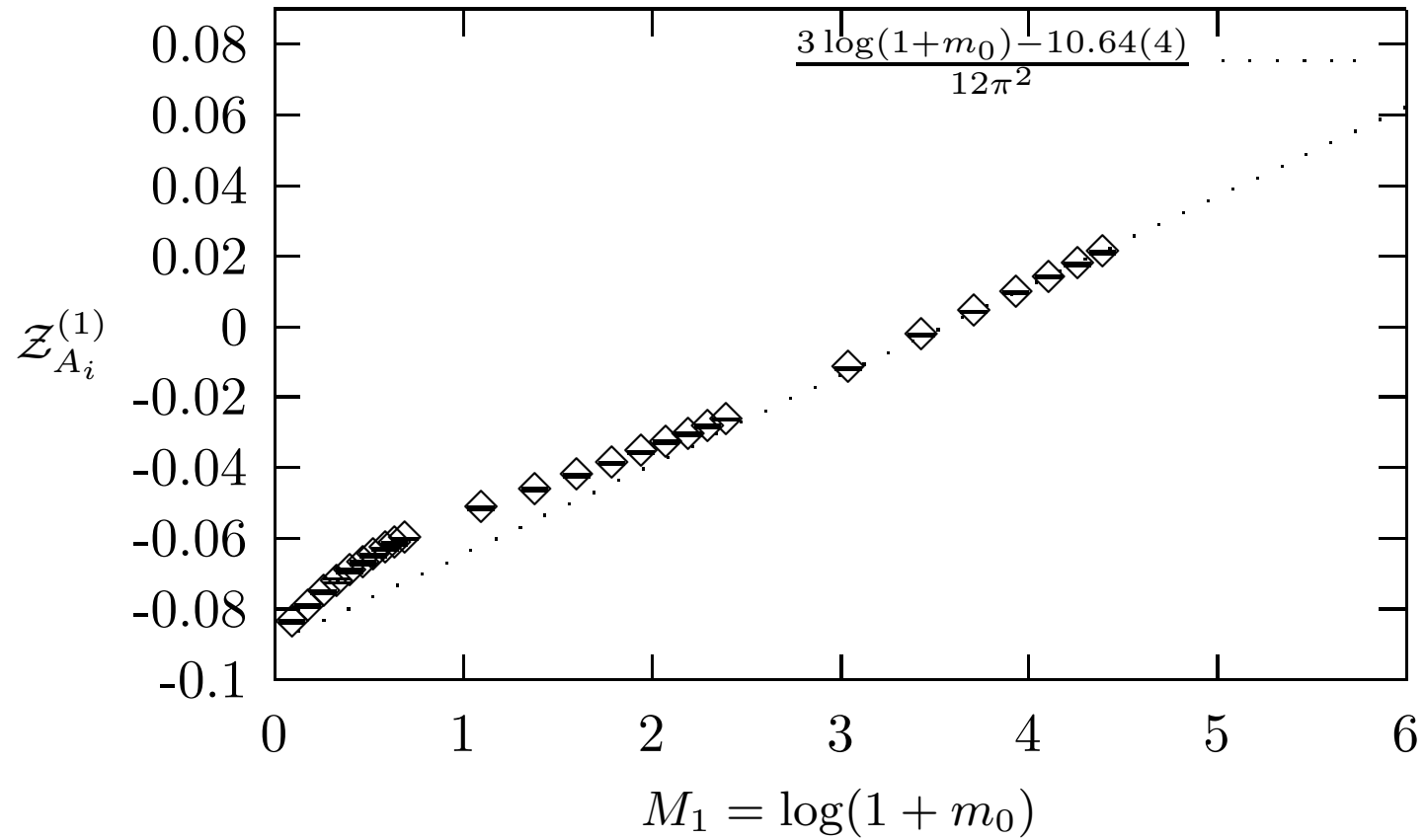
$$Z_{\Gamma} = \frac{z_{\Gamma} + 3 \log(1 + m_0)}{12\pi^2}$$

3. Agreement between two independent calculations for the self energy Z_2 with improved glue and asqtad quarks
 - All these checks are satisfied

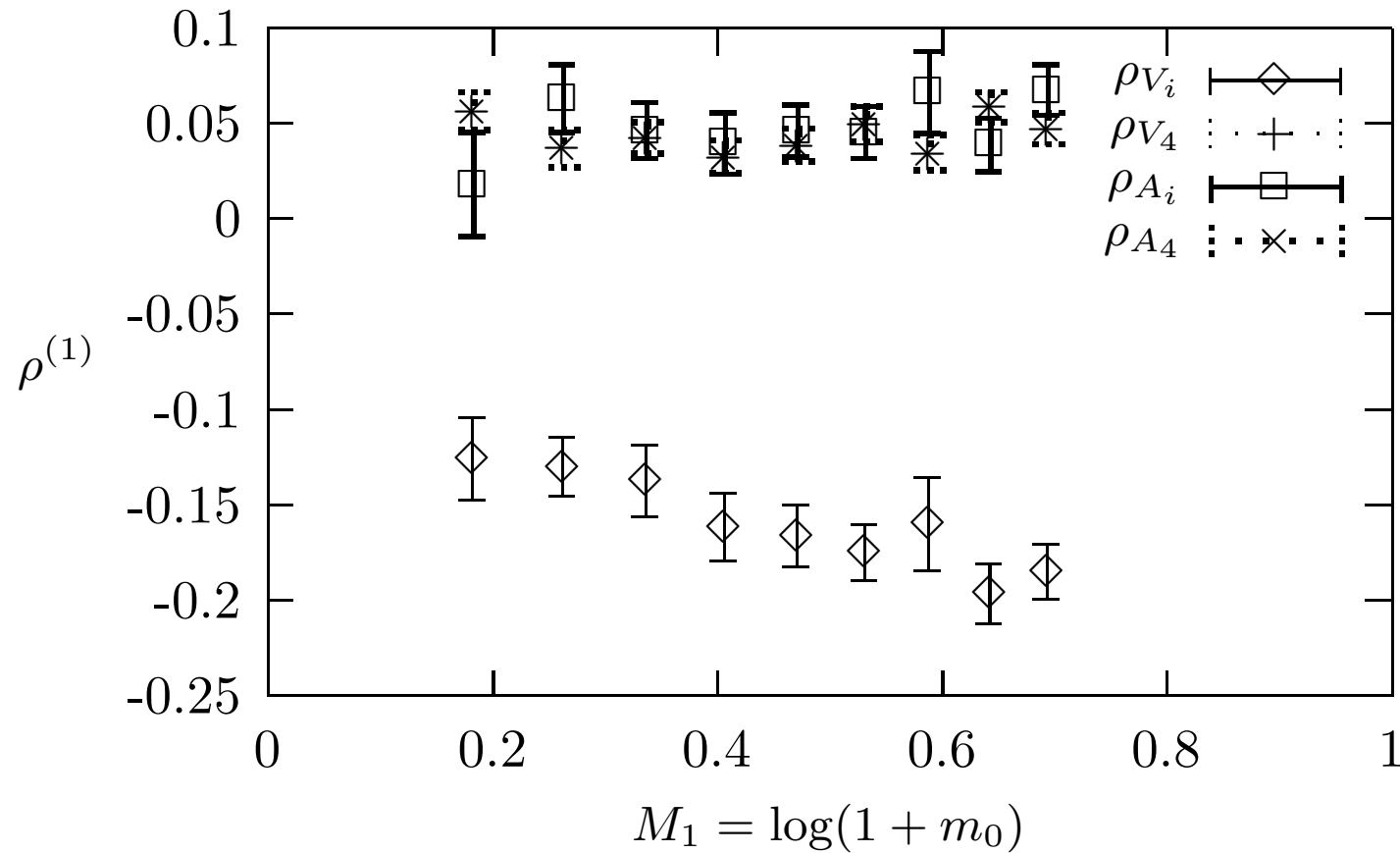
$Z_{V_4}^{hl}$ with Improved glue



$Z_{A_i}^{hl}$ with Improved glue

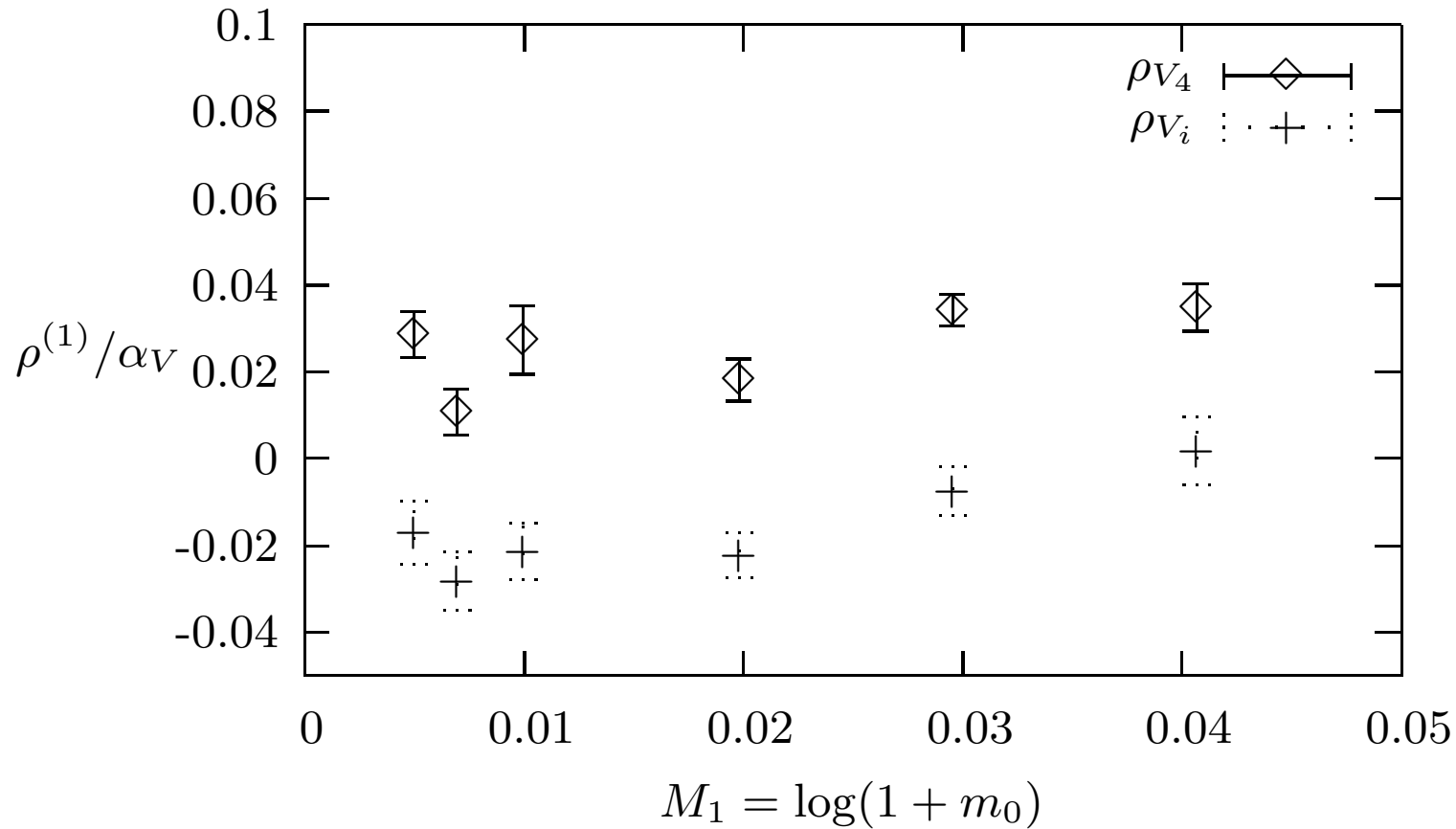


ρ factors for light clover quarks and improved glue



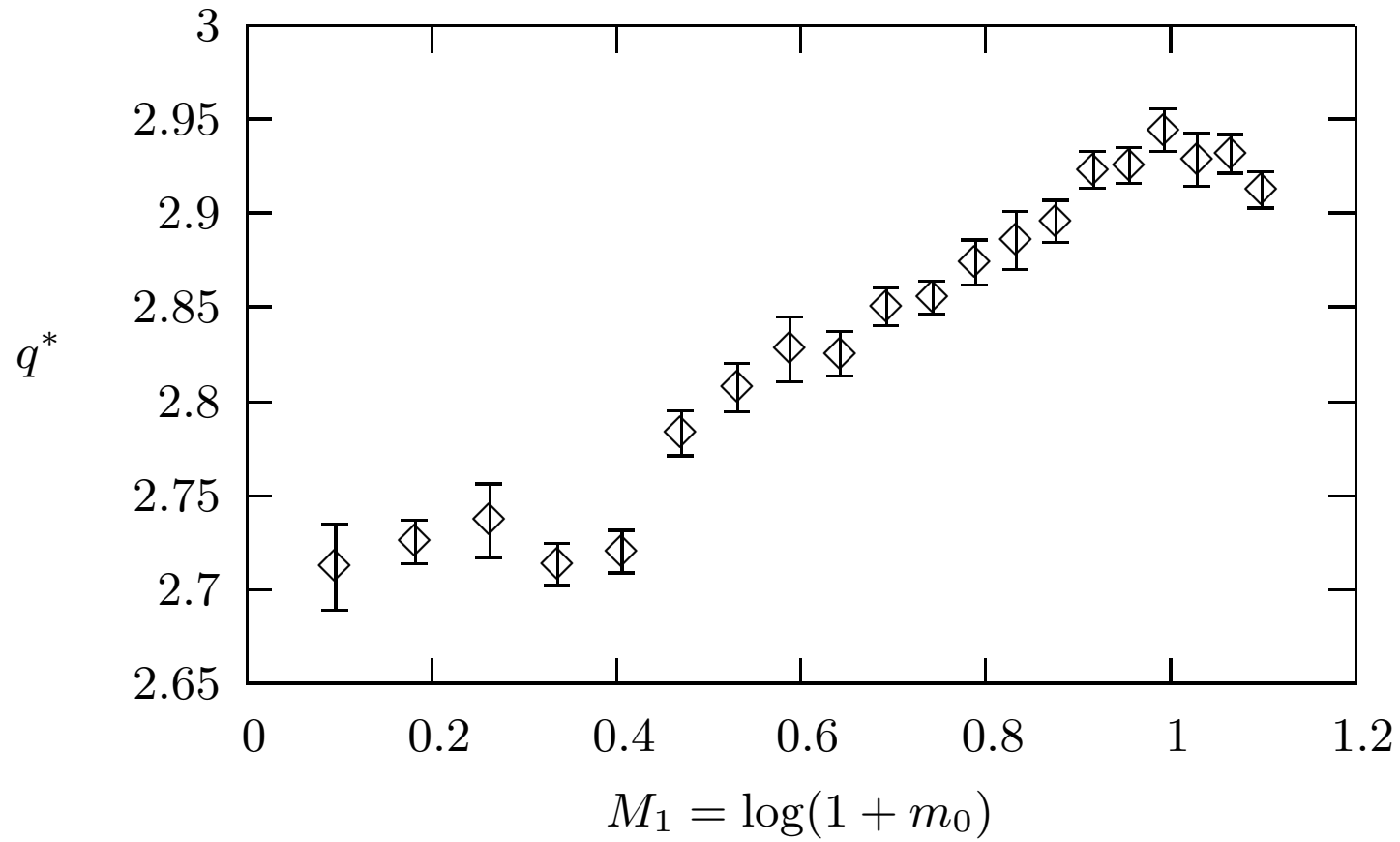
$m_{0,\ell} = 0.01$

ρ factors for light Asqtad quarks and improved glue



$$m_{0,c} = 0.575$$

q^* for Z_{V_4}



$$m_{0,\ell} = 0.001$$

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$$\rho_{V_\parallel} = 1.01 + \mathcal{O}(\alpha_V^2)$$

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Use of ρ factors greatly reduces perturbative errors due to current matching

Conclusions and Outlook

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- Automated perturbation theory works well on current matching calculations
- ρ factors were expected to have small perturbative corrections
- The one loop calculations presented here confirm this
- Smallness of one loop terms means that two loop PT is not needed
- Does a similar factorization work for NRQCD?
- Dominant error is now matching of parameters in the action
- These calculations are nearing completion (Fermilab and NRQCD)