

Mobius Fermions

Improved DW Chiral Fermions

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Outline

- Generalization of Shamir/Borici DW Fermions
- Modified EVEN/ODD Preconditioning
- M_{res} and the Ward-Takahashi identity at finite L_s
- Performance Test on Wilson $16^3 \times 32$, $\beta = 6.0$
- Future Developments

Chiral Breaking Operator For Overlap

Approximate Ginsparg-Wilson relation:

$$\gamma_5 D_{ov}(0) + D_{ov}(0) \gamma_5 - 2D_{ov}(0) \gamma_5 D_{ov}(0) = 2\gamma_5 \Delta_L$$

$$\nabla^\Gamma = \frac{\nabla}{J} (J - \epsilon_L^\Gamma(H))$$

Breaking term in Neether theorem:

$$(1-m)^{-1} \delta(\bar{\psi} D_{ov}(m) \psi) = m_q \bar{\psi} (\gamma_5 + \hat{\gamma}_5) \psi + 2\bar{\psi} \gamma_5 \Delta_L \psi$$

$$m_q = m/(1-m) \quad \text{and} \quad \hat{\gamma}_5 = \gamma_5(1-2D_0) = -\epsilon_L(H)$$

Chiral breaking operator, $\Delta_L(x,y)$, is “local”, not ultra local.

Overlap Solutions to GW Relation

$$D_{ov}(m) = \frac{1+m}{2} + \frac{1-m}{2} \gamma_5 \epsilon_L[\gamma_5 D(M_5)]$$

GOALS OF DOMAIN WALL IMPLEMENTATION

- Best choices for 4-d “kernel”:

$$H \sim \gamma_5 D(M_5)$$

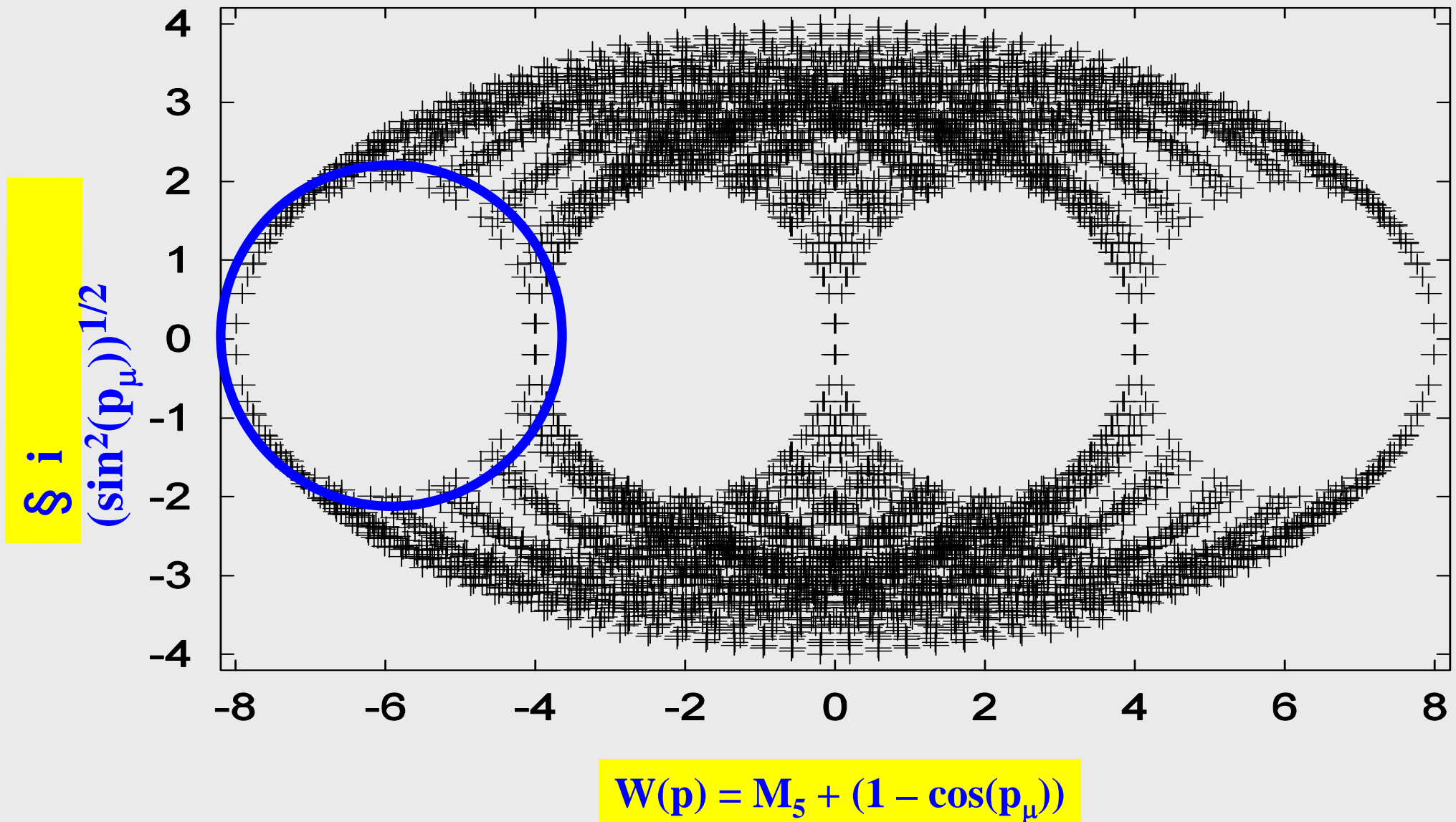
- Best approx of sign function:

$$\epsilon_L[\mathbf{x}] \sim \mathbf{x}/|\mathbf{x}|$$

- Define Chiral breaking :

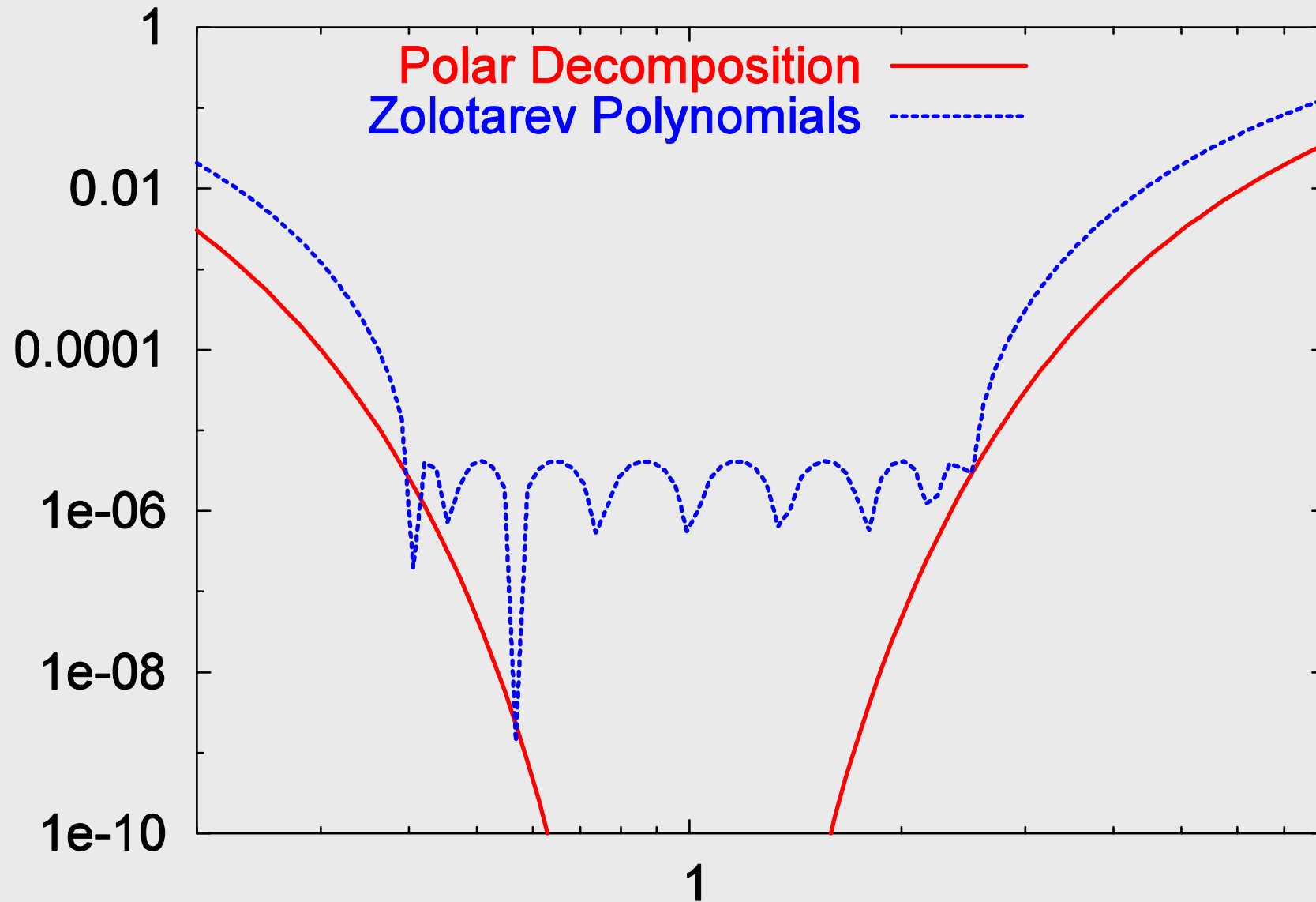
$$\Delta_L[H] = (1 - \epsilon_L^2[H])/4$$

$D_{\text{wilson}}(-4)$ Free Eigenvalues



$1 - \varepsilon_L[x]$ at $L_s = 16$ (polar) vs 8 (Zolotarev)

$\text{abs}(1 - \varepsilon_L(x))$



What is the best way to exploit the 5th time axis?

DW Construction: Shamir vs Borici

$$\text{Shamir: } D(M_5) = \frac{a_5 D_w(M_5)}{2 + a_5 D_w(M_5)}$$

or Borici: $D(M_5) = a_5 D_W(M_5)$

where the standard Wilson 4-d operator is

$$D_w(M) = (4 + M_5)\delta_{x,y} - \frac{1}{2} \sum_{\mu=1}^4 [(1 - \gamma_\mu)U_\mu(x)\delta_{x+\mu,y} + (1 + \gamma_\mu)U_\mu^\dagger(y)\delta_{x,y+\mu}]$$

implies

$$\epsilon_{L_s}[H] = \frac{T^{-L_s} - 1}{T^{-L_s} + 1} \text{ with } T = \frac{1 - H}{1 + H}, \quad H \equiv \gamma_5 D(M)$$

Mobius Generalization

$$\begin{aligned} D_{\text{Mobius}}(M_5) &= \frac{a + b D_w(-1)}{c + d D_w(-1)} \\ &\equiv \frac{(b_5 + c_5) D_w(M_5)}{2 + (b_5 - c_5) D_w(M_5)} \end{aligned}$$

Parameters: M_5 , $a_5 = b_5 - c_5$ and scale: $\alpha = b_5 + c_5$

- M_5 sets the zero for $H = \gamma_5 D(M)$
- a_5 moves the doublers down
- **NEW** α slides e.v. in the window $1 - \varepsilon_L[H]$

Domain Wall Implementation

$L_s \times L_s$ DW Matrix:

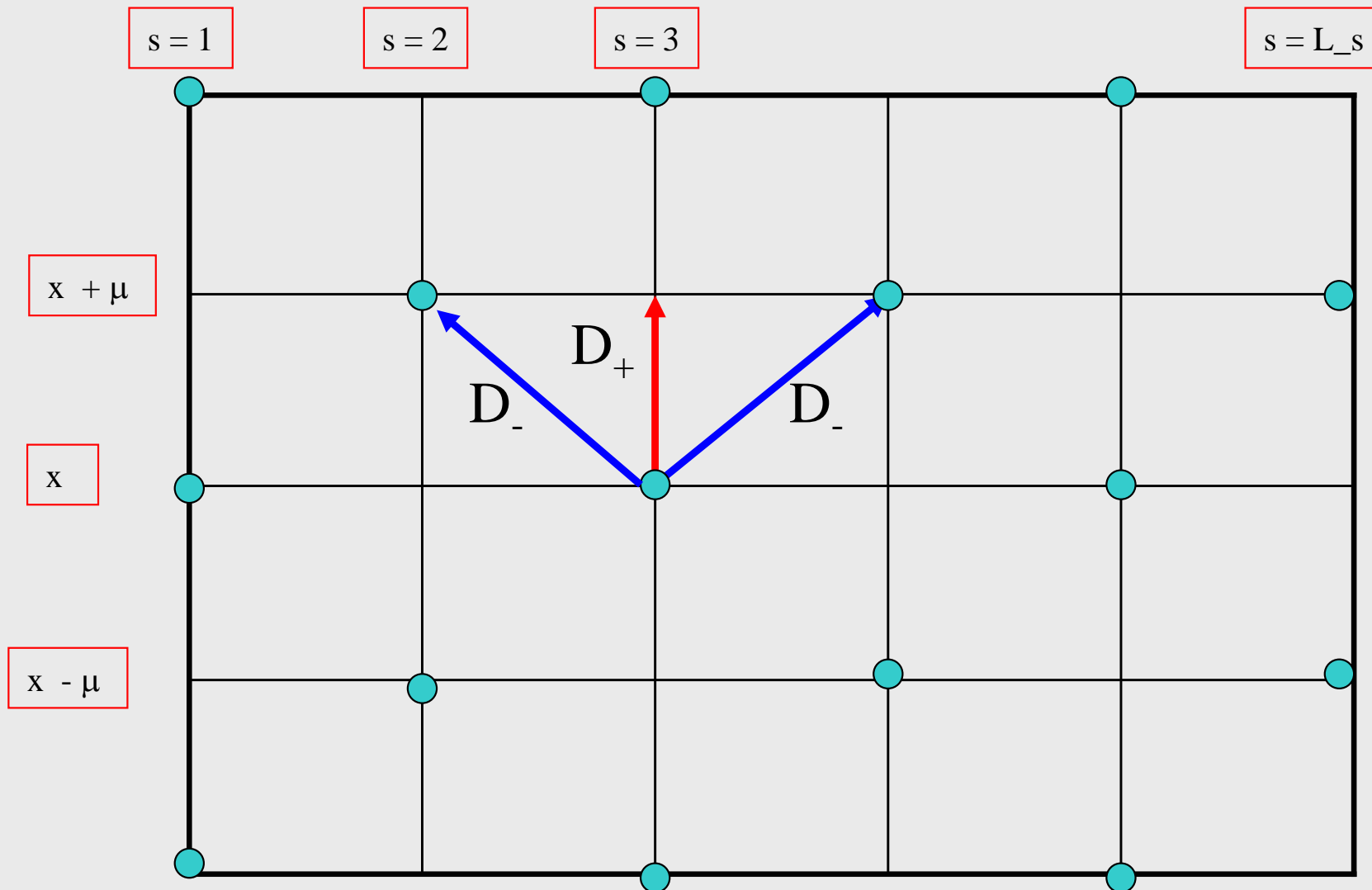
$$D_{s',s}^{DW} = \begin{bmatrix} D_+^{(1)} & D_-^{(1)} P_- & 0 & \dots & -m D_-^{(1)} P_+ \\ D_-^{(2)} P_+ & D_+^{(2)} & D_-^{(2)} P_- & \dots & 0 \\ 0 & D_-^{(3)} P_+ & D_+^{(3)} & \dots & 0 \\ \dots & \dots & \dots & \dots & \dots \\ -m D_-^{(L_s)} P_- & 0 & 0 & \dots & D_+^{(L_s)} \end{bmatrix}$$

with L^4 Wilson Operators:

$$D_+^{(i)} = b_5 \omega_i D_w(M) + 1, \quad D_-^{(i)} = c_5 \omega_i D_w(M) - 1$$

$$P_{\pm} \equiv \frac{1}{2}(1 \pm \gamma_5) \quad (\omega_i = 1, \text{ except for Chiu's Zolotarev form})$$

Mobius generalization of Shamir/Borici



Shamir: $b_5 = a_5, c_5 = 0$

Borici: $b_5 = c_5 = a_5$

$$D_+ = b_5 D_w(M) + 1, \quad D_- = c_5 D_w(M) - 1$$

Standard LDU \rightarrow 4-d Overlap Form

Define Boundary fields at $s = 1, L_s$

$$q_x = P_- \Psi_{x,1} + P_+ \Psi_{x,L_s}$$

$$\bar{q}_y = -[\bar{\Psi}_1 D_-^{(1)}]_y P_+ - [\bar{\Psi}_{L_s} D_-^{(L_s)}]_y P_-$$

$$\langle q_x \bar{q}_y \rangle = \frac{1}{1-m} [D_{ov}^{-1}(m) - 1]_{xy}$$

$$D_{ov} \equiv D_4^{-1}(1) D_4(m) = \left[\frac{1+m}{2} + \frac{1-m}{2} \gamma_5 \frac{T^{-L_s} - 1}{T^{-L_s} + 1} \right],$$

where $T^{-L_s} \equiv T_1^{-1} T_2^{-1} \dots T_{L_s}^{-1}$

$\varepsilon_L[H]$

Edwards & Heller use “Standard” UDL decomposition

$$' D^{DW} \mathcal{P} = U D L(m) '$$

$$D^{DW} \mathcal{P} = [\gamma_5 Q_- U Q_-^{-1}] \quad Q_- \begin{bmatrix} D_4(m) & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix} \quad [L(m)]$$

Step #1: Prepare the Pivots by Permute Columns

$$\mathcal{P} = \begin{bmatrix} P_- & P_+ & 0 & 0 \\ 0 & P_- & P_+ & 0 \\ 0 & 0 & P_- & P_+ \\ P_+ & 0 & 0 & P_- \end{bmatrix}$$

Step #2: Do Gaussian Elimination to get U matrix

$$U = \begin{bmatrix} 1 & -T_1^{-1} & -T_1^{-1}T_2^{-1} & -T_1^{-1}T_2^{-1}T_3^{-1} \\ 0 & 1 & -T_2^{-1} & -T_2^{-1}T_3^{-1} \\ 0 & 0 & 1 & -T_3^{-1} \\ 0 & 0 & 0 & 1 \end{bmatrix}$$

Step #3 Back substitution to get L matrix

$$L(m) = \begin{bmatrix} & -1 & & & 0 & 0 & 0 \\ & -T_2^{-1}T_3^{-1}T_4^{-1}c_+ & & & 1 & 0 & 0 \\ & -T_3^{-1}T_4^{-1}c_+ & & & 0 & 1 & 0 \\ & -T_4^{-1}c_+ & & & 0 & 0 & 1 \end{bmatrix}$$

where

$$Q_-^{-1} = \text{Diag}[(Q_-^{(1)})^{-1}(Q_-^{(2)})^{-1}(Q_-^{(3)})^{-1}(Q_-^{(4)})^{-1}]$$

$$Q_-^{(s)} = \gamma_5[D_-^{(s)}P_+ + D_+^{(s)}P_-] \quad c_- = P_- - mP_+$$

$$Q_+^{(s)} = \gamma_5[D_+^{(s)}P_+ + D_-^{(s)}P_-] \quad c_+ = P_+ - mP_-$$

Generalized γ_5 Hermiticity and All That

- To get all the nice identities for Borici, Chiu and Mobius

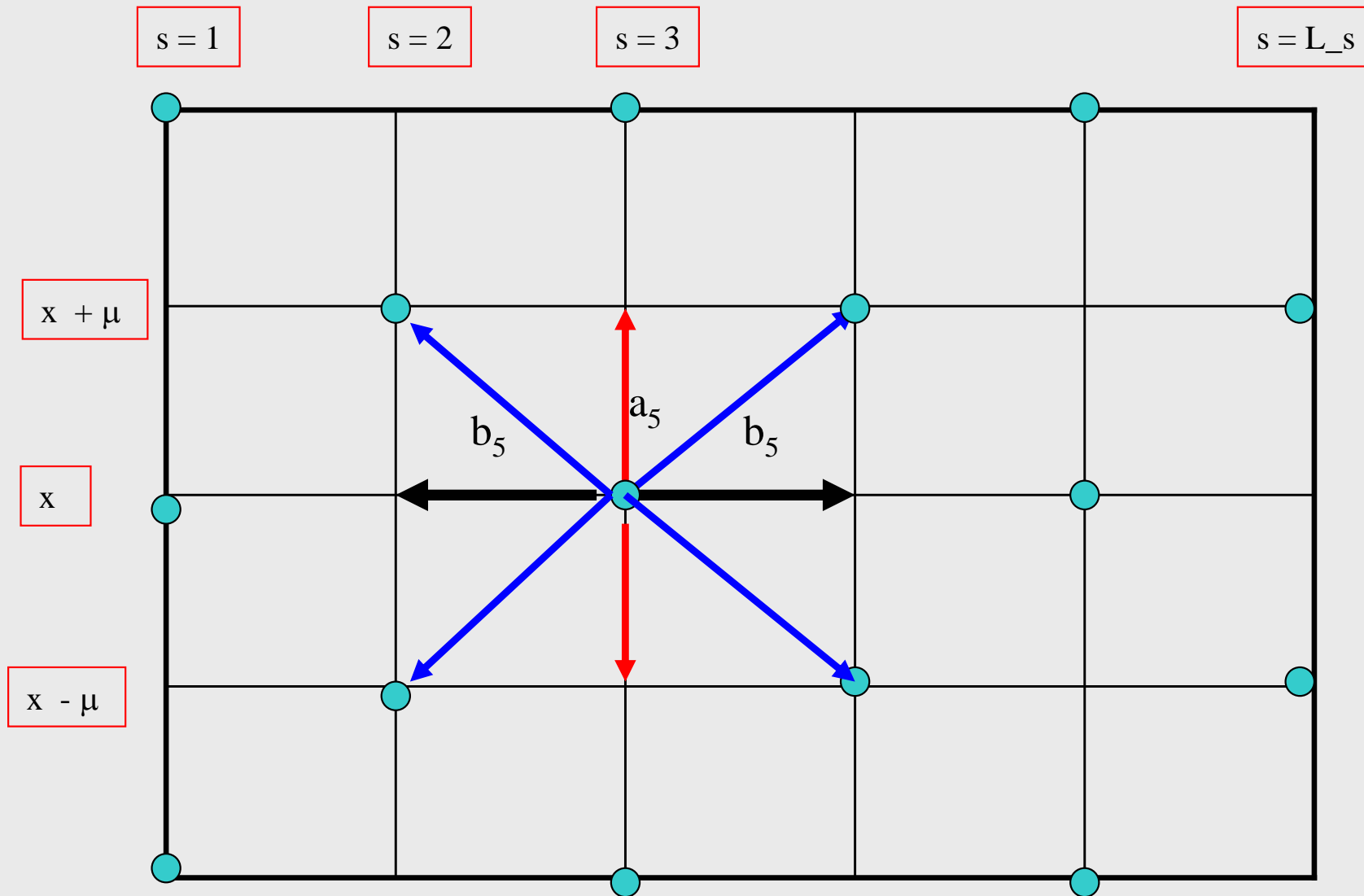
$$D_{DW} = D_- \times \tilde{D}_{DW} \equiv$$

$$\begin{bmatrix} D_-^{(1)} & 0 & 0 & 0 \\ 0 & D_-^{(2)} & 0 & 0 \\ 0 & 0 & D_-^{(3)} & 0 \\ 0 & 0 & 0 & D_-^{(4)} \end{bmatrix} \times \begin{bmatrix} D_+^{(1)}/D_-^{(1)} & P_- & 0 & -mP_+ \\ P_+ & D_+^{(2)}/D_-^{(2)} & P_- & 0 \\ 0 & P_+ & D_+^{(3)}/D_-^{(3)} & P_- \\ -mP_- & 0 & P_+ & D_+^{(4)}/D_-^{(4)} \end{bmatrix}$$

So $\mathcal{R}D_- \gamma_5$ acts like the DW "gamma 5"

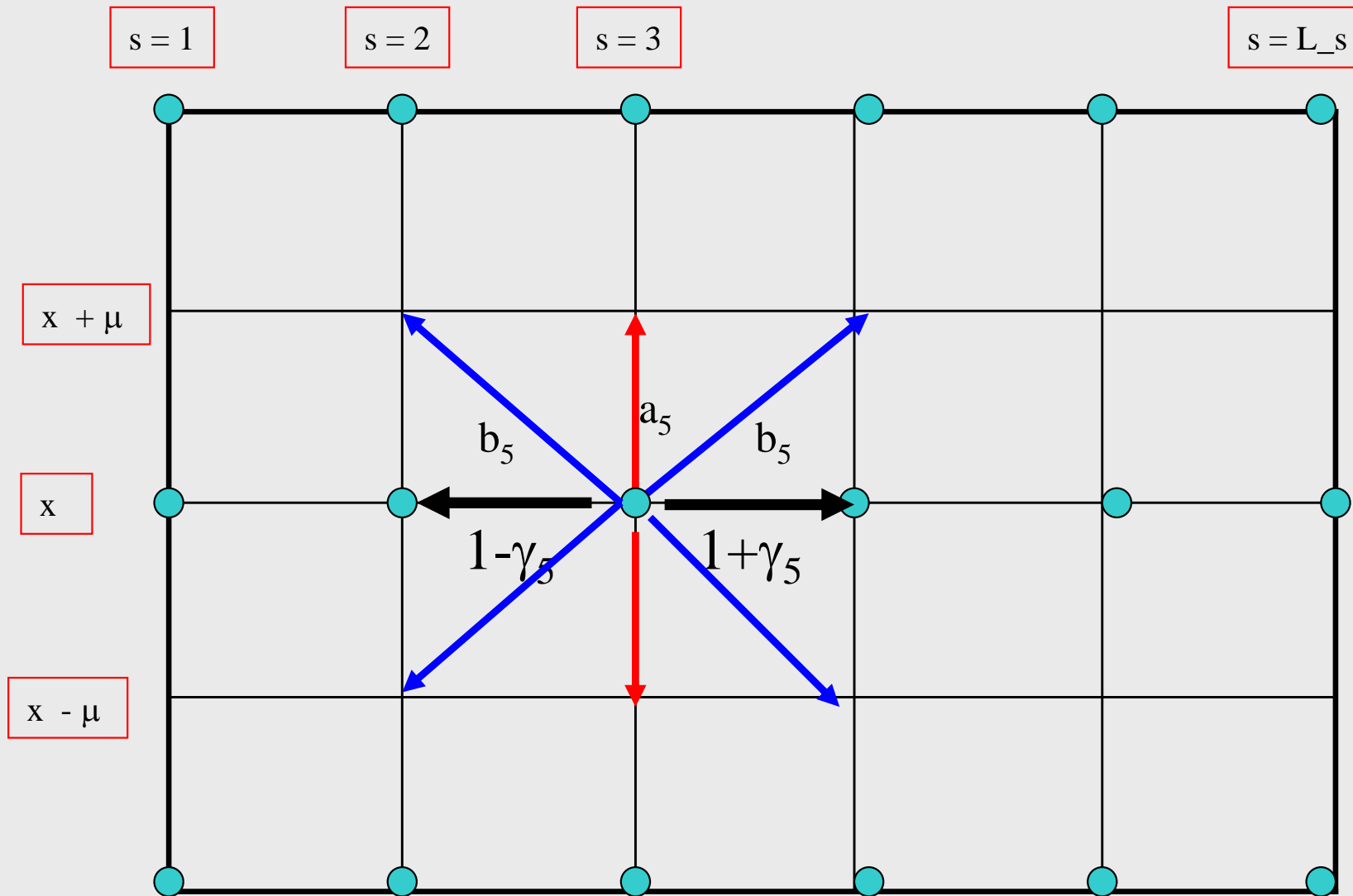
$$\gamma_5 \mathcal{R} \tilde{D}_{DW} = \tilde{D}_{DW}^\dagger \gamma_5 \mathcal{R}.$$

Standard Even/Odd \rightarrow 5d Checker board



Fails for Borici Form $b_5 > 0$

Modified Even/Odd \rightarrow 4-d Checkerboard



Even/Odd Partition of Matrix

$$D_w(M) = \begin{bmatrix} I_{ee} & D_{eo}^{DW'} \\ D_{oe}^{DW'} & I_{oo} \end{bmatrix}$$

I_{ee} and I_{oo} SHOULD be simple to invert

The Schur decomposition

$$D_w(M) = \begin{bmatrix} 1 & 0 \\ D_{oe}^{DW'} I_{ee}^{-1} & 1 \end{bmatrix} \begin{bmatrix} I_{ee} & 0 \\ 0 & I_{oo} - D_{oe}^{DW'} I_{ee} D_{eo}^{DW'} \end{bmatrix} \begin{bmatrix} 1 & I_{ee}^{-1} D_{eo}^{DW'} \\ 0 & 1 \end{bmatrix}$$

$$D_{preconditioned}^{DW} = 1 - I_{oo}^{-1} D_{oe}^{DW} I_{ee}^{-1} D_{eo}^{DW}$$

On $16^3 \times 32$, $\beta = 6.0$ lattice

- For Shamir: Both 4-d & 5-d Even/Odd give » 2.7 speed up.
- For Borici: Even/Odd gives » 2.7 speed up

Ward Takashi Identity For Domain Wall

$$\Delta_{\mu} J_{\mu}^{a, DW}(x) = 2m \bar{q}_x \lambda^a \gamma_5 q_x + 2\bar{Q}_x \gamma_5 \lambda^a Q_x$$

Noether's Theorem:

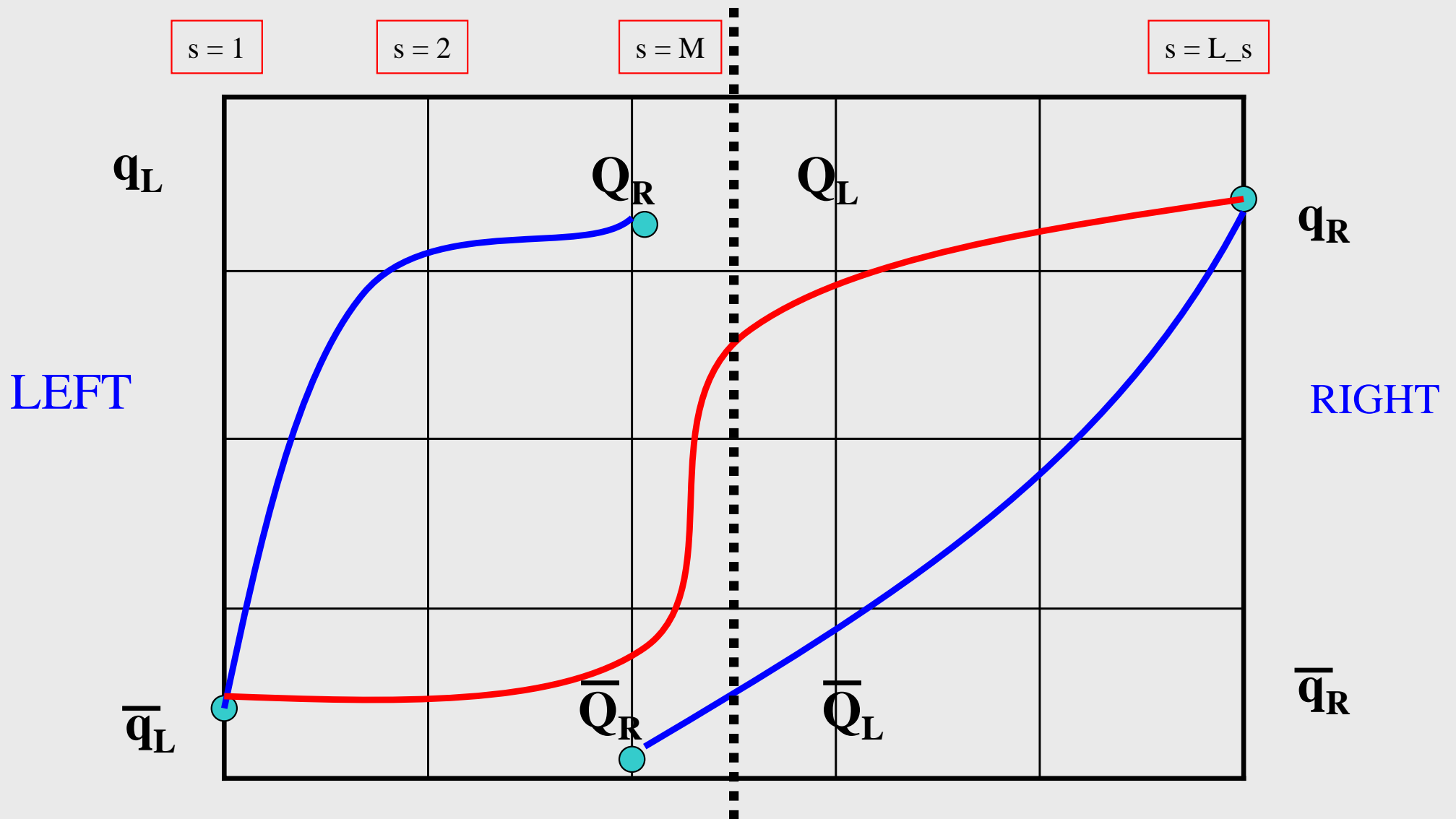
Rotate LEFT $\exp[i\theta_x]$ & RIGHT $\exp[-i\theta_x]$

- **Is this breaking term UNIQUE?**
- **Is it the same as the overlap operators $\Delta_L(\mathbf{x})$?**

Nice Definition of Overlap Axial Current:

$$\langle J_{\mu}(x) \psi_y \bar{\psi}_z \rangle_c \equiv (1 - m) \langle J_{\mu}^{DW}(x) q_y \bar{q}_z \rangle_c$$

Split Screen Correlators



$$\langle Q_x \bar{q}_y \rangle$$

$$\langle q_x \bar{q}_y \rangle = \frac{1}{1-m} [D_{ov}^{-1}(m) - 1]_{xy}$$

$$\langle q_x \bar{Q}_y \rangle$$

Split Screen Propagators

$$\langle Q_s \bar{q} \rangle = \frac{T_{s+1}^{-1} \cdots T_{L_s}^{-1}}{1 + T^{-L_s}} D_{ov}^{-1}(m)$$

$$\langle q \bar{Q}_s \rangle = D_{ov}^{-1}(m) \gamma_5 \frac{1}{1 + T^{-L_s}} [T_1^{-1} \cdots T_s^{-1}] \gamma_5$$

where $s = M$ plane

$$Q_s = P_- \Psi_{s+1} + P_+ \Psi_s$$

$$\bar{Q}_s = -\bar{\Psi}_{s+1} D_-^{(L_s/2+1)} P_+ - \bar{\Psi}_{L/2} D_-^{(s)} P_-$$

y if $s = L_s/2$: $\gamma_5 \langle Q_s \bar{q} \rangle^\dagger \gamma_5 = \langle q \bar{Q}_s \rangle$

Measuring the Operator Δ_{Ls}

(use Plateau region away from sources)

$$m_{res}(t) \equiv \frac{\sum_x \langle \bar{Q}_{t,x} \gamma_5 Q_{t,x} \bar{q}_0 \gamma_5 q_0 \rangle_c}{(1 - m)^2 \sum_x \langle \bar{q}_{t,x} \gamma_5 q_{t,x} \bar{q}_0 \gamma_5 q_0 \rangle_c}$$

Sum over t \rightarrow Measure Matrix element of Δ_L operator

$$m_{res} \equiv \frac{Tr[\Delta_L(H) D_{ov}^{-1} D_{ov}^{\dagger -1}]}{Tr[D_{ov}^{-1} D_{ov}^{\dagger -1}]} = \sum_{\lambda} \rho(\lambda) \Delta_L(\lambda)$$

$|\lambda \rangle$ in the Eigen basis of $\mathbf{H} = \gamma_5 \mathbf{D}(-\mathbf{M})$

Model for m_{res} dependence on L_s

$$m_{res} = \sum_{\lambda} \rho(\lambda) \Delta_L(\lambda)$$

$$\rho(\lambda) = \frac{\langle \lambda | G_{ov} G_{ov}^\dagger | \lambda \rangle}{\sum_{\lambda} \langle \lambda | G_{ov} G_{ov}^\dagger | \lambda \rangle} \geq 0$$

$$\Delta_L(\lambda) = \langle \lambda | \Delta_L(H) | \lambda \rangle = \frac{4}{2 + [\frac{1+\lambda}{1-\lambda}]^{-L} + [\frac{1+\lambda}{1-\lambda}]^L} \geq 0$$

$\rightarrow e^{-L \log[(1+\lambda)/(1-\lambda)]}$ for $O(L^{-1}) < \lambda < O(L)$

Scaling Model: Assume that $\rho(\lambda)$ has negligible dependence on α and L_s

$$m_{res} \simeq \int dn(\lambda) \rho(\lambda) \Delta_L(\alpha\lambda)$$

Derivation:

$$m_{res} \equiv \frac{\sum_x \langle \bar{Q}_x \gamma_5 Q_x \bar{q}_0 \gamma_5 q_0 \rangle_c}{(1 - m)^2 \sum_x \langle \bar{q}_x \gamma_5 q_x \bar{q}_0 \gamma_5 q_0 \rangle_c} \Rightarrow$$

$$\begin{aligned} & \frac{\sum_x Tr[\langle Q_x \bar{q}_0 \rangle \gamma_5 \langle q_0 \bar{Q}_x \rangle \gamma_5]}{(1 - m)^2 \sum_x Tr[\langle q_x \bar{q}_0 \rangle \gamma_5 \langle q_0 \bar{q}_x \rangle \gamma_5]} = \frac{\sum_x Tr[\Delta_{xy}^R \langle q_y \bar{q}_0 \rangle \langle q_z \bar{q}_0 \rangle^\dagger \Delta_{zx}^L]}{\sum_x Tr[\langle q_x \bar{q}_0 \rangle \langle q_x \bar{q}_0 \rangle^\dagger]} \\ & = \frac{\sum_{zy} \langle \bar{q}_z \gamma_5 \Delta_{zy} q_y \bar{q}_0 \gamma_5 q_0 \rangle}{\sum_x \langle \bar{q}_x \gamma_5 q_x \bar{q}_0 \gamma_5 q_0 \rangle} = \frac{Tr[\Delta_L D_{ov}^{-1} D_{ov}^{\dagger-1}]}{Tr[D_{ov}^{-1} D_{ov}^{\dagger-1}]} \simeq \frac{\langle 0 | \bar{q} \gamma_5 \Delta q | \pi \rangle}{\langle 0 | \bar{q} \gamma_5 q | \pi \rangle} \end{aligned}$$

where

$$\begin{aligned} \Delta^L \Delta^R &= \frac{(T_{L_s/2+1} \dots T_{L_s})^{-1}}{1 + T^{-L_s}} \times \frac{(T_1 \dots T_{L_s/2})^{-1}}{1 + T^{-L_s}} \\ &= \frac{T^{-L_s}}{(1 + T^{-L_s})^2} = \frac{1}{2 + T^{L_s} + T^{-L_s}} = \Delta_{L_s} \end{aligned}$$

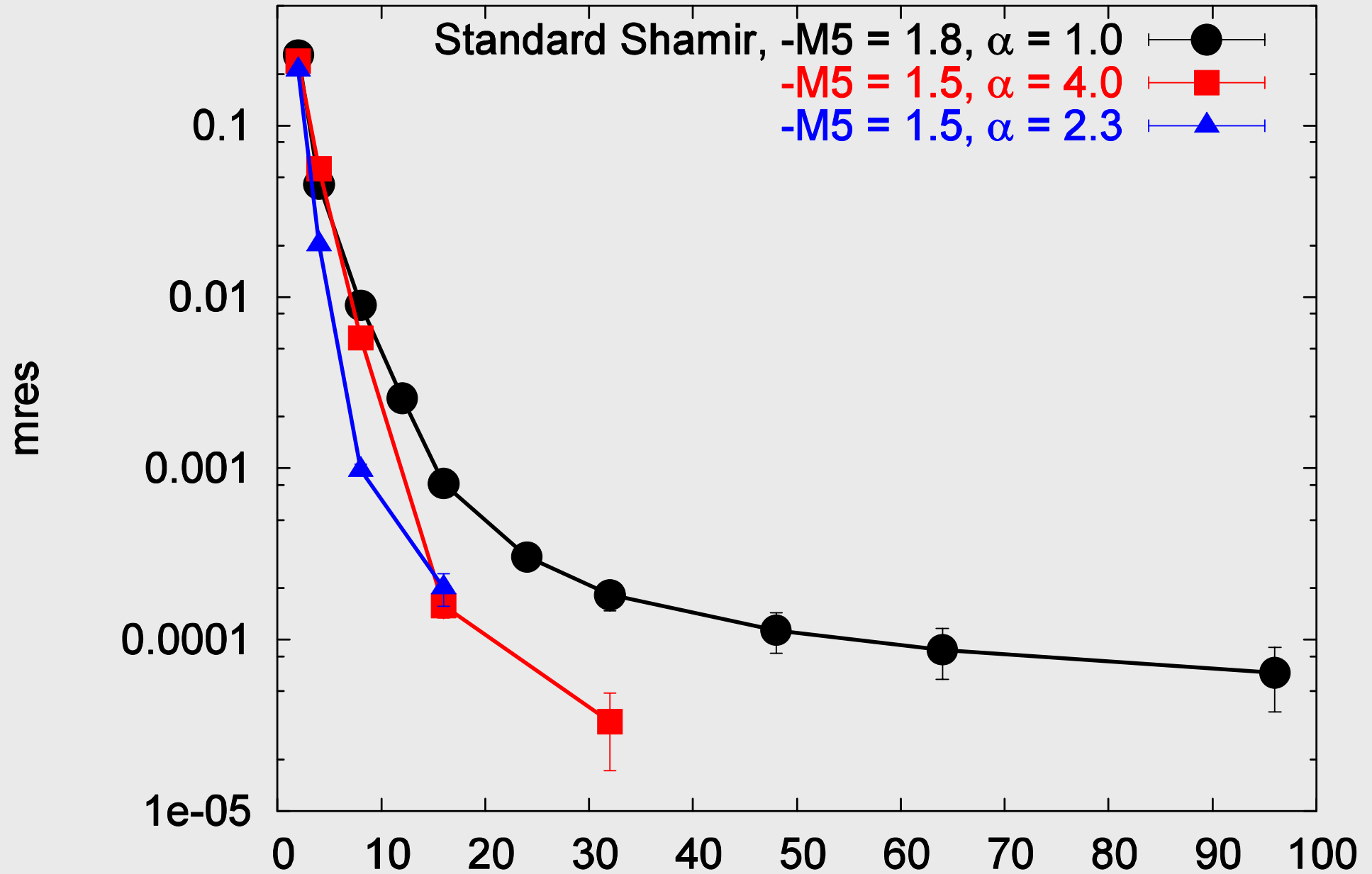
M_{res} is Good Measure of Chiral Symmetry

- M_{res} is independent of Placement of Split Screen
- Δ_L is local positive definite operator.
- $m_{\text{res}} = 0 \rightarrow$ GW and Chiral Symmetry is EXACT
- Low energy matrix elements of Δ_L correct the effective Chiral Lagrangian. To leading $O(a)$ this is a just a shift in the quark mass.

Preliminary Performance Tests

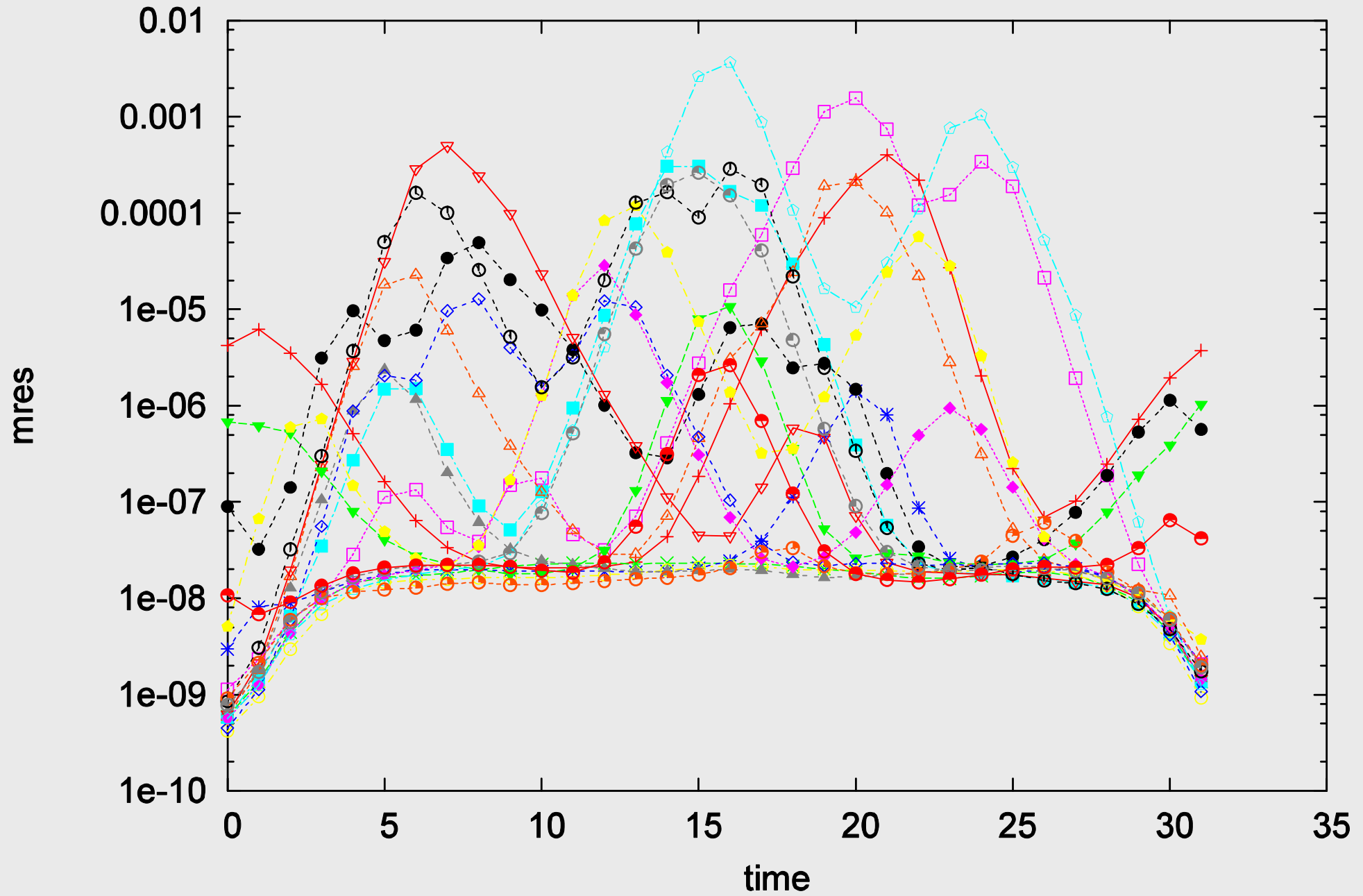
- 20 NERSC Gauge lattices: $16^3 \times 32$ at $\beta = 6.0$
- Set $m_q = 0.06$ for Shamir ($a_5 = 1, M_5 = 1.8$)
- As you vary M_5 and $\alpha = b_5 + c_5$ to fixed m_π

Shamir vs Mobius



$m_{\text{res}}(t)$ for large α

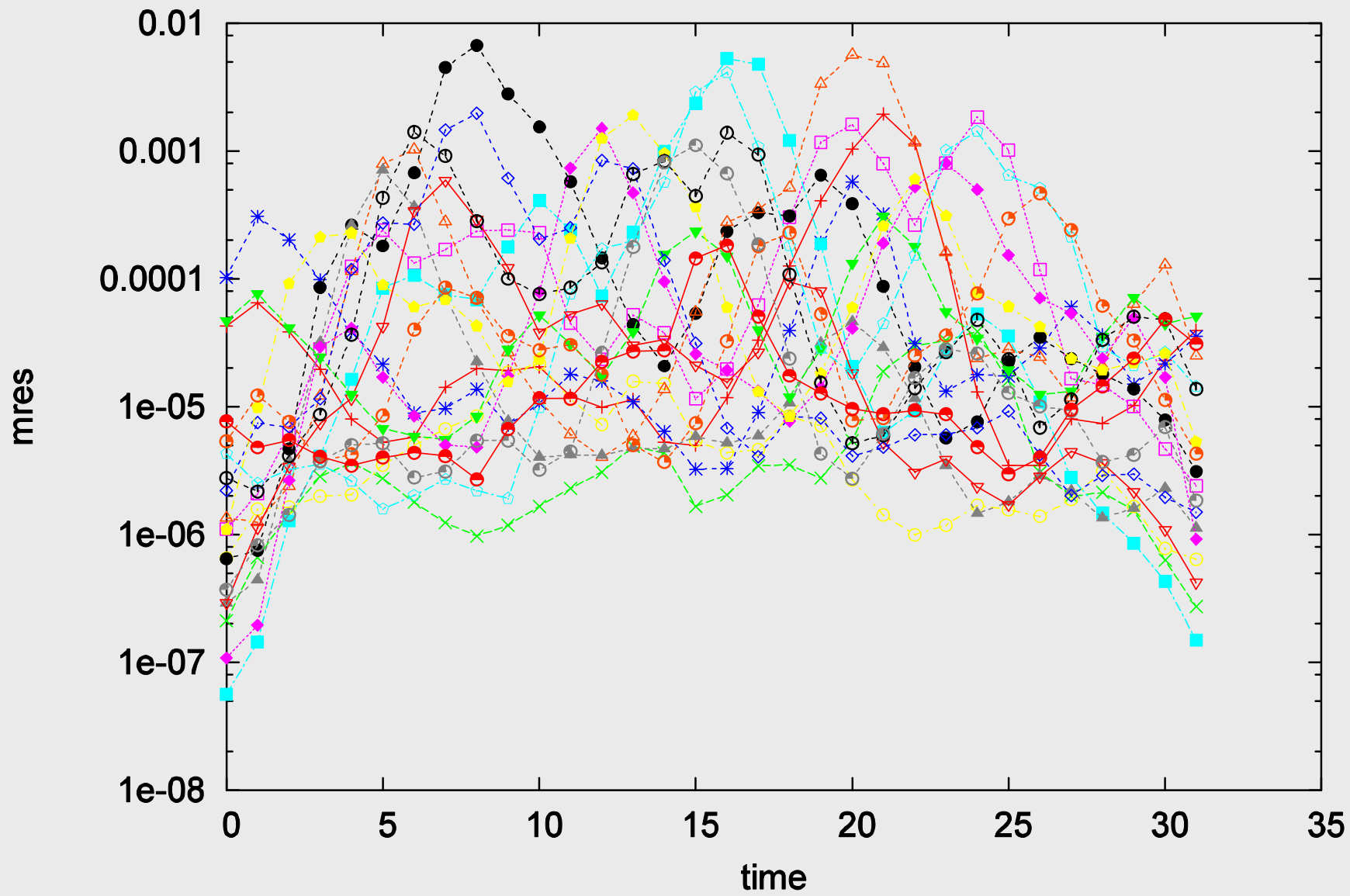
$L_s=32, \alpha = 4.0$



Future Refinements and Projects

- Further tuning in M_5 , a_5 and **scale** parameters
- Heuristic for m_{res} : Error $|1 - \varepsilon_L(\lambda)|$ for lowest e. v. ?
- Test on HMC Lattices (RBC collaboration)
- Next to nearest neighbor ?
- Multi-grid in L_s ? (with C. Rebbi & David Keyes)

$L_s=16, \alpha = 2.3$



Low Statistics Large L_s Behavior

