

Neutron Electric Dipole Moment(NEDM)

Toward the calculation on the lattice

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for CP-PACS collaboration

theta term in QCD

CP odd

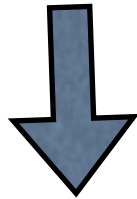
$$i\theta \frac{1}{32\pi^2} \epsilon^{\mu\nu\alpha\beta} F_{\mu\nu}(x) F_{\alpha\beta}(x) \equiv i\theta q(x)$$

NEDM: experimental bound

$$|\vec{d}_n| \leq 6.3 \times 10^{-26} e \cdot cm$$

Model estimate

$$|\vec{d}_n|/\theta \simeq 10^{-15} \sim 10^{-17} e \cdot cm$$



$$\theta = \theta_{\text{QCD}} + \theta_{\text{EW}} \leq O(10^{-8}) \quad \text{Strong CP problem}$$

1st principle calculation of NEDM from lattice QCD

Some History

[Aoki-Gocksch](#) quenched QCD with Wilson fermion $me^{i\theta\gamma_5}$

Incorrect  [Aoki, Gocksch, Manohar, Sharpe](#)

 twisted mass QCD

 some of model estimates are also incorrect
New estimate by ChPT ([Aoki-Hatsuda](#))

[Kuramashi](#) (unpublished) Wilson fermion $\bar{\psi}(x)\gamma_5\psi(x) \Leftrightarrow q(x)$
No signal !

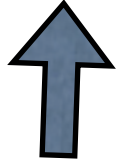
Anomalous flavor singlet chiral WT identities

theta term \Leftrightarrow complex mass term $me^{i\theta\gamma_5}$
 $|\vec{d}_n| = 0$ at $m = 0$

Chiral symmetry is important \Rightarrow DW/GW fermions !

Method for small theta

$$\vec{d}_n \propto \theta \int d^3x \langle N | \vec{x} J_0^{EM}(\vec{x}) Q | N \rangle$$



$$\langle N(p) | J_\mu^{EM}(q = p - p') Q | N(p') \rangle \quad \text{Form factor}$$

How can I extract this form factor from lattice data ?

Naive reduction formula can not be used. (Ishizuka)

Formulation at finite θ

Parity-even Nucleon interpolating operators

$$\begin{aligned}\chi_+(\vec{0}, 0) &= Z_+^\theta N_+^\theta(\vec{0}, 0) + Z_-^\theta \gamma_5 N_-^\theta(\vec{0}, 0) \\ &= Z_+ N_+(\vec{0}, 0) + Z_- \gamma_5 N_-(\vec{0}, 0)\end{aligned}$$

$$N_\pm^\theta(\vec{x}, t) = \sum_{\vec{p}, s} [a_\pm^\theta(\vec{p}, s) u_\pm^\theta(\vec{p}, s) e^{i p x} + \dots]$$

$$N_\pm(\vec{0}, 0) = \sum_{\vec{p}, s} [a_\pm(\vec{p}, s) u_\pm(\vec{p}, s) + \dots]$$

Assumption

$$e^{-i\theta Q_5} a_\pm(P) e^{i\theta Q_5} = a_\pm(P) - i\theta C_\mp(P, K) a_\mp(K) + O(\theta^2)$$

$$\text{Ex. } C_\pm(P, K) = c \delta_{P.K}$$

Q_5 : Anomalous chiral charge

$$\text{(Note } a_\pm^\theta(P) = e^{i\theta Q_5} a_\pm(P) e^{-i\theta Q_5} \text{)}$$

Consequence of the assumption on 2-pt function

$$|\theta\rangle = e^{i\theta Q_5} |0\rangle$$

$$\begin{aligned} G_{NN}^\theta(\vec{p}, t) &\equiv \sum_{\vec{x}} \langle \theta | \chi_+(\vec{0}, 0) \bar{\chi}_+(\vec{x}, t) | \theta \rangle e^{i\vec{p}\vec{x}} \\ &= \sum_s \langle \theta | \chi_+(\vec{0}, 0) | N_+^\theta(P) \rangle \langle N_+^\theta | \bar{\chi}_+(\vec{0}, 0) | \theta \rangle e^{-E_+^\theta t} + \dots \end{aligned}$$

$$\begin{aligned} \langle \theta | \chi_+(\vec{0}, 0) | N_+^\theta(P) \rangle &= Z_+^\theta u_+^\theta(P) \\ &= Z_+ \langle \theta | N_+ | N_+^\theta(P) \rangle + Z_- \gamma_5 \langle \theta | N_- | N_+^\theta(P) \rangle \end{aligned}$$

$$\langle \theta | N_+ | N_+^\theta(P) \rangle = \langle 0 | e^{-i\theta Q_5} N_+ e^{i\theta Q_5} | N_+(P) \rangle = u_+(P) + O(\theta^2)$$

$$\begin{aligned} \langle \theta | N_- | N_+^\theta(P) \rangle &= \langle 0 | e^{-i\theta Q_5} N_- e^{i\theta Q_5} | N_+(P) \rangle \\ &= -i\theta C_-(P, K) u_-(K) + O(\theta^3) \end{aligned}$$

CP symmetry at $\theta = 0$ \longrightarrow $Z_+^\theta = Z_+ + O(\theta^2)$

Therefore

$$u_+^\theta(P) = u_+(P) - i\theta\gamma_5 D(P, K)u_-(K) + O(\theta^2)$$

$$D(P, K) \equiv \frac{Z_-}{Z_+} C_-(P, K)$$

After a little algebra

$$G_{NN}^\theta(\vec{p}, t) = G_{NN}(\vec{p}, t) + i\theta G_{NN}^Q(\vec{p}, t) + O(\theta^2)$$

$$G_{NN}(\vec{p}, t) = |Z_+|^2 e^{-E_+ t} \frac{E_+ \gamma_0 + m_+ - i\vec{\gamma} \cdot \vec{p}}{2E_+}$$

$$G_{NN}^Q(\vec{p}, t) = - |Z_+|^2 e^{-E_+ t} \gamma_5 [D_P + i\gamma_0 D_0 - \vec{\gamma} \cdot \vec{p} D_k]$$

$$G_{NN}^Q(\vec{p}, t) = \langle 0 | \chi(0) \bar{\chi}(t) Q | 0 \rangle$$

$D_{P,0,k}$: real constants

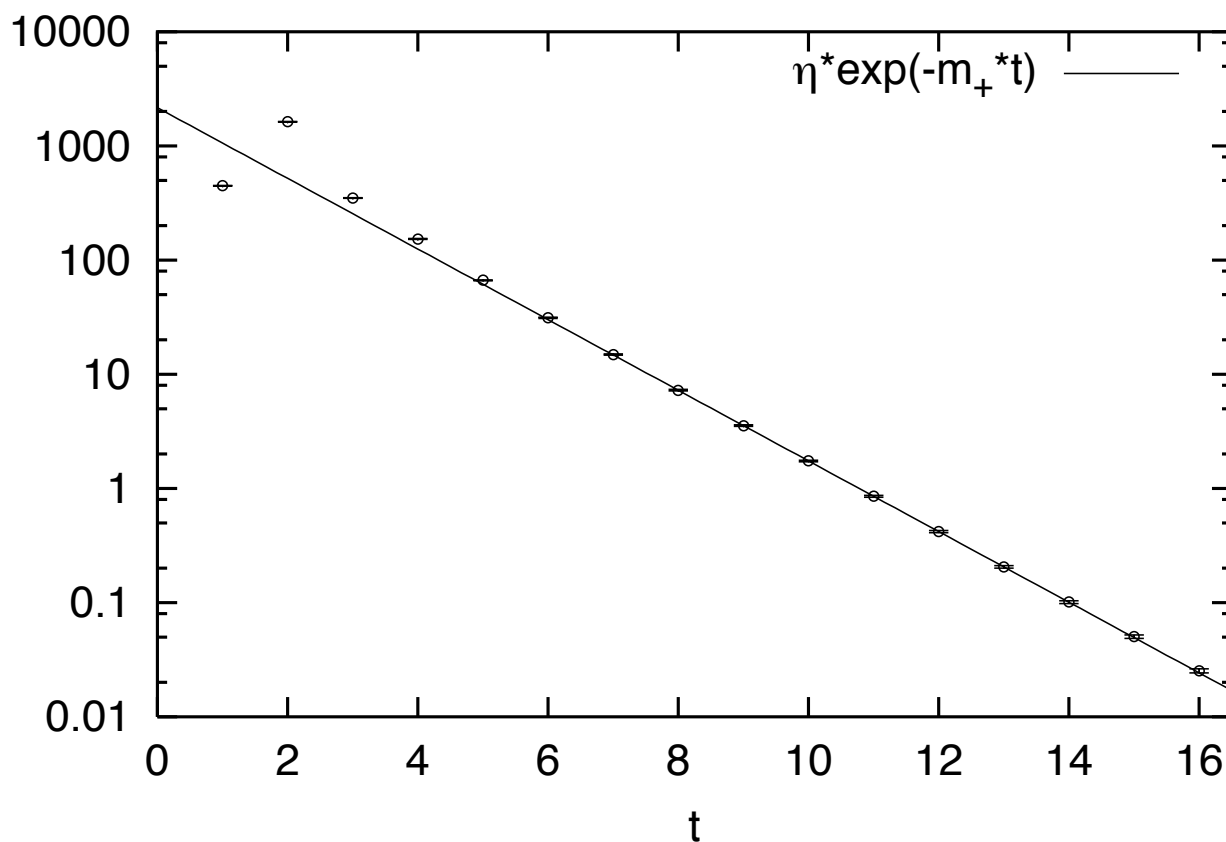
Numerical Results

Iwasaki-RG(Quenched) at $\beta = 2.6(a \simeq 0.1\text{fm}) + \text{DWF}$

$16^3 \times 32 \times 16$, $m_f a = 0.03$, 716 conf.

$$G_{NN}(\vec{0}, t)$$

$\langle N_+ N_+ \rangle(p)$, $p^2=0$, fit range [9:13]

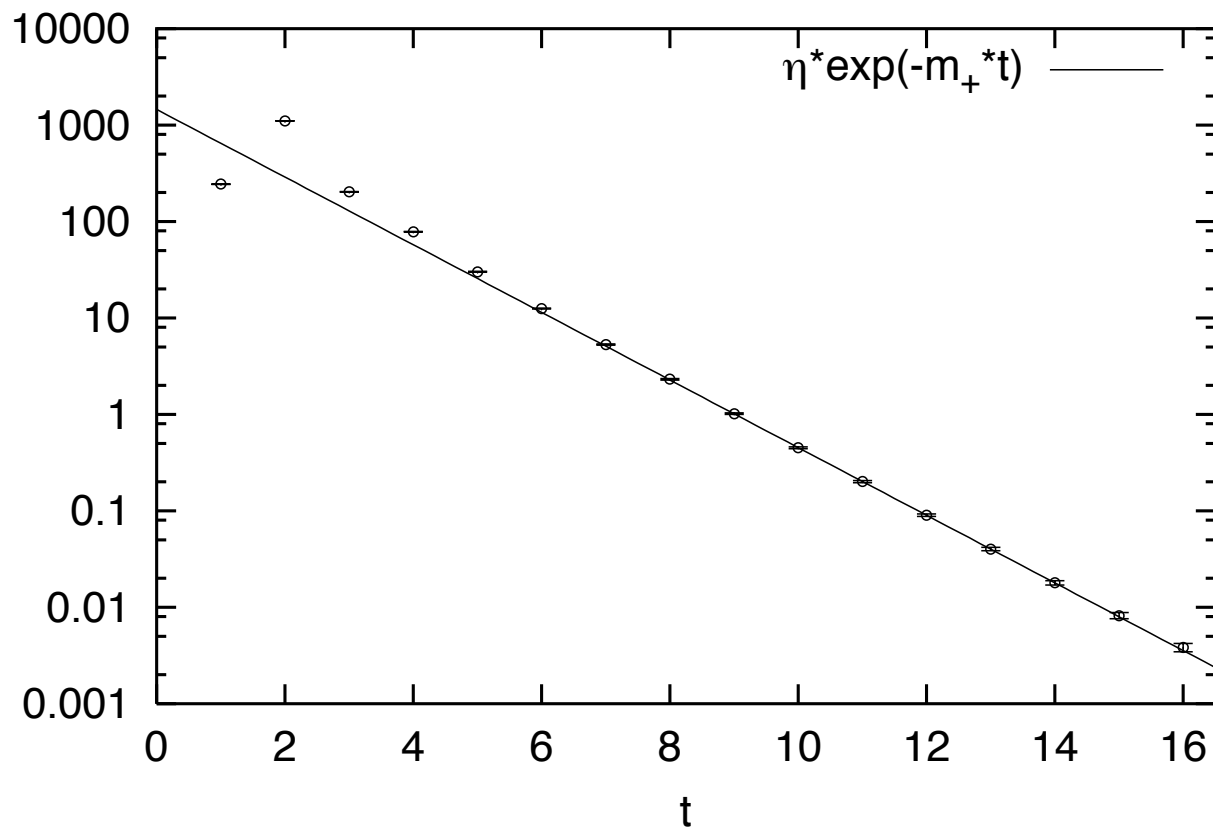


$$m_+ = 0.712(4)$$

$$Z_+^2 = 2162(87)$$

$$G_{NN}(\vec{p}, t) \quad |\vec{p}| = \frac{2\pi}{L}$$

$\langle N_+ N_+ \rangle(p), p^2=1, \text{ fit range } [9:13]$



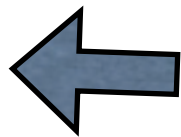
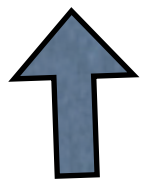
$$E_+ = 0.807(8)$$

$$(\sqrt{m_+^2 + \vec{p}^2} = 0.813)$$

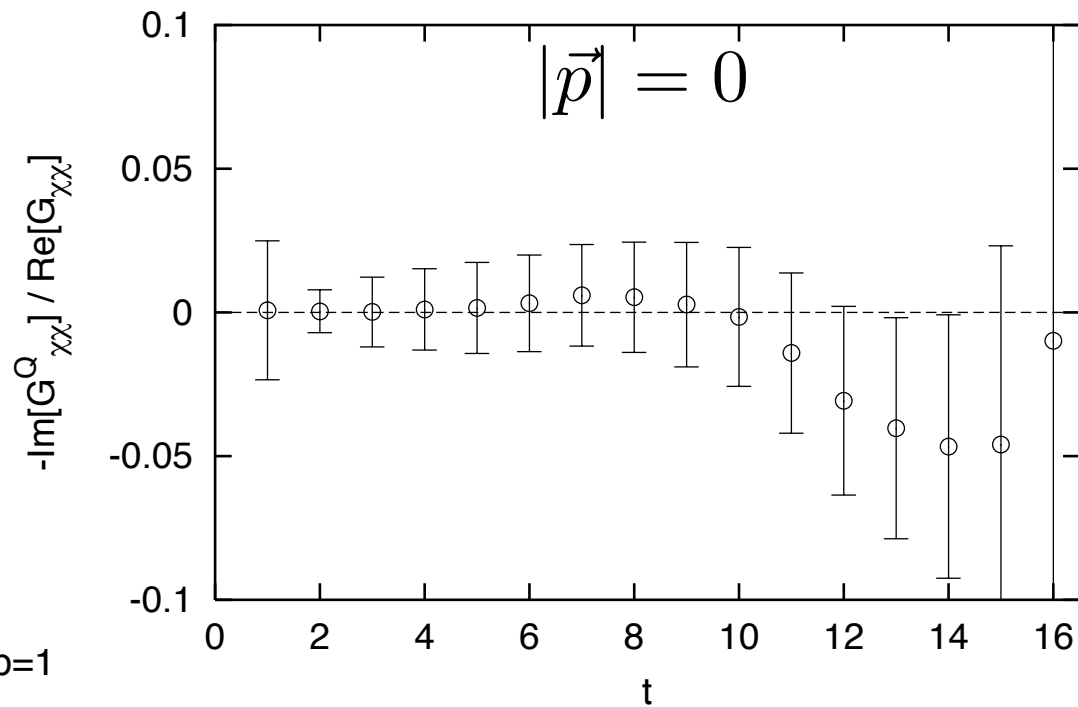
$$Z_+^2 = 1453(104)$$

$$\text{Im} \left\{ \frac{\text{tr}[G_{NN}^Q(\vec{p}, t)\gamma_5\Gamma_0]}{\text{tr}[G_{NN}(\vec{p}, t)\Gamma_0]} \right\}$$

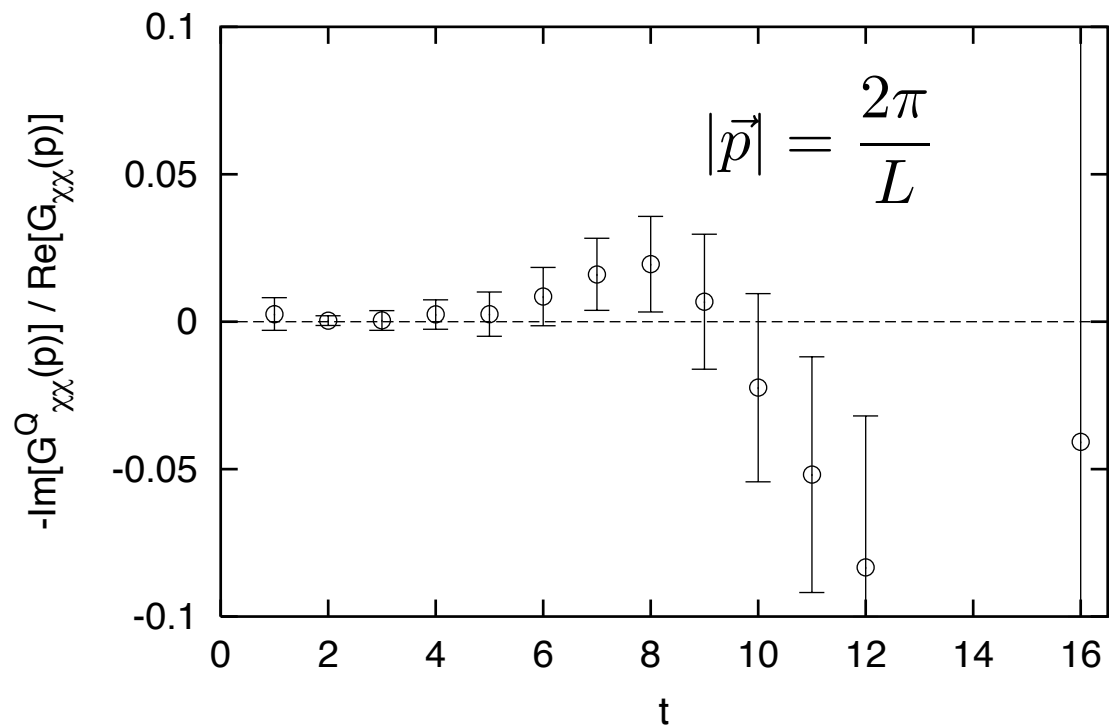
$$D_0 \simeq 0$$



RG $\beta=2.6, 16^3 \times 32 \times 16, A=1.28, B=0.40$



RG $\beta=2.6, 16^3 \times 32 \times 16, A=1.28, B=0.40, p=1$



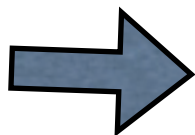
Similarly

$$D_k \simeq 0$$

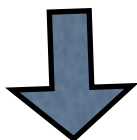
$$G_{NN}^Q(\vec{p}, t) = -|Z_+|^2 e^{-E_+ t} \gamma_5 D_P$$

$$|\vec{p}| = 0$$

Effective mass

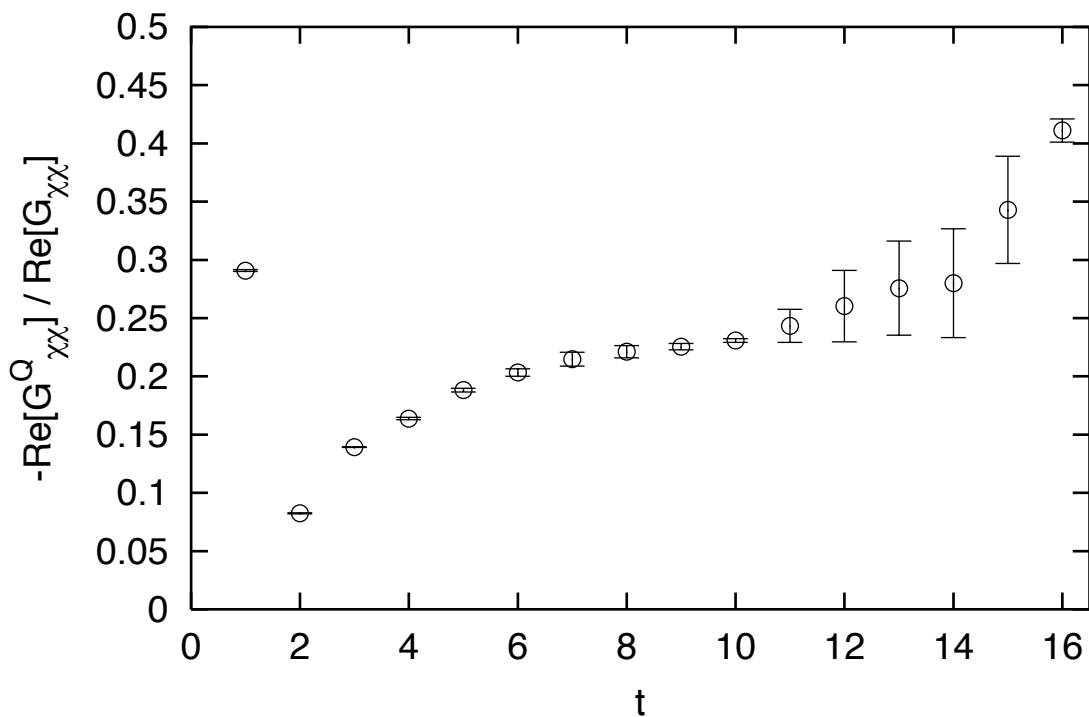
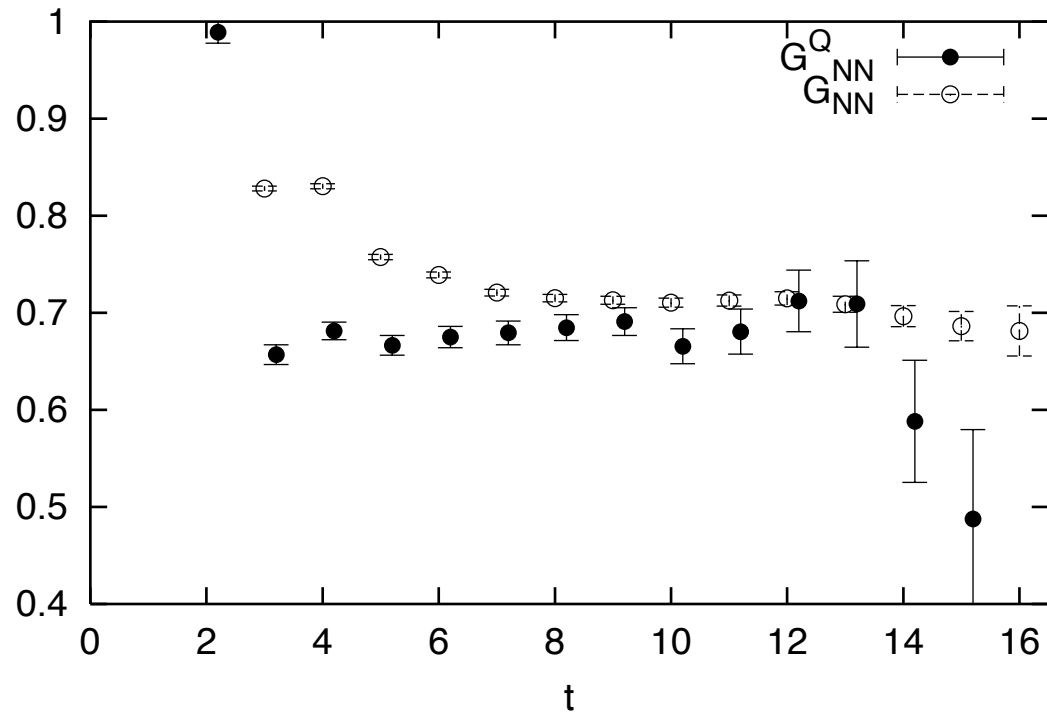


Ratio



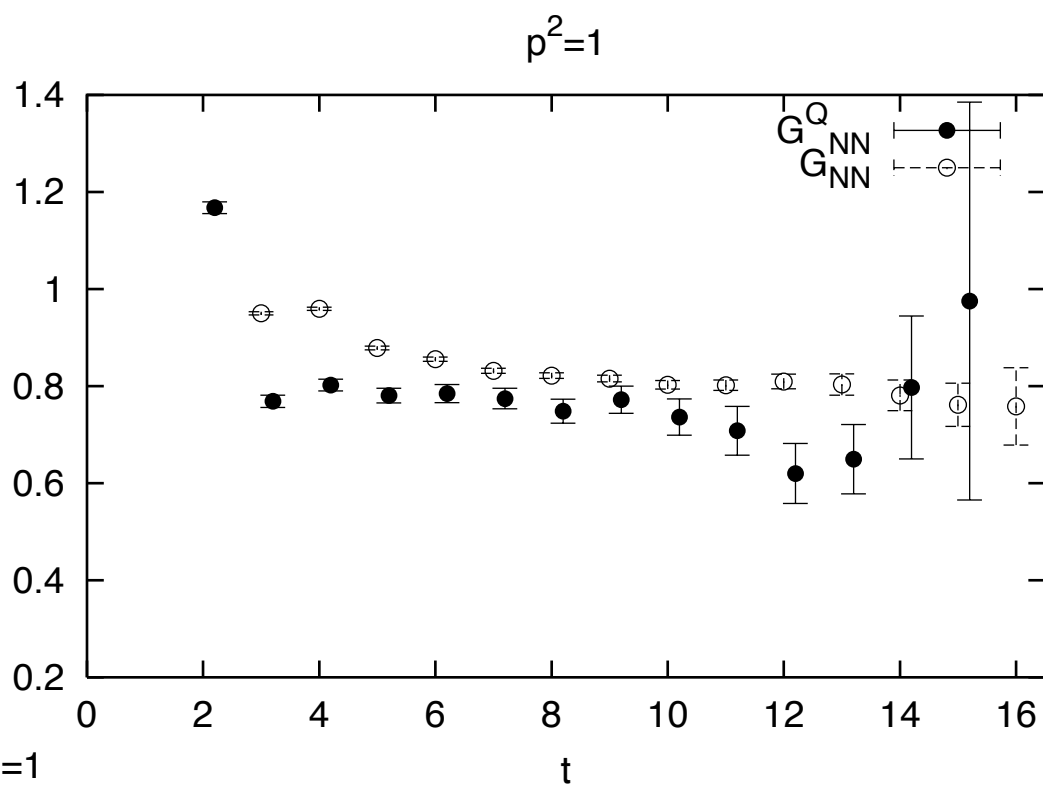
RG $\beta=2.6, 16^3 \times 32 \times 16, A=1.28, B=0.40$

$p^2=0$

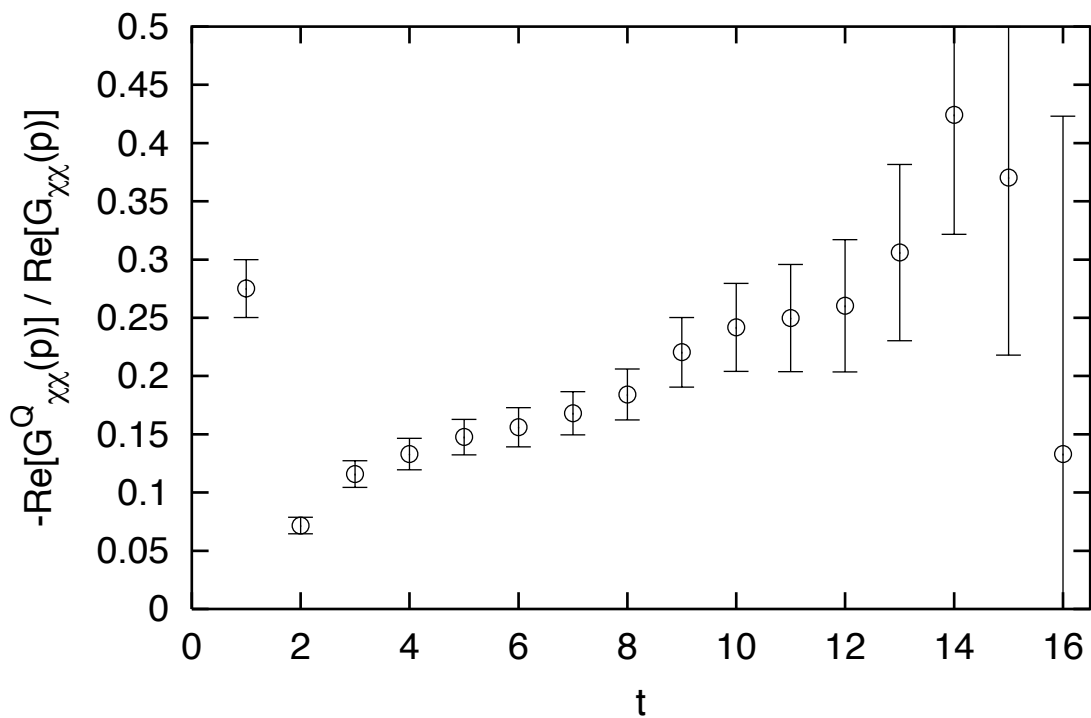


Signal for $D_P \neq 0$!

$$|\vec{p}| = \frac{2\pi}{L}$$



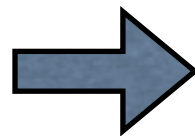
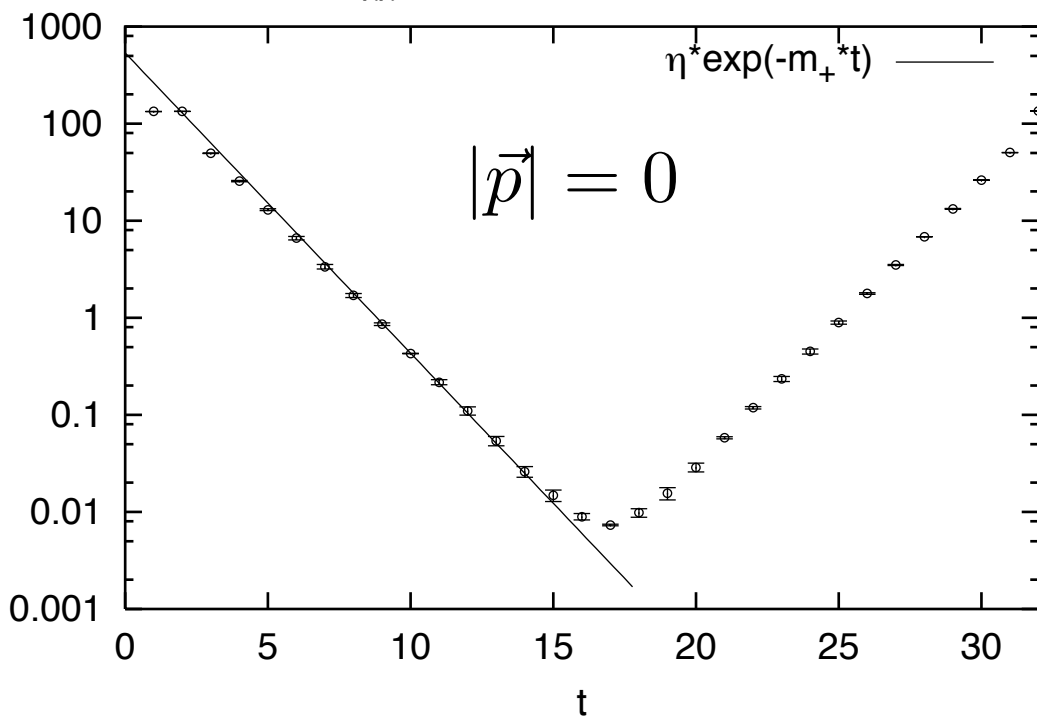
RG $\beta=2.6, 16^3 \times 32 \times 16, A=1.28, B=0.40, p=1$



$$G_{NN}^Q(\vec{p}, t) = -|Z_+|^2 e^{-E_+ t} \gamma_5 D_P$$

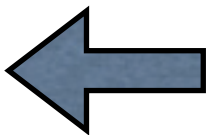
Fixed E_+

$G_{\chi\chi}^Q$ smear, fit range [9:14], $p=0$

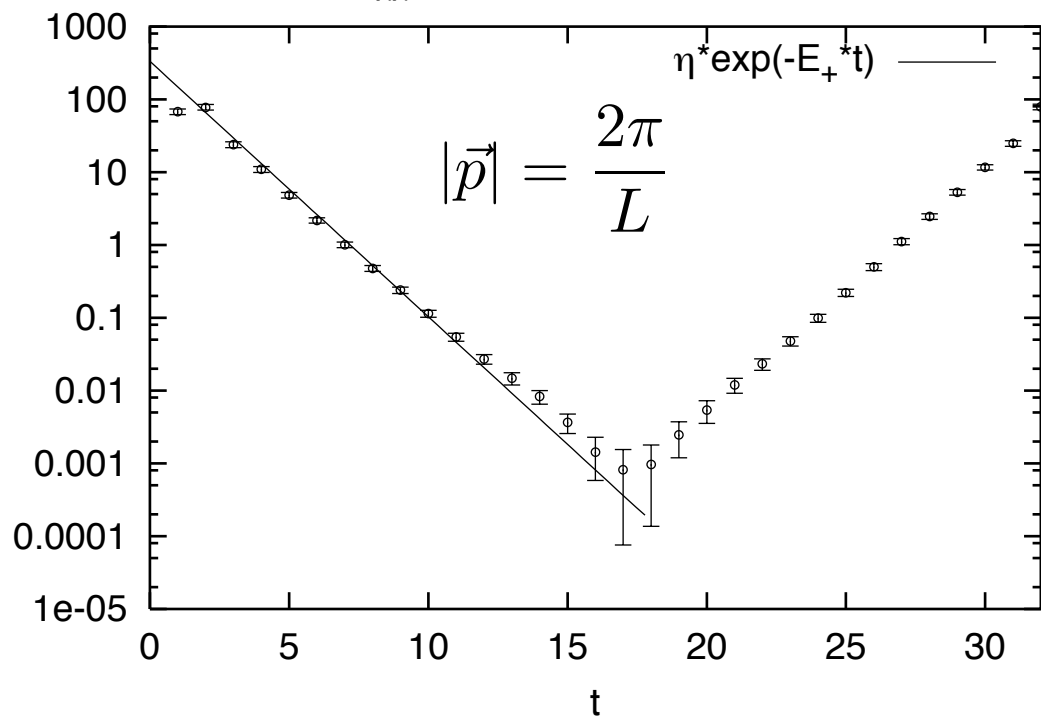


$$D_P = 0.248(17)$$

$$D_P = 0.216(18)$$



$G_{\chi\chi}^Q$ smear, fit range [9:14], $p=1$



Correlation between Q and quark is very important to get signals

$$\text{tr}[G_{NN}^Q(\vec{p}, t)\gamma_5] = -D_P |Z_+|^2 e^{-E_+ t} \Leftrightarrow \text{tr}[G_{NN}(\vec{p}, t)\gamma_5] = 0$$



Same E_+



$$\text{tr}[G_{NN}(\vec{p}, t)] = |Z_+|^2 e^{-E_+ t}$$

use of DWF

enough sampling of Q

3-pt function and form factor

$$\langle N_+^\theta(P') | J_\mu(q) | N_+^\theta(P) \rangle = \bar{u}_+^\theta(P') W_\mu^{\theta++}(q) u_+^\theta(P) \quad \text{Target}$$

$$W_\mu^{\theta++}(q) = W_\mu^{0++}(q) + i\theta W_\mu^{1++}(q) + O(\theta^2)$$

$$W_\mu^{0++}(q) = F_1(q^2) \gamma_\mu + F_2(q^2) \frac{q_\nu \sigma_{\mu\nu}}{2m_+}$$

$$W_\mu^{1++}(q) = F_3(q^2) \gamma_5 \frac{q_\nu \sigma_{\mu\nu}}{2m_+} + G_2(q^2) (\gamma_\mu q^2 - q_\mu \gamma \cdot q)$$

Lattice calculation

$$\langle \theta | \chi_+(\vec{p}', 0) J_\mu(q, \tau) \bar{\chi}_+(\vec{p}, t) | \theta \rangle = |Z_+^\theta|^2 e^{-E_+^{\theta'} \tau} e^{-E_+^\theta (t-\tau)}$$

$$\times \sum_{s, s'} u_+^\theta(P') \bar{u}_+^\theta(P') W_\mu^{\theta++}(q) u_+^\theta(P) \bar{u}_+^\theta(P) + \dots$$

By expanding in terms of θ in both sides

$$\langle 0|\chi(0)J_\mu(\tau)\bar{\chi}(t)|0\rangle = |Z_+|^2 e^{-E'_+\tau} e^{-E_+(t-\tau)} \\ \times \frac{-i\gamma \cdot p' + m_+}{2E'_+} W_\mu^{0+++}(q) \frac{-i\gamma \cdot p + m_+}{2E_+}$$

$$\langle 0|\chi(0)J_\mu(\tau)\bar{\chi}(t)Q|0\rangle = |Z_+|^2 e^{-E'_+\tau} e^{-E_+(t-\tau)} \\ \times \left[\frac{-i\gamma \cdot p' + m_+}{2E'_+} W_\mu^{1+++}(q) \frac{-i\gamma \cdot p + m_+}{2E_+} \right. \\ - \frac{-i\gamma \cdot p' + m_+}{2E'_+} W_\mu^{0+++}(q) D_P(\gamma_5 X_{-+} + X_{+-} \gamma_5) \\ \left. - D_P(\gamma_5 X_{-+} + X_{+-} \gamma_5)' W_\mu^{0+++}(q) \frac{-i\gamma \cdot p + m_+}{2E_+} \right]$$

$$X_{\pm\mp} = \sum_s u_\pm(P)\bar{u}_\mp(P)$$

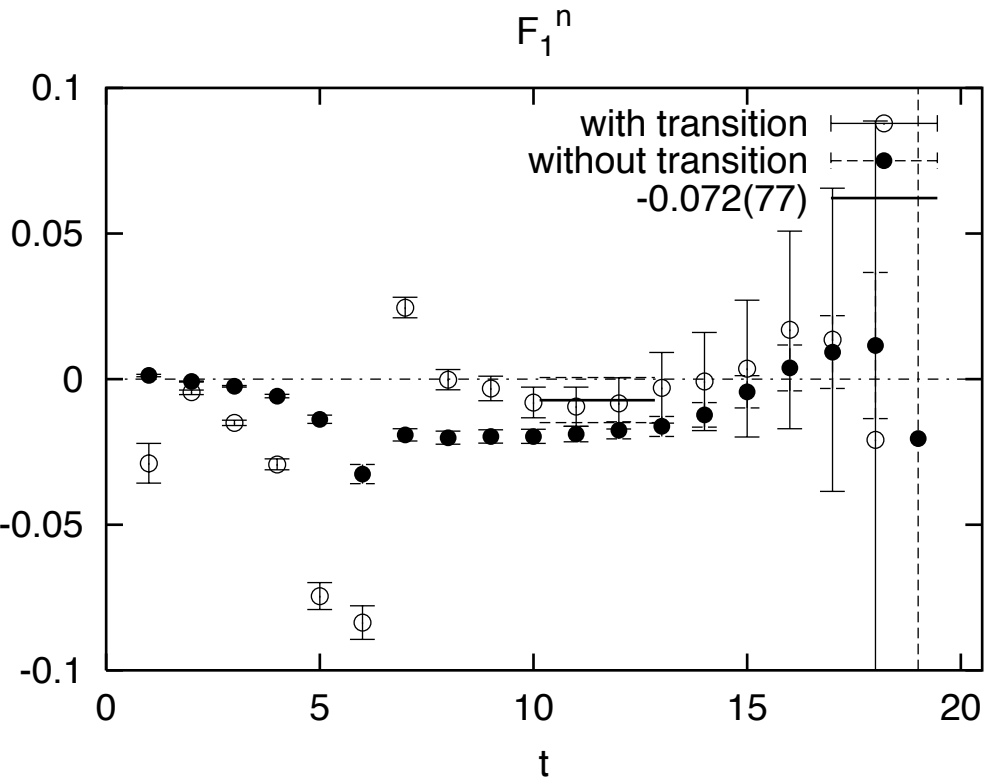
We have to subtract the effect of mixing due to
from

$$\langle 0 | \chi(0) J_\mu(\tau) \bar{\chi}(t) Q | 0 \rangle \quad \text{to obtain} \quad W_\mu^{1^{++}}(q)$$

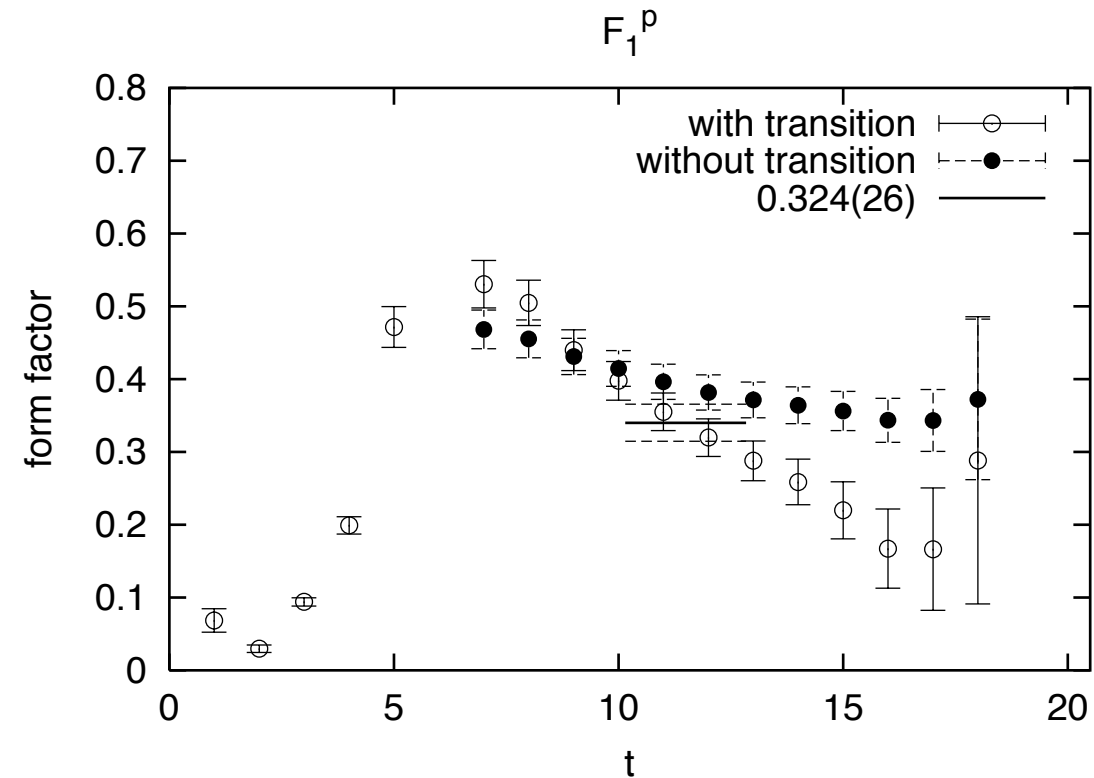
Work in progress (Shintani)

$$W_\mu^{0^{++}}(q)$$

$$f_1(q^2)$$



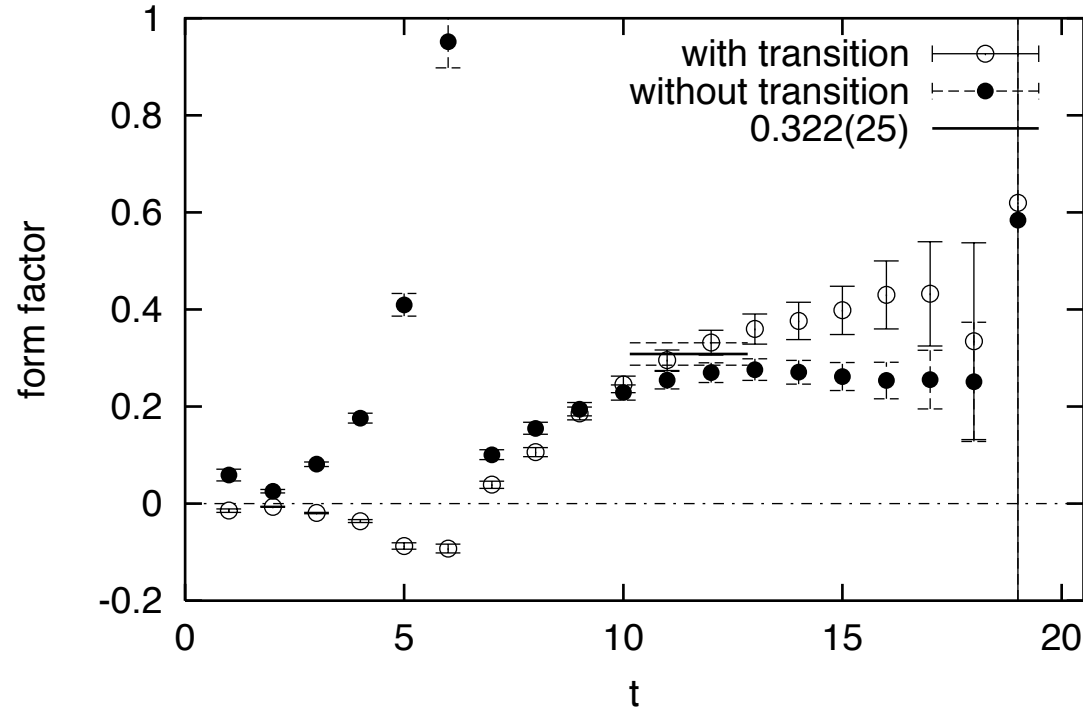
Neutron



Proton

$$W_{\mu}^{0^{++}}(q)$$

$$f_2(q^2)$$

 F_2^p 

Neutron

Proton

 F_2^n 