



Charm Physics

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Flavour physics

- Heavy flavour sector
 - Charm **not so well served** by HQET
 - CLEO-c and B-Factories → lots of data
 - f_{D_s} critical quantity
- Charmed hadrons accessible on lattice
 - $am_c < 1$
 - Simulate **directly** on lattice
 - Easy to control **effective theory techniques**



Outline

- Past calculations - Quenched
 - Ds spectrum – New states
 - A. Dougall, R.D. Kenway, C.M., C. McNeile Phys.Lett. B569 (2003) 41-44
 - Charm quark mass
 - A. Dougall, C.M., C. McNeile hep-lat/0409089
- Future calculation – dynamical fermions
 - f_{D_s} and Ds spectrum: Overlap on KS sea
 - Heavy quarks on a GW sea



New states in D_s spectrum

- B. Aubert *et al.* (BaBar) Phys.Rev.Lett 90 242001, 2003 Cited 162 times!
- Discovery of new narrow state
 - Mass ~ 2.32 GeV
 - Decays to $D_s^+ \pi^0$ $J^P=0^+(?)$
 - Is it $c\bar{s}$ meson?
- Confirmed by Belle and CLEO
 - Other narrow states discovered



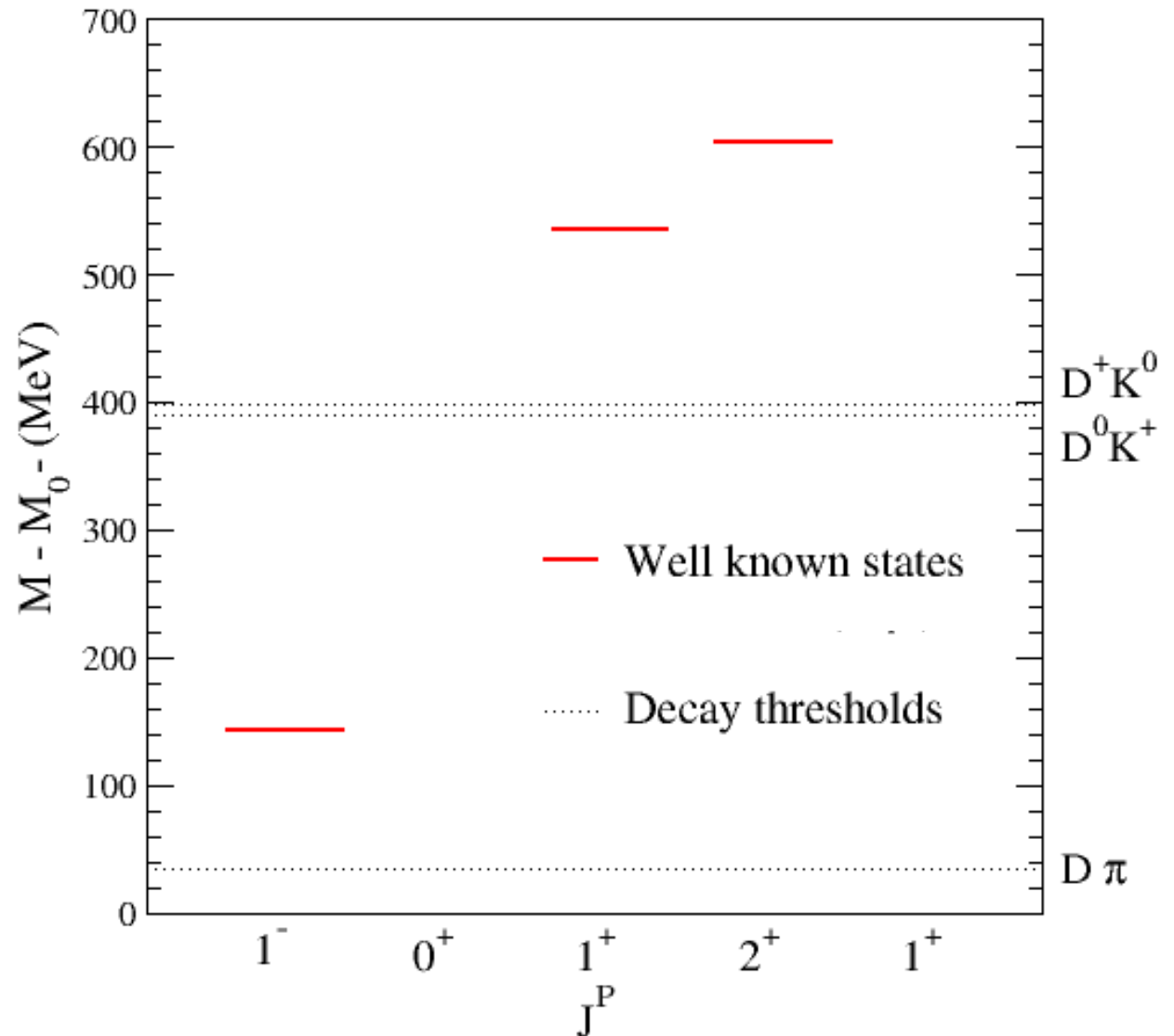
Ds Spectrum

Potential model
mass is high,
resonance is wide

BaBar: Narrow
resonance in $D+\pi$
channel $J^P=0^+$

speculation that
new state is not
a meson but 4
quark molecule

$$0 \rightarrow M_{D_s}(Ps) = 1.9685 \text{ GeV}$$





Problems for theory?

- These states “*not predicted by theory*”
- Some **melodramatic** statements in popular press
- “*This sends theorists back to their drawing boards*”
 - SLAC director Jonathan Dorfan – **BBC**
- Predictions come from **Quark Model** not **QCD**



Correlation functions

Local correlation functions

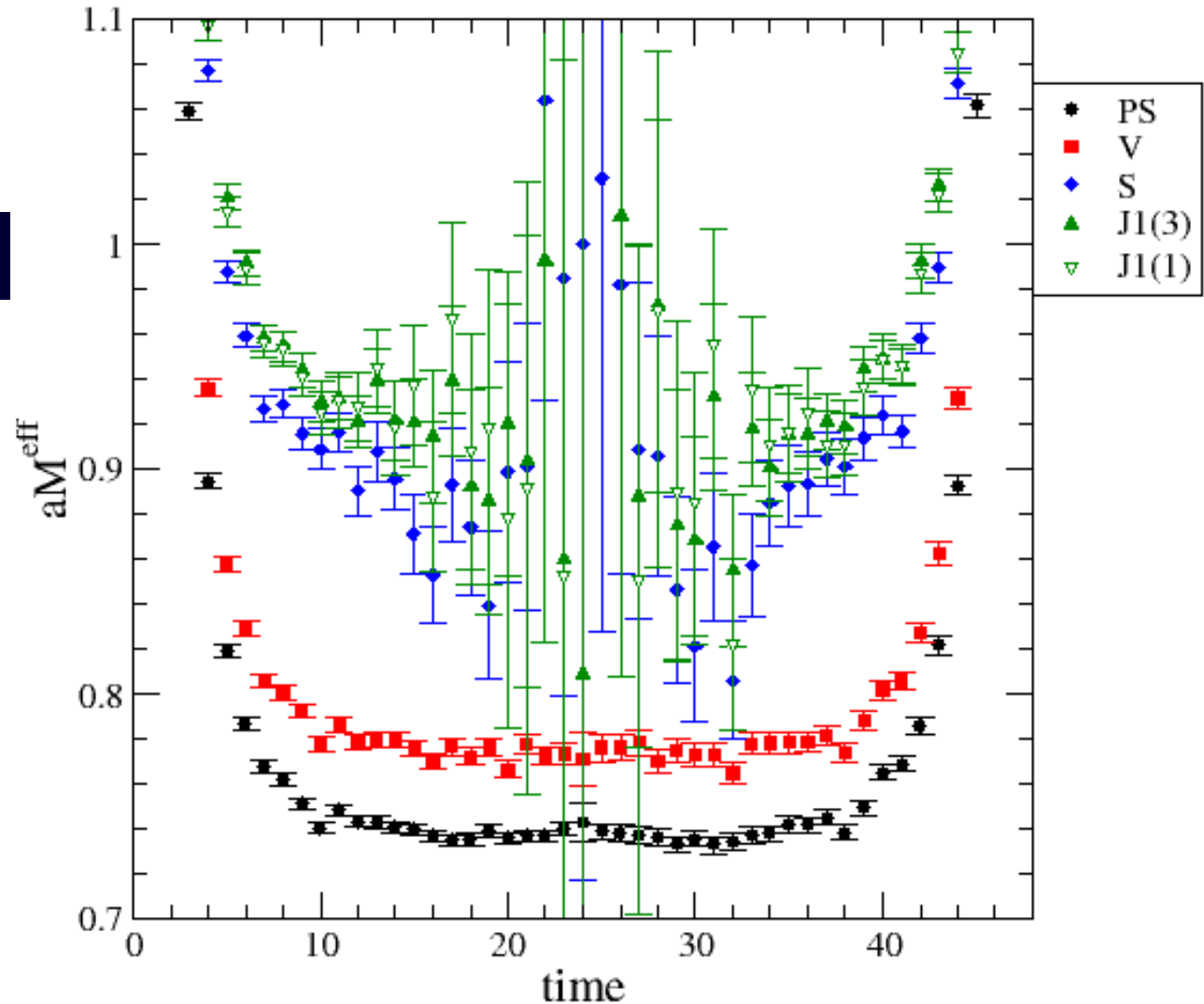
$$\Omega_{ij}(x) = \psi_i(x) \Gamma \psi_j(x)$$

$$\Gamma = \{1, \gamma_\mu, \gamma_5, i\gamma_5 \gamma_\mu, \sigma_{\mu\nu}\}$$

Quenched $\beta=6.2$
Valence quark masses roughly

$c\bar{s}$

$$C(\tilde{p}, t) = \sum_{\tilde{x}} e^{i\tilde{p} \cdot \tilde{x}} \langle \Omega_X(\tilde{x}, t) \Omega_X(\tilde{0}, 0) \rangle = \sum_n A e^{-m_n t}$$





Data sets

*Wilson Glue, Clover-Wilson quarks NP c_{SW}
Quenched and $N_F=2$*

$\{ \beta , \kappa \}$	V	a^{-1} (GeV)	#cfgs	M_{PS}/M_V
6.2,0	$24^3 \times 48$	2.91	216	N/A
6.0,0	$16^3 \times 48$	2.12	496	N/A
5.93,0	$16^3 \times 32$	1.90	623	N/A
5.2,0.1350	$16^3 \times 32$	1.91	202	0.70(1)

Scale set by $r_0 = 0.5 \text{ fm}$



P wave states and symmetry

- States labelled by **spin** of **light quark** j

$$L = 1 \Rightarrow j = \left\{ \frac{1}{2}, \frac{3}{2} \right\}$$

- Plus **spin** of **heavy quark** \rightarrow 2 doublets

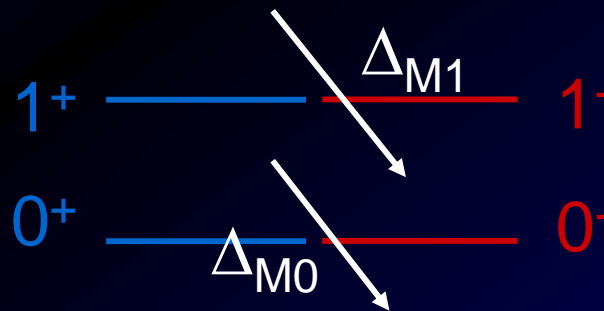
$$j = \frac{1}{2} \ni J = \{0, 1\} \quad j = \frac{3}{2} \ni J = \{1, 2\}$$

- $J=1$ states no **definite C**
 - Mixing \rightarrow see **lowest state** on lattice



Heavy and Chiral symmetry

- Double limit \rightarrow parity doublets $\{0^-, 1^-\}$
 $\{0^+, 1^+\}$ degenerate
- Spontaneous χ -sym breaking splits
 - $\Delta_{M1}(1^+-1^-) = \Delta_{M0}(0^+-0^-)$



- CLEO $\Delta_{M1}=351(2)$, $\Delta_{M0}=350(1)$ MeV



Mass splittings

$$\frac{C_i(t)}{C_j(t)} = \lim_{t \rightarrow \infty} A_{ije}^{- (m_i - m_j)t}$$

- Fitting Δm can help reduce stat errors
- Interpolate mass splittings
 - $\Delta_{M0}=0^+-0^-$, $\Delta_{M1}=1^+-1^-$, $\Delta_{MH}=1^- - 0^-$
 - $(m_{PS})^2$ for lights – kaon for strange
 - $1/m_{PS}$ for heavies – M_{Ds} for charm



Continuum extrapolation

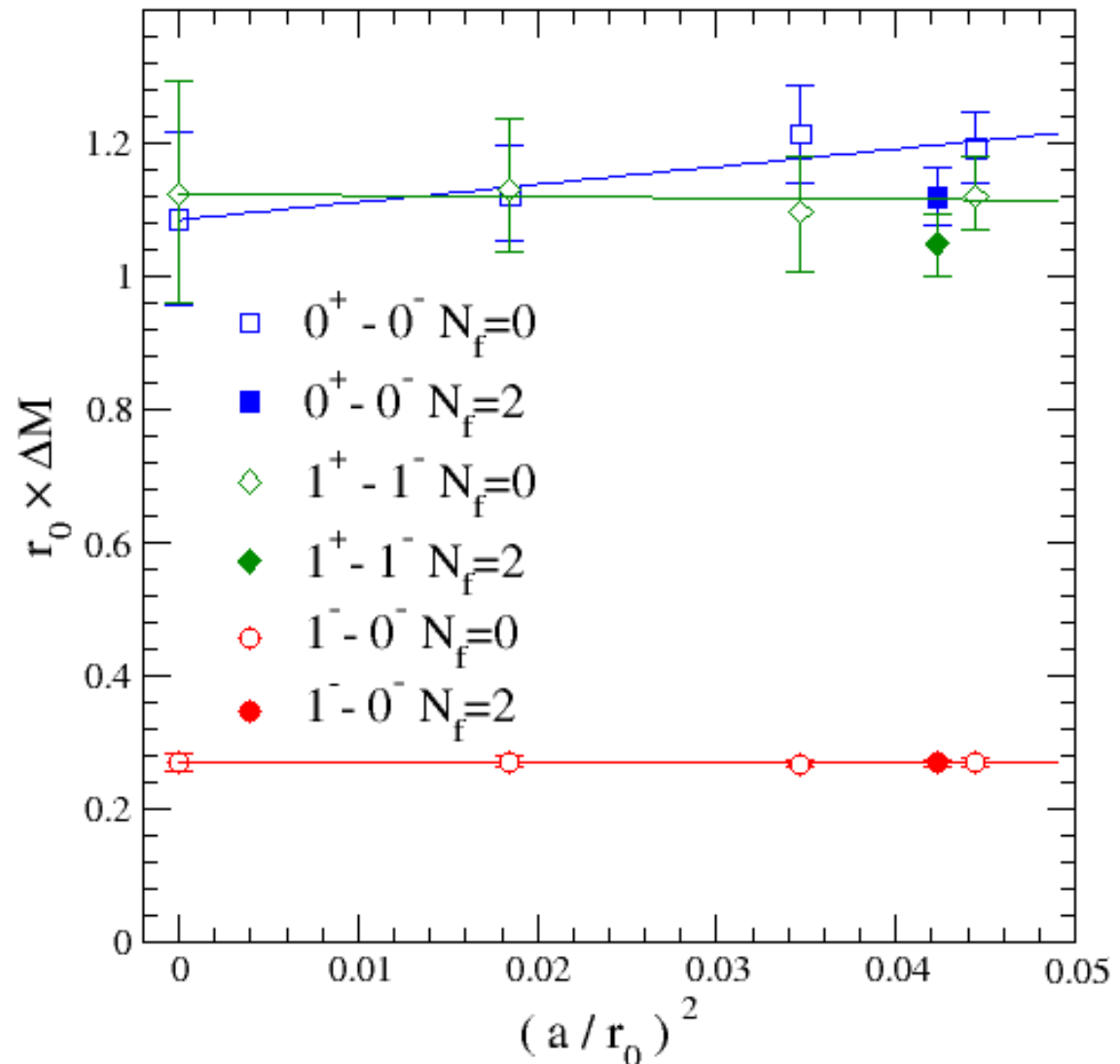
Green and blue
Symbols show

Δ_{M1} and Δ_{M0}

Red is Δ_{MH}

$$\Delta_{M0} = \Delta_{M1}$$

No statistically
significant
effect for DF

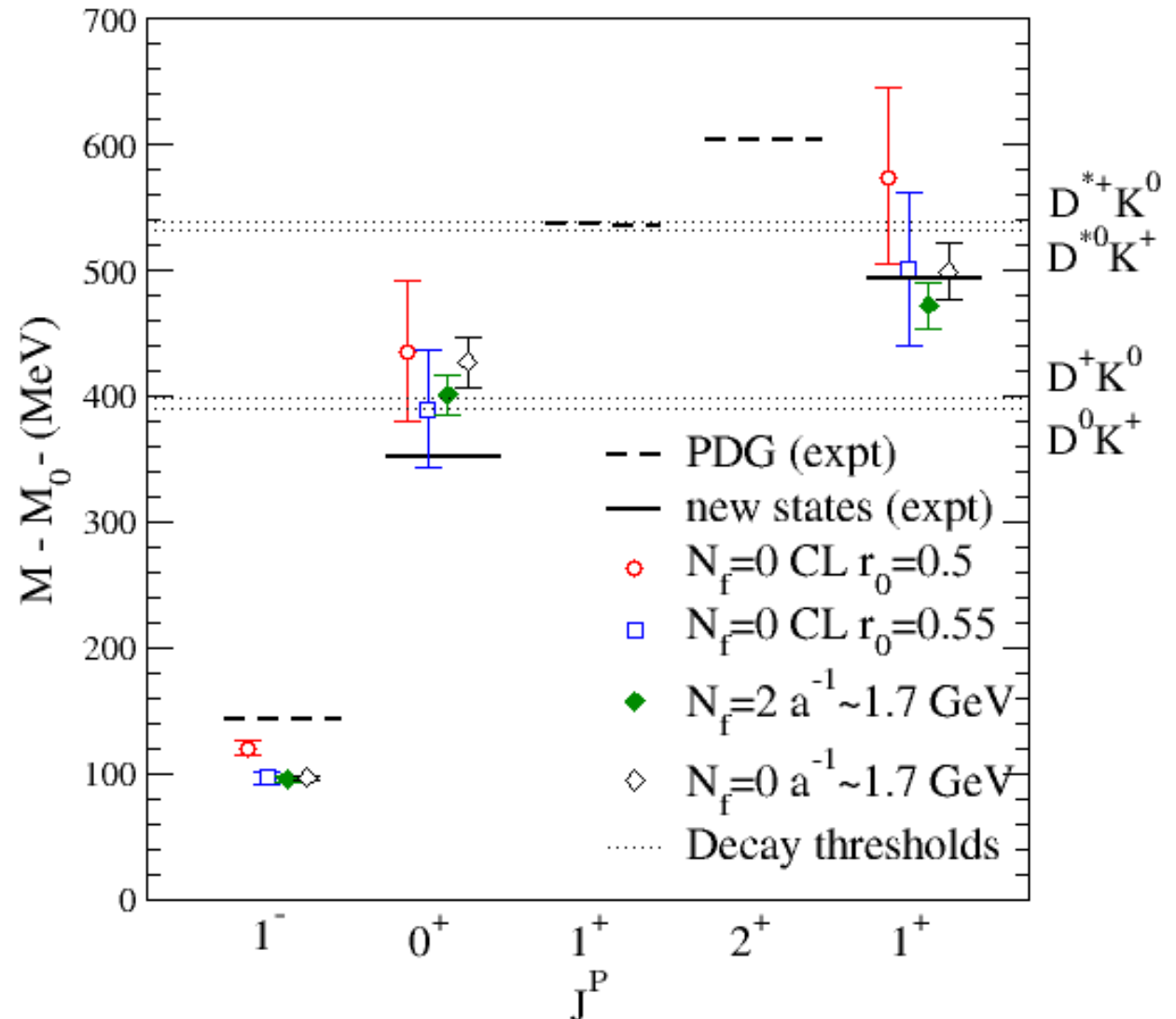




Ds Spectrum II

Results consistent with mesonic interpretation

Narrow width as state is below threshold and decay to pion channel is isospin suppressed





Charm quark mass

- $M_c \sim M(\eta_c)/2 - \Lambda_{\text{QCD}} \sim 1.3 \text{ GeV}$
- M_c is both scale and scheme dependent
 - Not well known quantity

$$1.0 < \overline{m}_{\text{charm}}^{\overline{MS}}(\overline{m}_{\text{charm}}) < 1.4 \text{ GeV} \quad \text{PDG2002}$$

- Many quenched calculations
 - Require dynamical fermions to do better



General strategy

- Coefficients $\{Z_J, b_J, c_J\}$ not known NP on coarse lattice
 - Use tadpole improved PT Bhattacharya *et al* hep-lat/0009038 in MS scheme
 - Lattice artefacts $O(\alpha_s a)$ and $O(a^2)$
 - coarsest lattice $\rightarrow am_q \sim 0.7$
 - Compare to FNAL method
- Interpolate quark masses to charm (Ds)
- Scale $m_c(\mu=a)$ to $m_c(\mu=m_c)$ with RunDec



Axial Ward Identity

$$m_{ij} = \frac{\partial_4 \langle A_4^I \mathcal{O} \rangle}{2 \langle P \mathcal{O} \rangle} = \frac{m_i + m_j}{2}$$

$$A_4^I = A_4 + ac_A \partial_4 P$$

$$A_4(x) = \bar{\psi}_i(x) \gamma_4 \gamma_5 \psi_j(x)$$

$$P(x) = \bar{\psi}_i(x) \gamma_5 \psi_j(x)$$

$$m_A = \frac{Z_A}{Z_P} (1 + [b_A - b_P] am) m$$



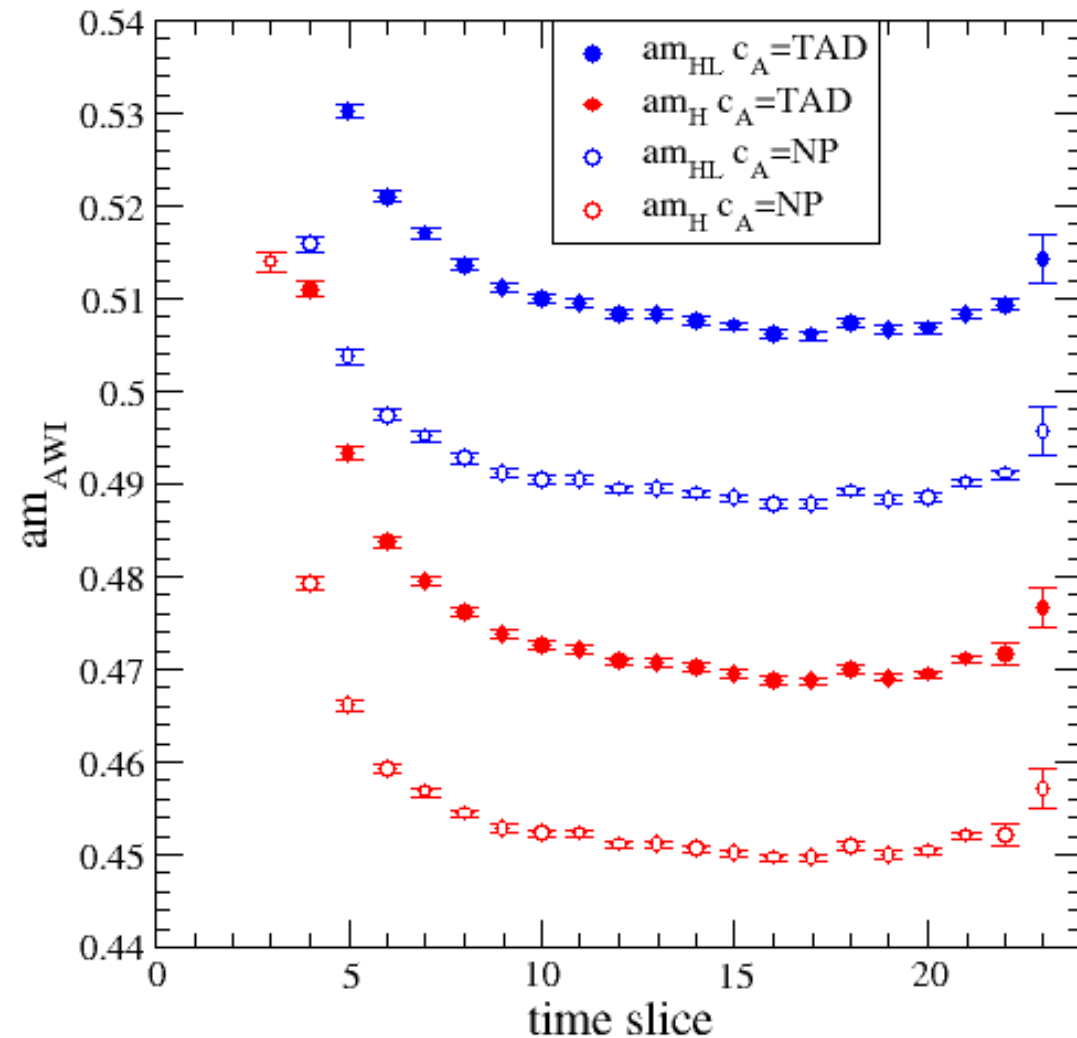
Bare axial mass

Effective mass
plot for axial Ward
identity mass

Open symbols
show the effect of
NP C_A

Effect of NP C_A
half as big as
 am_L

$$\beta=6.2 \quad \kappa_H=0.1200 \quad \kappa_L=0.1346$$





Hopping parameter mass

$$m_q = \frac{1}{2} \left(\frac{1}{\kappa} - \frac{1}{\kappa_{\text{crit}}} \right)$$

$$m_V = Z_M (1 + b_m a m_q) m_q$$

$$b_m = -1/2 \text{ at tree-level}$$



FNAL Method

Aim to include **all** $O(am)^n$ lattice artefacts

Tree level definitions ...

$$m_2 = \frac{e^{m_1} \sinh m_1}{1 + \sinh(m_1)}$$

$$m_1 = \ln(1 + m_q) \sim m_q - \frac{1}{2}m_q^2 + \dots \quad (2)$$

m_1 same as m_V to $O(am)$

Use one-loop definitions in Mertens *et al*
hep-lat/9712024

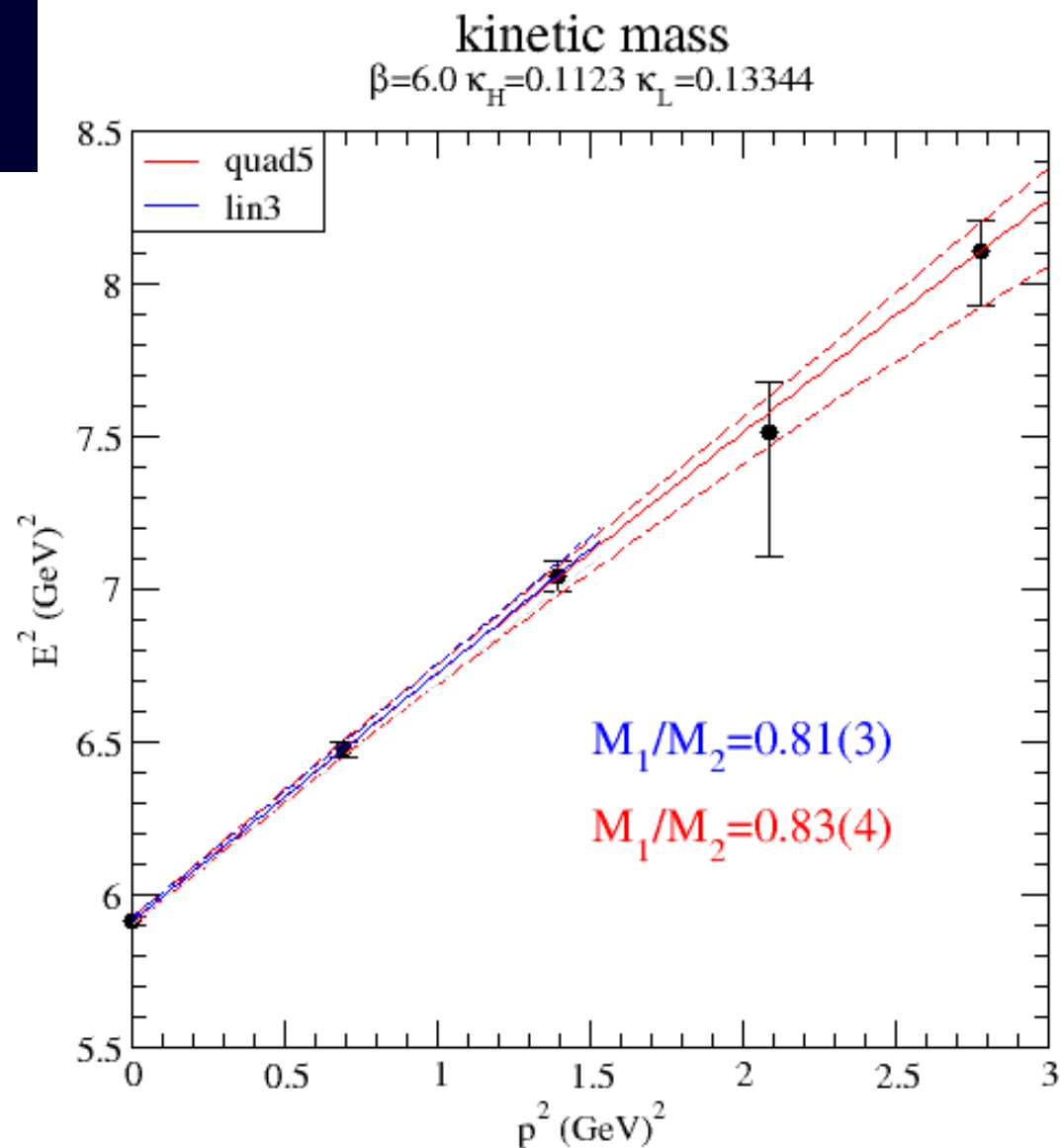


Dispersion relation

$$E^2 = M_1^2 + \frac{M_1}{M_2} \vec{p}^2 + \mathcal{O}(\vec{p}^4)$$

$$\frac{1}{M_2} = \frac{\partial^2 E}{\partial p_k^2}$$

2004/9/30





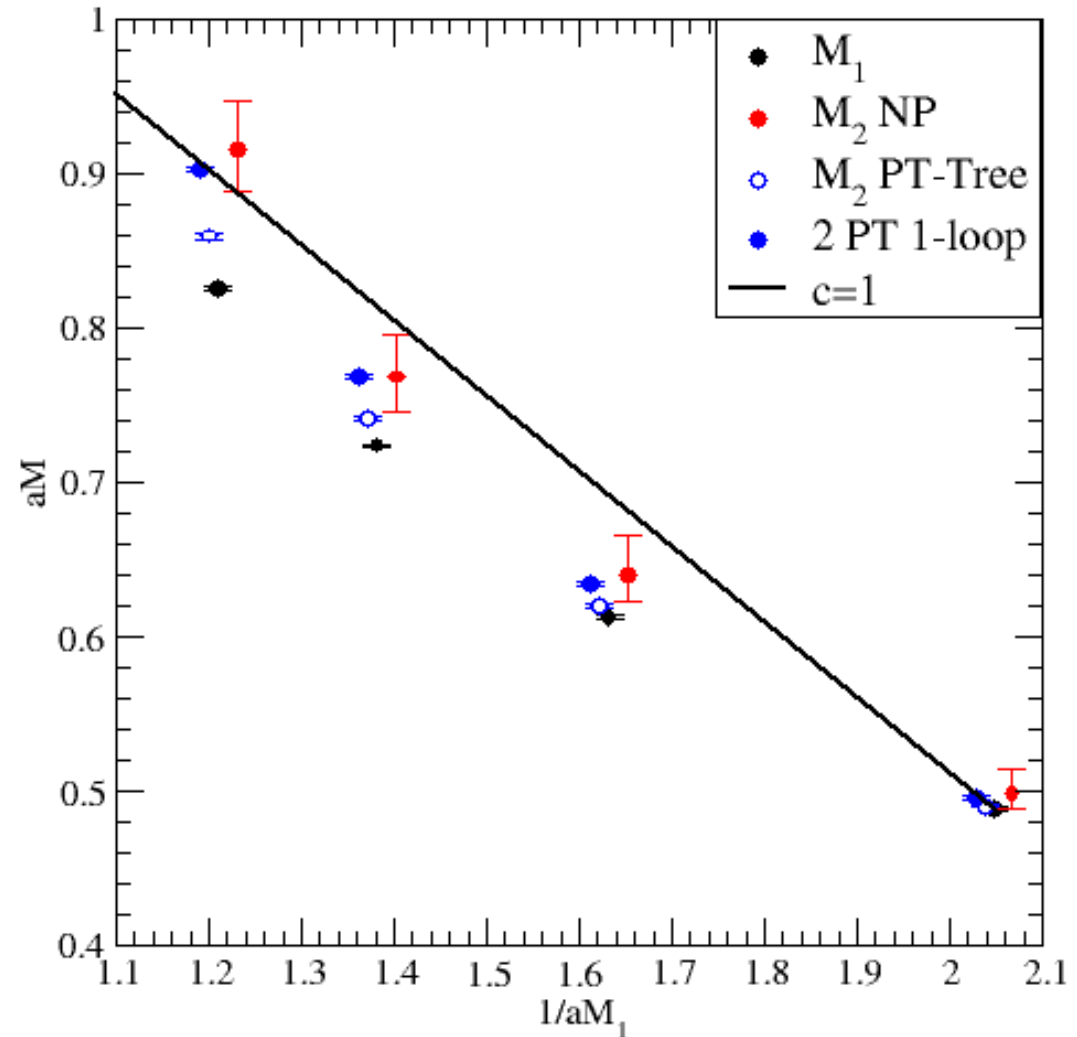
M_1 vs M_2 Hadron masses

$$M_2 = M_1 + (m_2 - m_1)$$

Tree-level M_2 is a poor estimate of NP M_2

1-loop M_2 track NP M_2

Use 1-loop M_2





Functional form of interpolation

HQS

$$f(x) = b_0 + \frac{b_1}{x} + \dots$$

But mass have $O(am)$ corrections

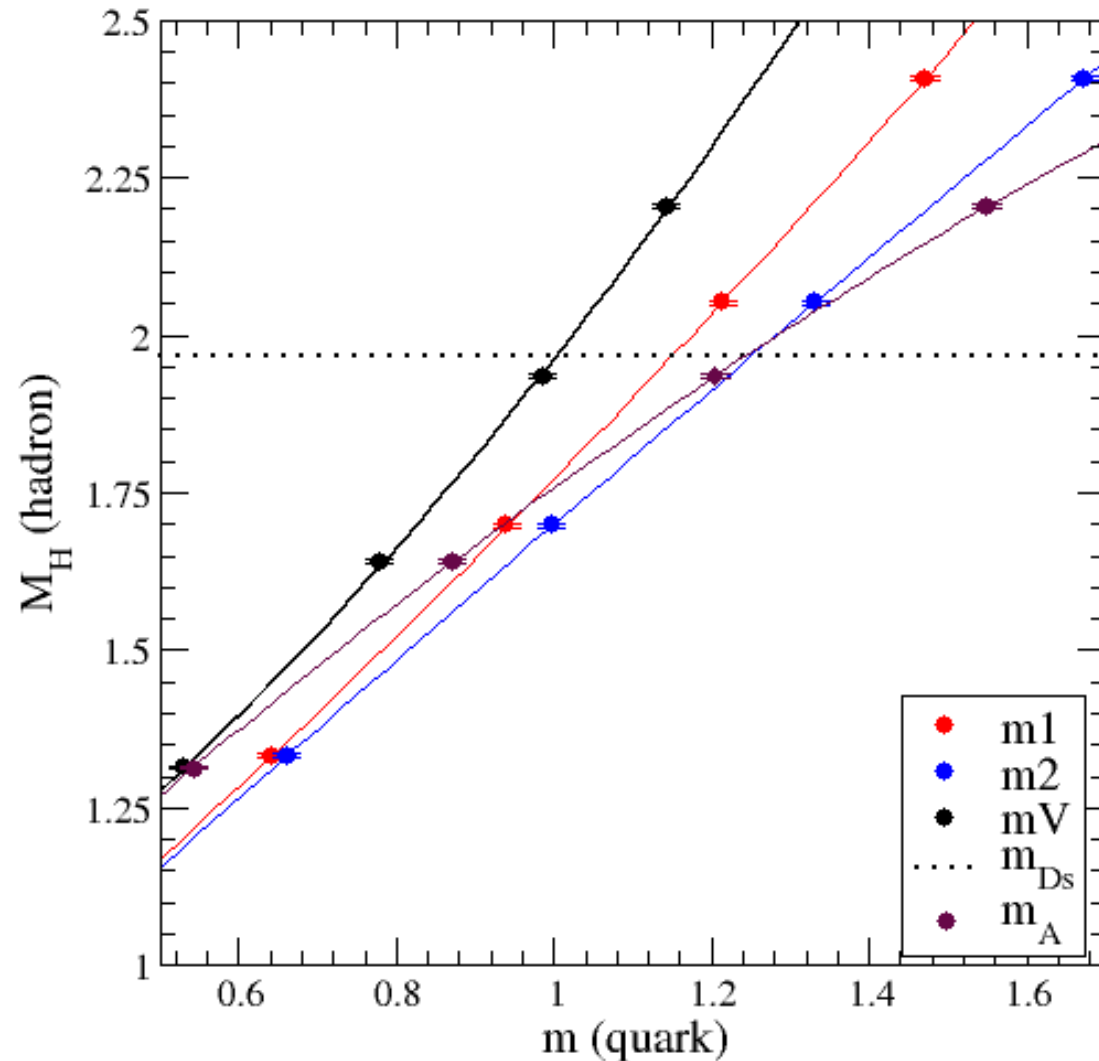
$$f(x) = a_0 + a_1x + \dots$$

Same in fixed range

$$\frac{1}{4} < x < \frac{3}{4}$$

Fixing the charm quark mass

$\beta=6.2$





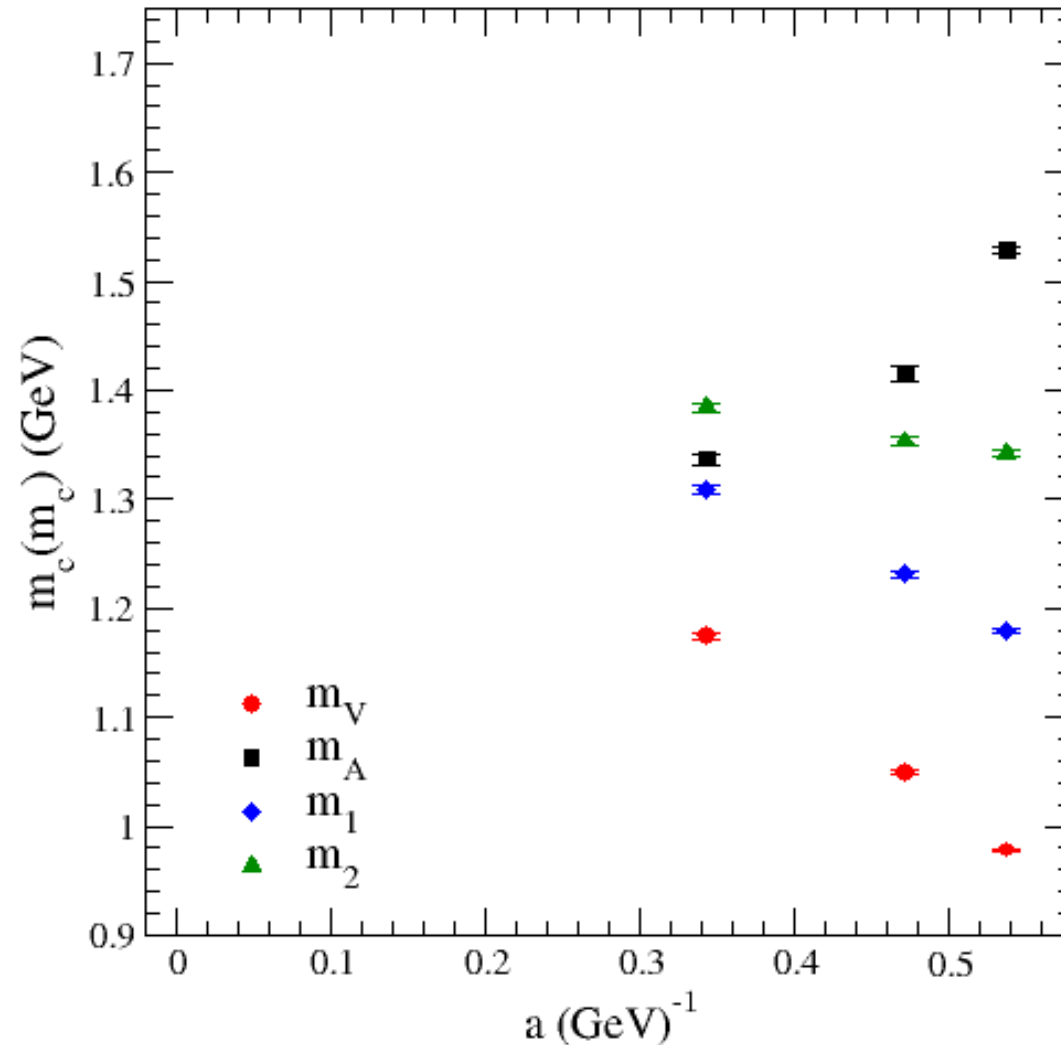
Quenched data

Lattice artefacts

$O(a)$ and $O(\alpha_s a)$

But $O(a^2)$ from
data

Not linear!

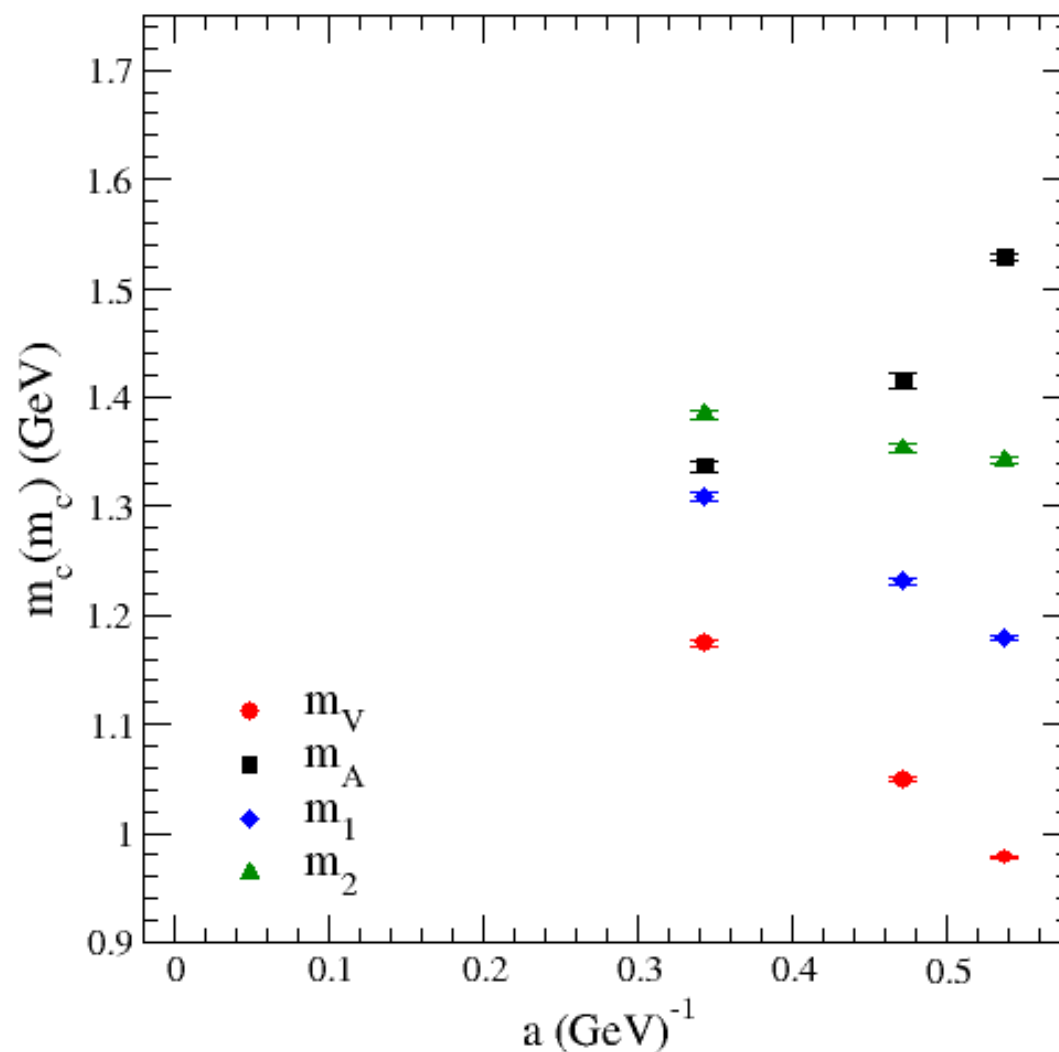


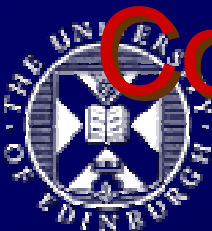


Continuum limit of “ALPHA” masses

Consistent defⁿ
of quark masses
should give
same continuum
limit (Rolf and Sint
JHEP 0212 (2002) 007)

Simultaneous fit
to m_A and m_V
(quadratic)

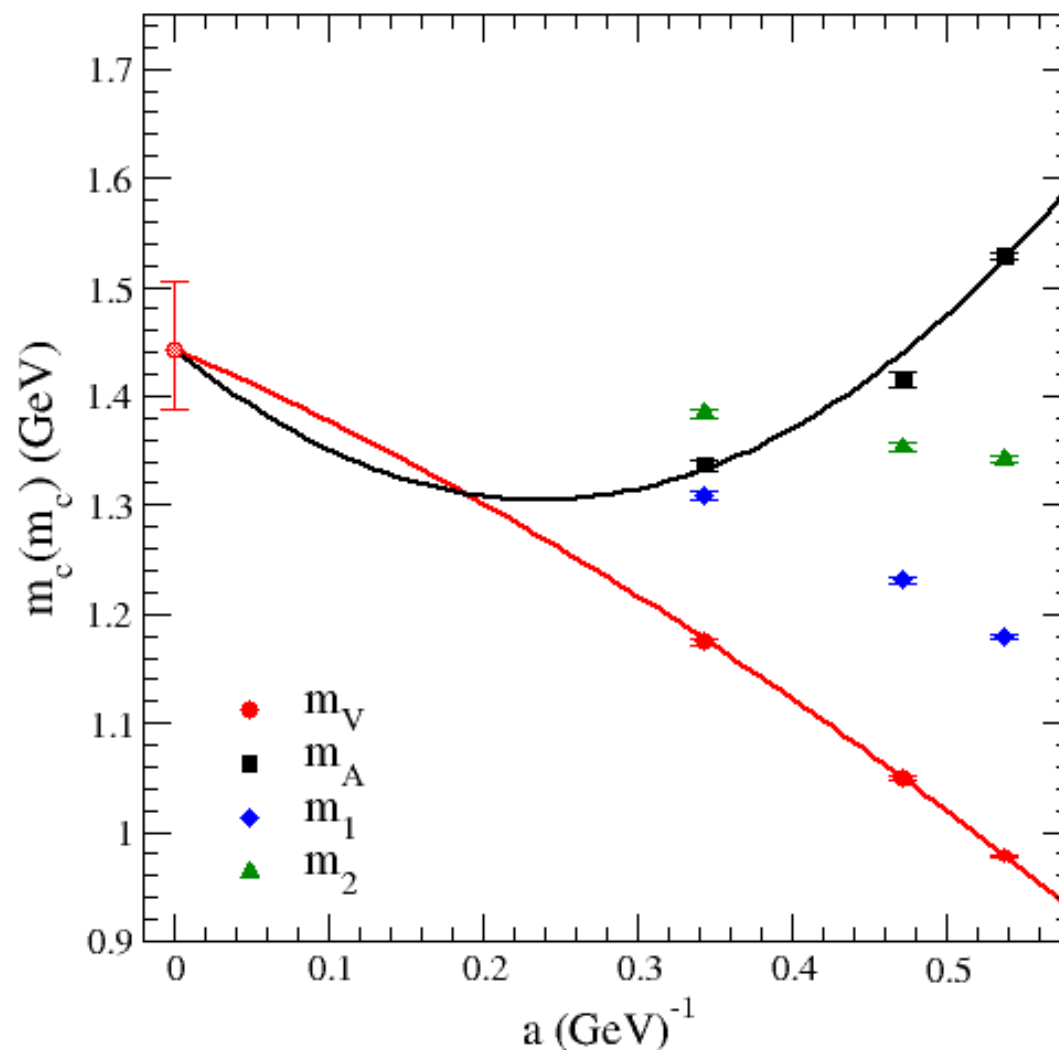




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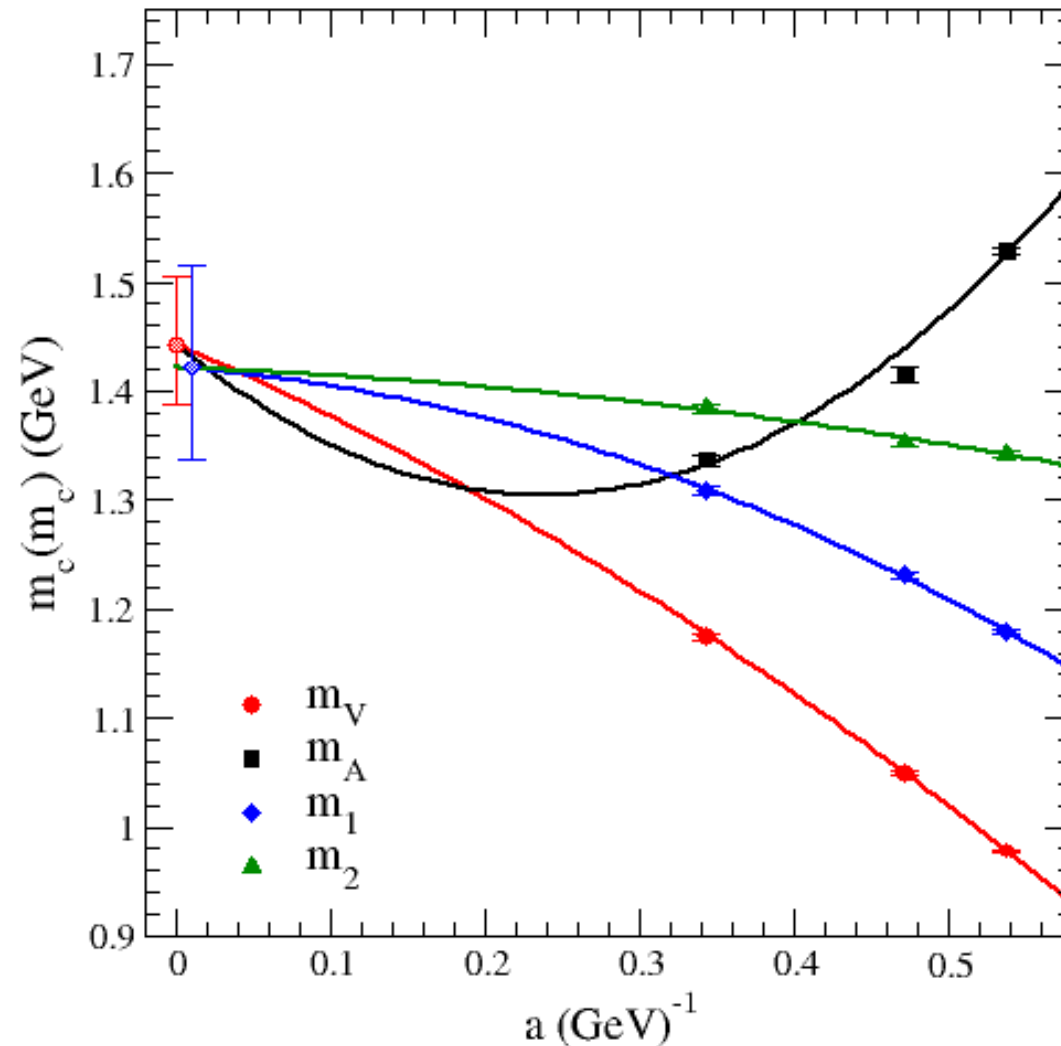




Continuum limit of FNAL masses

Simultaneous fit
to fnal masses

Continuum
limits agree!



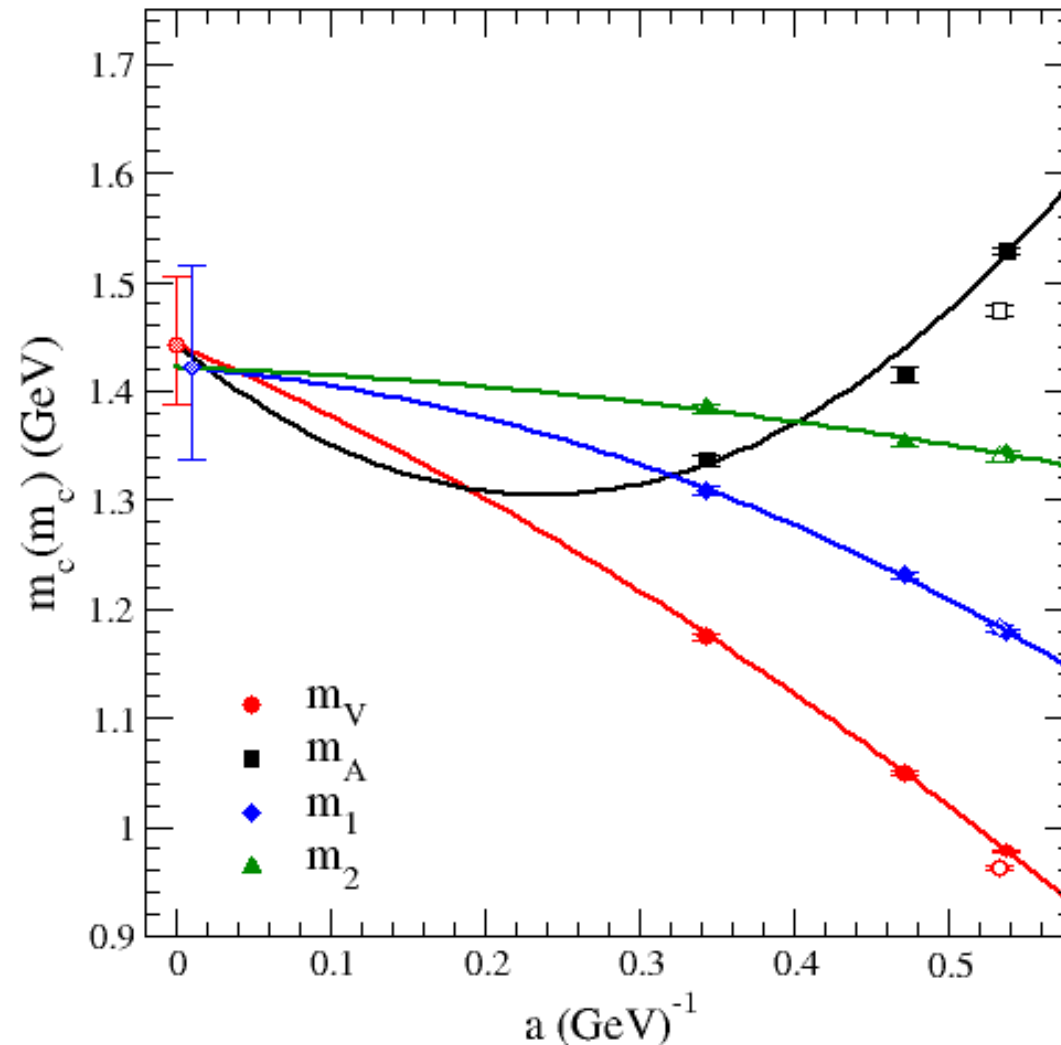


DF data at fixed lattice spacing

Open symbols show DF data

Alpha data quenched is too high at fixed lattice spacing. Lattice artefacts?

No effect for FNAL masses





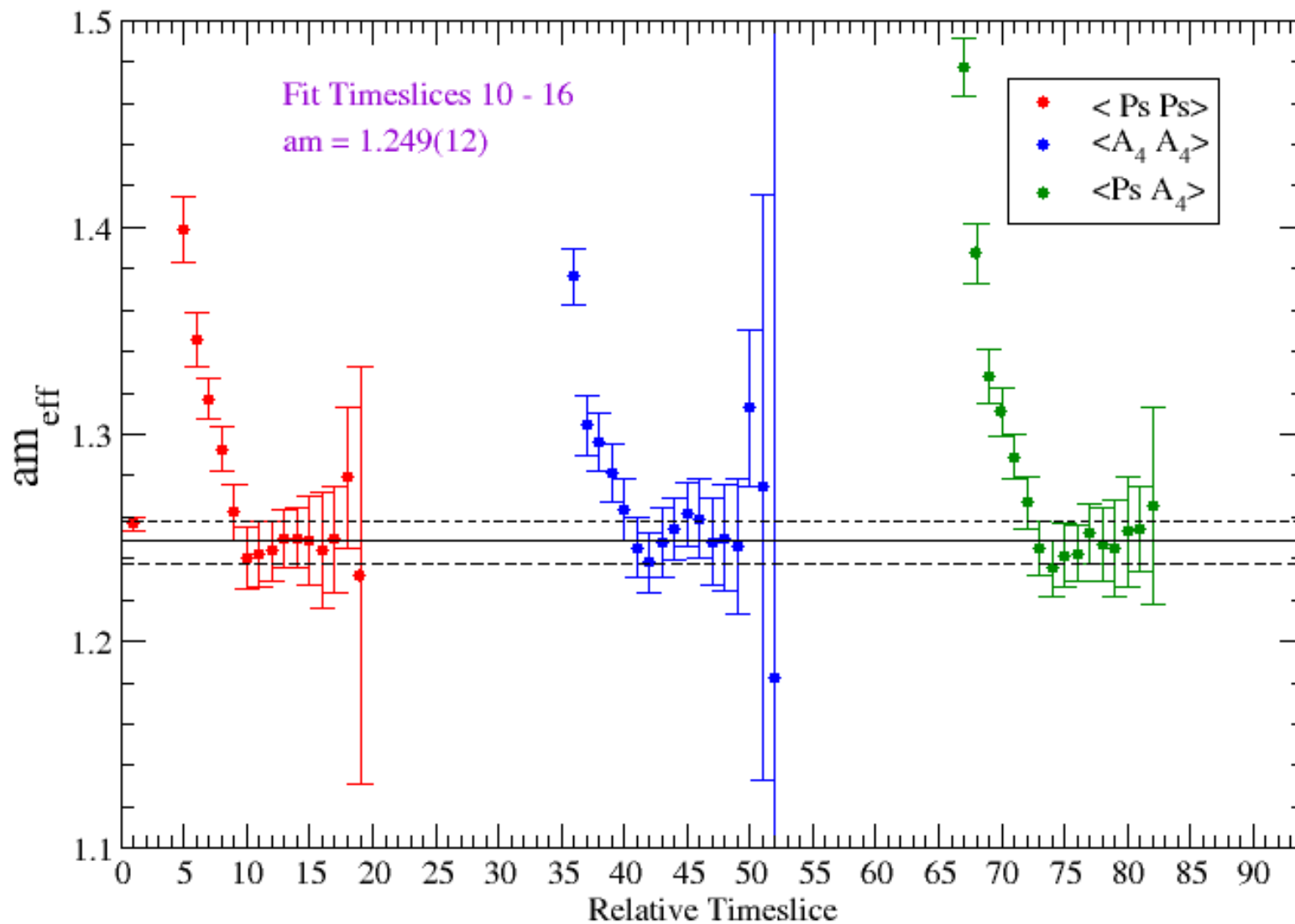
fDs: Overlap on KS sea

- See talk by RDK for background
- $a \sim 0.125\text{fm} \rightarrow am_c \sim 1$
 - Rather coarse lattice spacing
- Overlap lattice artefacts $O(am)^2$
 - \rightarrow largest $(am)^2 \sim 3/4$
- Multi-mass solver
 - Charm quark propagator is free!



Simultaneous Uncorrelated Fit to Three Correlators

10 configs : $am_{sea} = 0.02/0.05$: $a\mu_H = 3000$ $a\mu_1 = 200$





Extracting f_{D_s}

- From Axial Ward Identity
 - Chiral symmetry \rightarrow only Z_A needed

$$Z_A \langle \partial_\mu A_\mu \mathcal{O} \rangle = 2m_q \langle P \mathcal{O} \rangle$$

Can define Z_A in terms of C_{PP} and C_{A4P}

This **cancels** in the definition of f_{D_s}

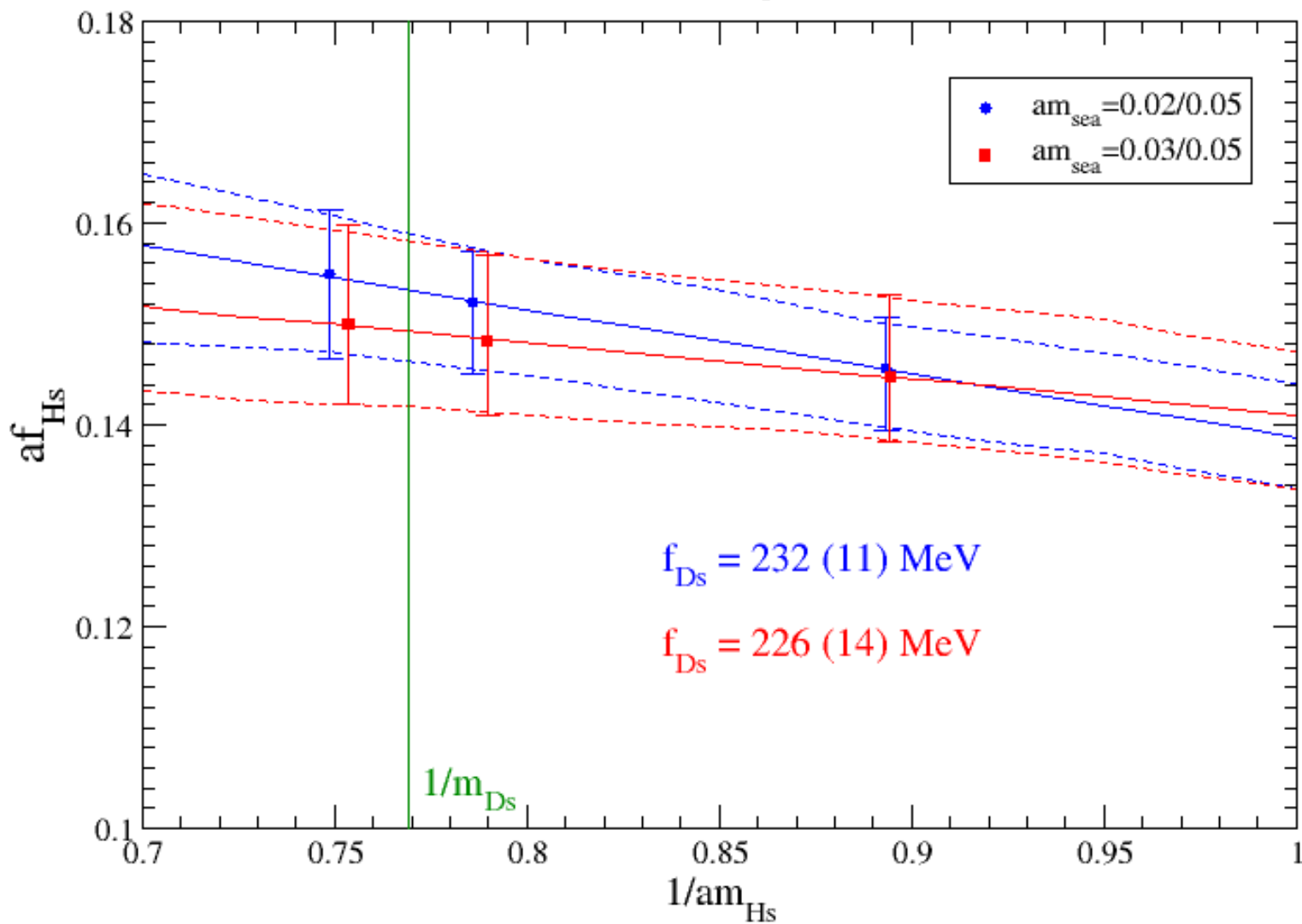
Determine renormalised f_{D_s} from P 's correlator

Operating mixing is under control



Heavy Quark Mass Dependence of f_{D_s}

10 configs





Summary

- QCL: Ds mesons consistent with new charm states
- ALPHA and FNAL m_c same QCL
 - Stat and sys error large from coarse a
- Minimal evidence for sea quark effects
 - Heavy sea / fixed lattice spacing
- HQ from overlap very promising
- GW on GW all above, plus semileptonics