

Actions for dynamical fermion simulations: are we ready to go?

Karl Jansen



- Conceptual questions
- Algorithmic questions
- questions in (chiral) perturbation theory
- Practical questions
- some results
- Con(bl)(f)usions





Actions

- **Wilson fermions**
 - non-perturbatively improved fermion action
 - various gauge actions (Plaquette, Symanzik, RG-improved)
 - **Staggered fermions**
 - improved fermion action (Asqtad)
 - various gauge actions
 - **Domain wall and overlap fermions**
 - RG improved gauge actions
 - fermion actions with eigenvalue projection
 - **Designer actions**
 - FLIC, Hypercube (various versions) + many more
- have to agree in continuum limit: provide valuable cross check
- don't waste resources

rigorous actions

- reflection positivity, Osterwalder-Schrader positivity, positive transfer matrix \Rightarrow reconstruction theorem
 - Wilson action
Lüscher, Commun.Math.Phys.54:283,1977; for $r = 1, \kappa < 1/6$
(– tmQCD)
 - (naive) staggered fermions:
Sharatchandra, Thun, Weisz, Nucl.Phys.B192:205,1981; Smit, Nucl.Phys.Proc.Suppl.20:542-545,1991 ; Palumbi, hep-lat/0208005
positive transfer matrix for 2 lattice spacings

not rigorous but local actions

- **no proof of reflection positivity or construction of positive transfer matrix**
- **ultra local actions**
 - Designer actions, I will take as example FLIC*
 - Symanzik improved actions
 - truncated perfect action
- **exponentially localized**
 - overlap
 - domain wall
 - perfect action

* *Fat Link Irrelevant Clover fermions*

$$D_{\text{FLIC}} = \frac{1}{u_0} \nabla_\mu \gamma_\mu + \frac{1}{2u_0^{(fl)}} \left(\Delta^{(fl)} - \frac{1}{2u_0^{3(fl)}} \sigma \cdot \mathcal{F}^{(fl)} \right)$$

Non-local actions?

candidate: taking square root of staggered fermion matrix

test following [Hernández, Lüscher, K.J.](#)

Source point

$$\eta_\alpha(x) = 1 \text{ for } x = 1, \alpha = 1 \quad \eta_\alpha(x) = 0 \text{ else}$$

compute for some operator $A^\dagger A$

$$\Psi(x) = \sqrt{A^\dagger A} \eta(x)$$

test whether couplings of the operator decay exponential

$$f(r) = \max \{ \|\Psi(x)\|; \|x - y\|_{\text{taxi}} = r \}$$

test for fixed value of lattice spacing a ; positive outcome:

$$f(r) = e^{-r/r_{\text{local}}}$$

locality in continuum limit?

possibility I

$$r_{\text{local}} \cdot m_{\pi} = \text{constant}; \text{ for } a \rightarrow 0, m_{\pi} \text{ fixed}$$

\Rightarrow obtain a continuum theory with $r_{\text{local}} \propto \xi_{\pi}$ non-local theory on the scale of pion Compton wave length \Rightarrow unacceptable

possibility II

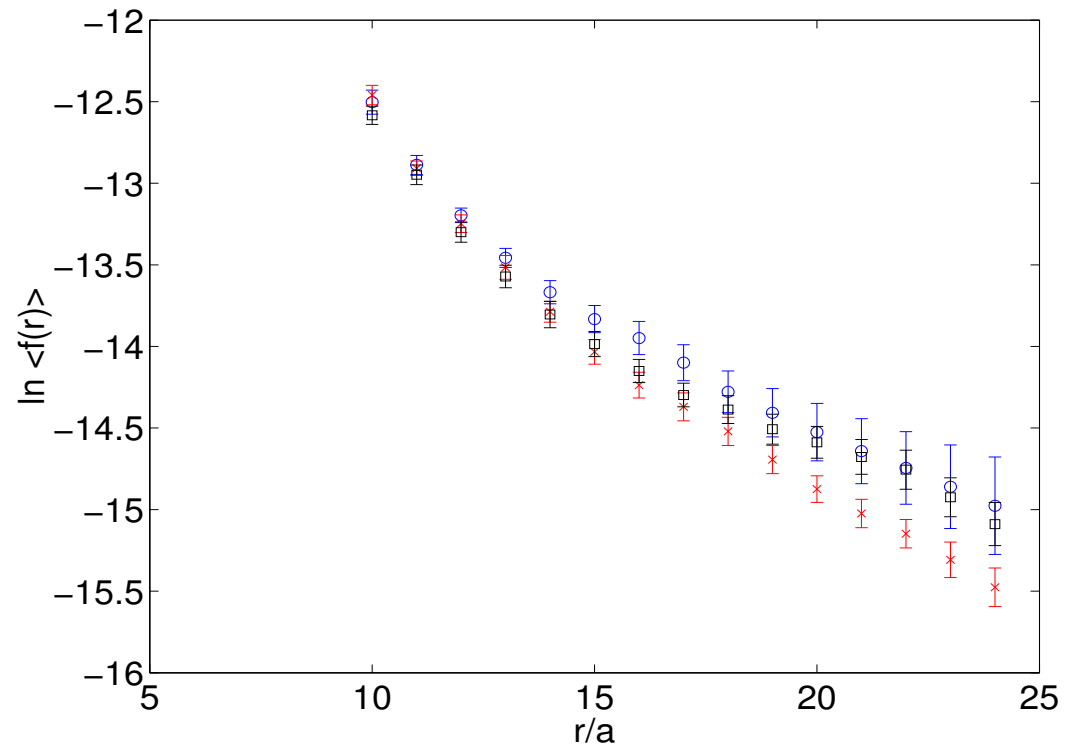
$$r_{\text{local}} \cdot m_{\pi} \rightarrow 0 \text{ for } a \rightarrow 0, m_{\pi} \text{ fixed}$$

$\Rightarrow r_{\text{local}}/a = \text{const}$ obtain a point local continuum theory

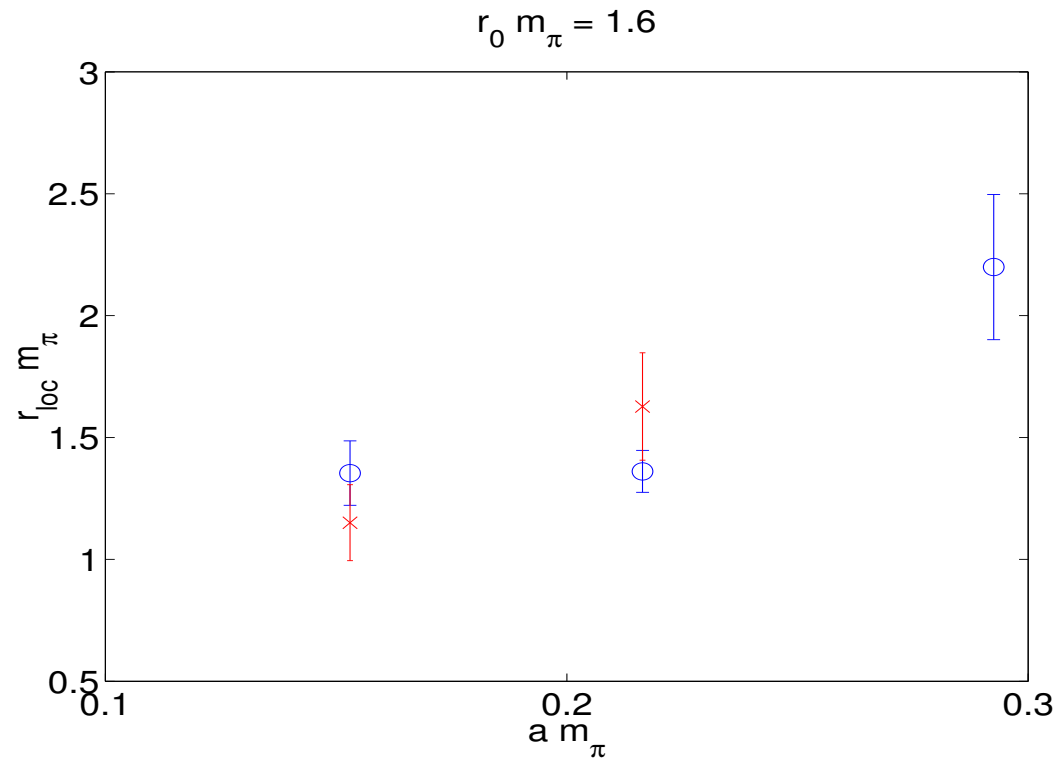
A first look

use (F. Knechtli, K.J.): $A =$ Wilson operator, $\sqrt{A^\dagger A} = P_{n,\epsilon}(A^\dagger A)$

fix $r_0 \cdot m_\pi = 1.6$, various $\beta = 6, 6.2, 6.45$



$$r_{\text{local}} \cdot m_{\pi}$$



red crosses: take r_{local} at $\beta = 6.0$ and scale it according to change of lattice spacing

My personal wishlist I

precise check for localization of staggered fermions

work in progress, Della Morte, Knechtli, K.J.

C, P & T

A warning from M. Creutz

spontaneous CP violation might be possible for $m_u \rightarrow 0$ tuning it negative

← miss this possibility when taking square roots?

← miss interesting part of physics?

A warning from Klinkhamer and J. Schilling

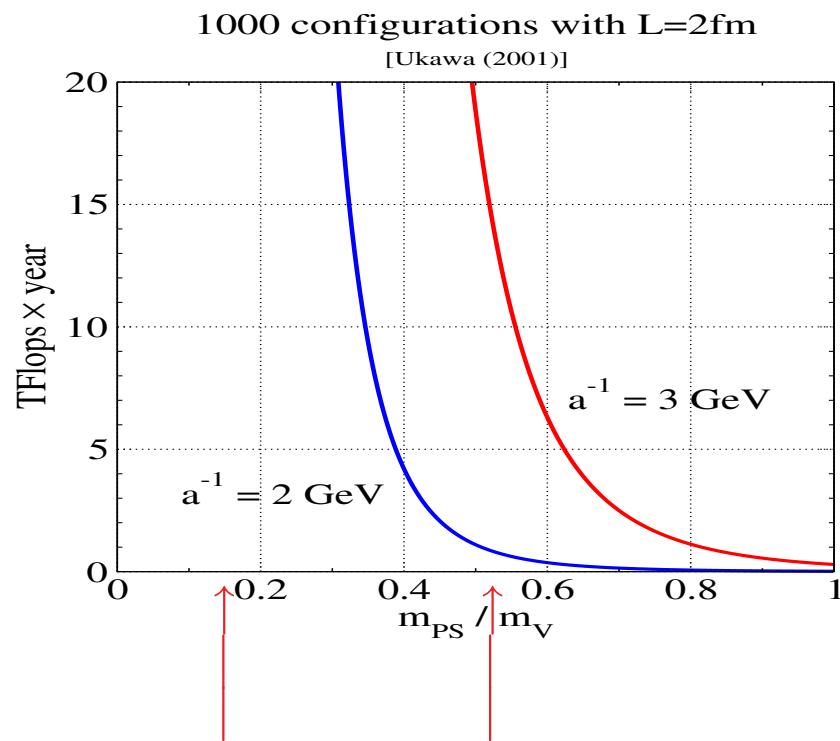
for a special class of gauge fields ($U_4(\mathbf{x}, x_4) = 1, U_m(\mathbf{x}, x_4) = U_m(\mathbf{x})$) *chiral gauge theories* from overlap fermions not CPT invariant

← violation of reflection positivity? Consequences?

see also Fujikawa, Ishibashi, Suzuki

Costs of dynamical fermions simulations

see panel discussion in Lattice2001, Berlin, 2001



$$\text{formula } C \propto \left(\frac{m_{\pi}}{m_{\rho}} \right)^{-z_{\pi}} (L)^{z_L} (a)^{-z_a}$$

$$z_{\pi} = 6$$

$$z_L = 5$$

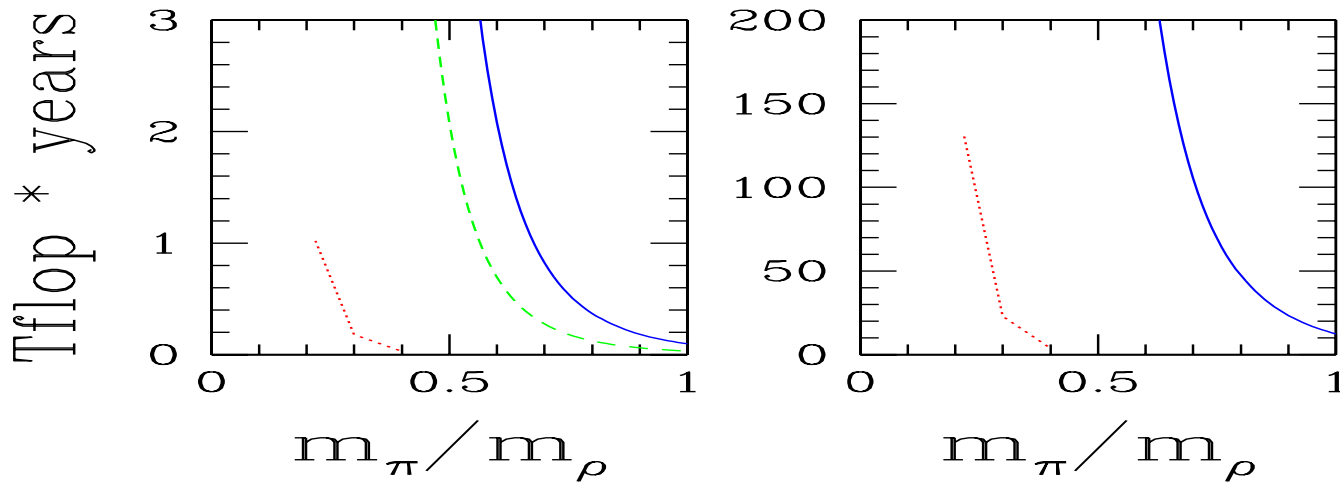
$$z_a = 7$$

physical
point

contact to
 χPT (?)

\Rightarrow use chiral perturbation theory (χPT) to extrapolate to physical point

Wilson versus staggered at fixed box length $L = 2.5$ fm



$a = 0.09$ fm
staggered: measured

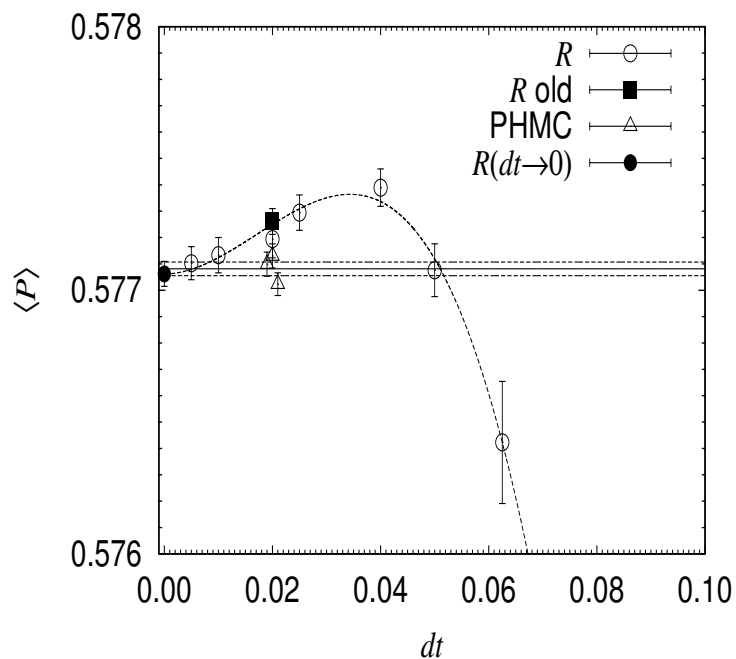
$a = 0.045$ fm
staggered: extrapolated

full line: Wilson; dashed line: staggered; dashed line: Wilson/3

MILC data, thanks to S. Gottlieb

Exact vs. inexact: why inexact?

Exact algorithm PHMC algorithm for $N_f = 3$ Aoki et.al. (JLQCD) hep-lat/0208058
(see also T. Kennedy)



- extrapolation to $\delta\tau = 0$ difficult
- treat $(A^\dagger A)^{1/n}$ by polynomial
- noisy Metropolis step or correction factor inversion of $(A^\dagger A)^{1/n}$ by Lancsoz method
- cost of exact algorithm \approx in-exact

My personal wishlist II

Use and test exact odd flavour algorithms

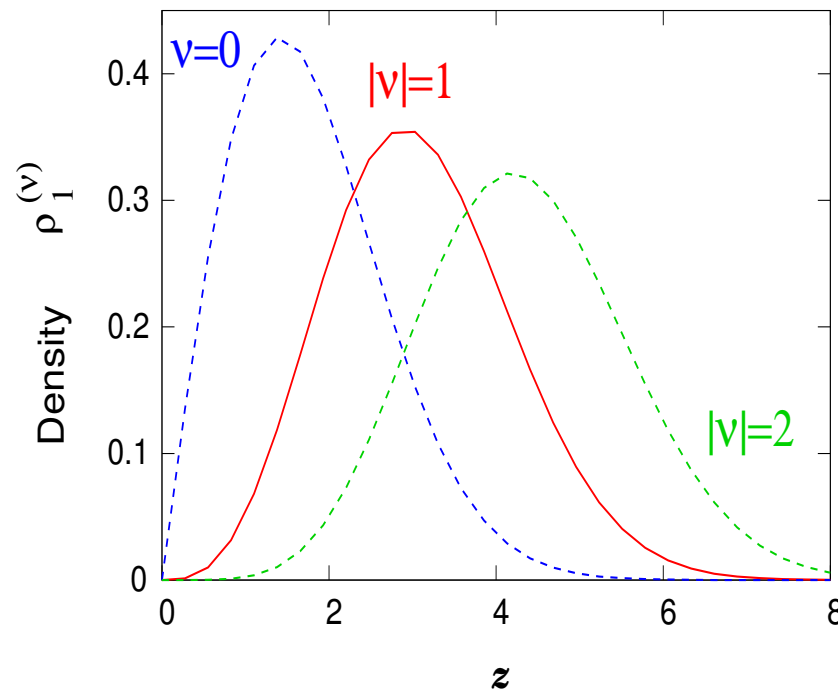
fair comparison of exact algorithms, continuum approach

How to simulate a designer action

→ complicated interactions, fattening

- first way a la Hasenbusch; Hasenfratz and Knechtli + many others
 - i) $U \rightarrow U'$ according to gauge field action
 - ii) $\det(A'^{\dagger}A')/\det(A^{\dagger}A) \rightarrow$ accept/reject; correction factor
 - iii) needs smearing/fattening
improvements: break up of determinant, ultraviolet filtering, ...
- second way a la W. Kamleh
re-unitarization through $X/\sqrt{X^{\dagger}X}$
expand $1/\sqrt{X^{\dagger}X}$
use chain rule to go from fattened link $U^{(n)}$ to original link $U^{(0)}$

A problem of principle:
the eigenvalue distribution from Random matrix Theory



- \Rightarrow small eigenvalues have to appear, checks in quenched simulations
Bietenholz, Shcheredin, K.J., QCDSF, Weisz et.al.
- \Rightarrow can lead to large statistical fluctuations or difficulties in the simulations when approaching the physical point

Perturbation theory
(review Capitani, hep-lat/0211036)

Analysis for Wilson fermions Bochicchio, Maiani, Martinelli, Rossi, Testa

Analysis for staggered Sharatchandra, Thun, Weisz; Goltermann, Smit, Vink

Designer actions

more links of course more complicated but doable

fattening/smearing/blocking \rightarrow

$$\int \frac{d^4q}{(2\pi)^4} I(q) \rightarrow \int \frac{d^4q}{(2\pi)^4} \left(1 - \frac{c}{6} \hat{q}^2\right)^{2N} I(q)$$

$c < 1$ smearing coefficient, N number of smearing steps

tadpole contribution substantially reduced:

$$12.23g_0^2/(16\pi^2)C_F \rightarrow 0.35g_0^2/(16\pi^2)C_F$$

Reisz Power Counting Theorem

(Reisz, Lüscher)

statement is that the lattice integral

$$I = \int_{-\pi/a}^{\pi/a} \frac{d^4 k}{(2\pi)^4} \frac{V(k, m, a)}{C(k, m, a)}$$

exists in the continuum limit, if (among others) the condition

$$|C(l, m, a)| \geq A(\hat{l}^2 + m^2)$$

is fulfilled for a small enough and some positive value of A

$$\text{Wilson}(r = 1) \quad C = (1 + am)\hat{p}^2 + m^2 + \frac{1}{2}a^2 \sum_{\mu < \nu} \hat{p}_\mu^2 \hat{p}_\nu^2$$

$$\text{Staggered} \quad C = \sum_{\mu} \sin^2 k_{\mu} + m^2 = \sum_{\mu} \hat{k}^2 - \frac{a^2}{4} \sum_{\mu} \hat{k}^4 + m^2$$

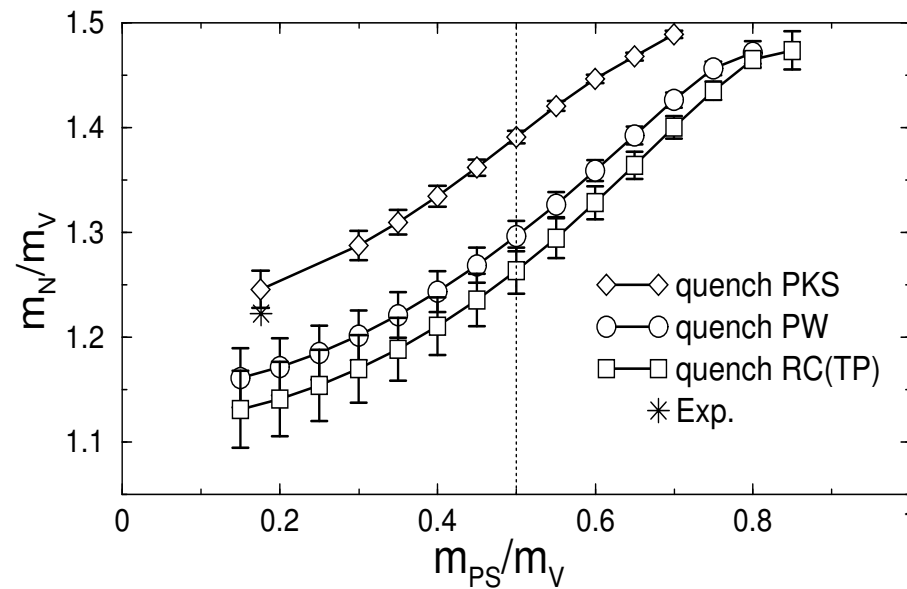
My personal wishlist III

construct a “Reisz theorem” for staggered fermions

Inconsistencies?

S. Aoki, hep-lat/0011074, Lattice2000 review

- PKS: plaquette action, staggered fermions
- PW: plaquette action, Wilson fermions
- RC(TP): RG gauge action, tadpole improved Wilson

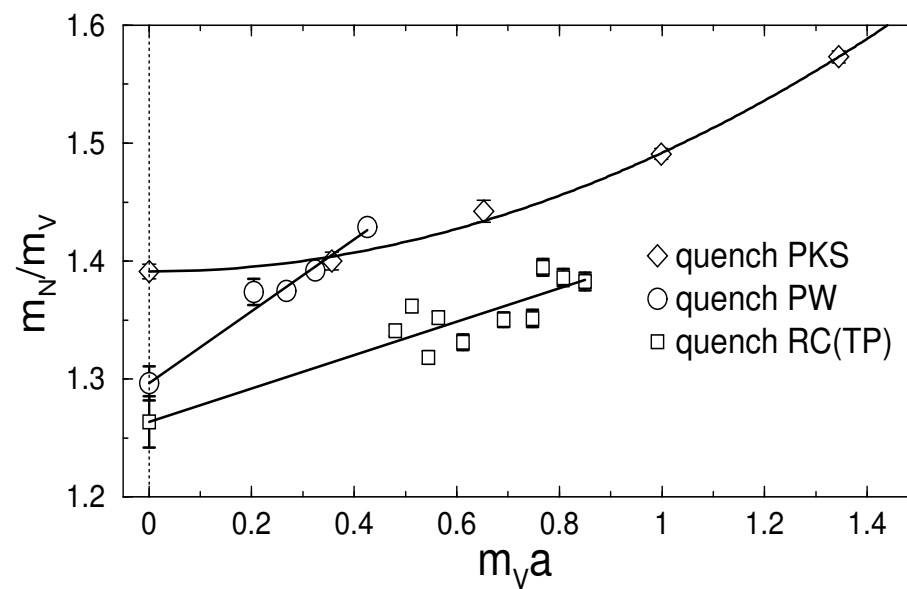


→ different continuum results even at large masses!

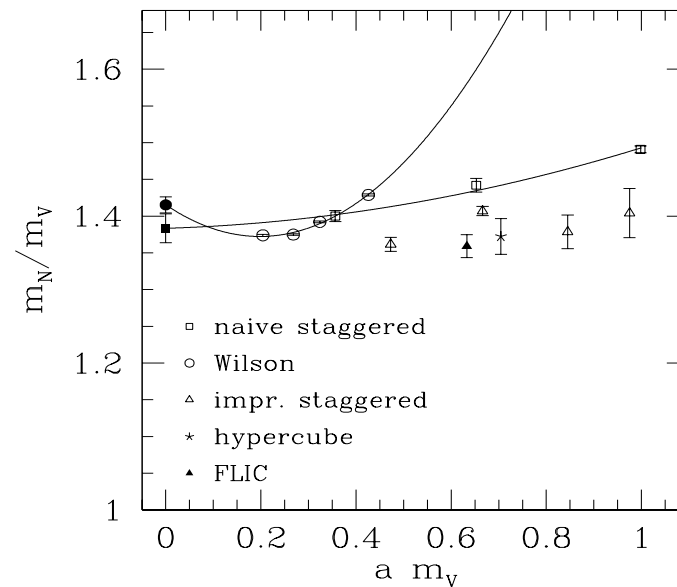
The continuum extrapolation

S. Aoki, hep-lat/0011074, Lattice review

$$m_\pi/m_\rho = 0.5$$



Large lattice artefacts/alternative fits?



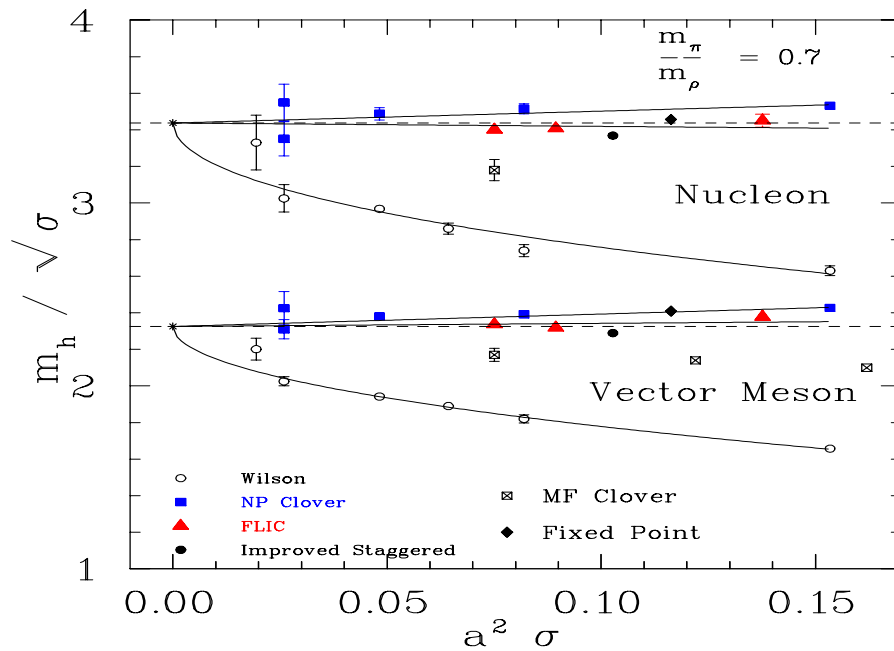
K.J. and J. Zanotti

fit may not not be the final one, but it is a possibility

My personal wishlist **IV**

precise scaling analysis for various fermion actions in the quenched approximation

A scaling plot

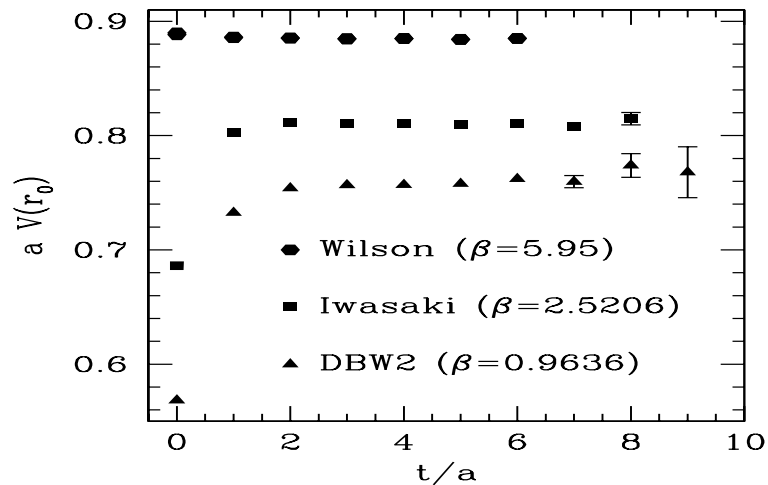


thanks to J. Zanotti

(talks by A. & P. Hasenfratz for scaling tests of Hyp, Asqtad, CI and TP)

Problems in practical dynamical simulations:
Gauge actions

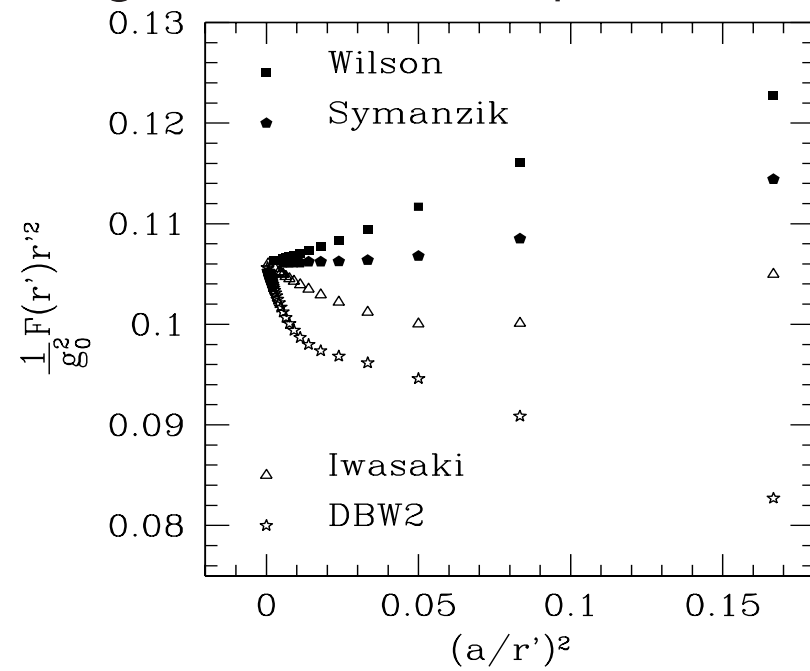
I: RG action \rightarrow not reflection positive
 \Rightarrow complex energies



Necco, Sommer

free field analysis: $t \gg t_{\min} = \begin{cases} 0.5 & \text{Symanzik} \\ 0.9 & \text{Iwasaki} \\ 1.7 & \text{DBW2} \end{cases}$

II: large lattice artefacts possible

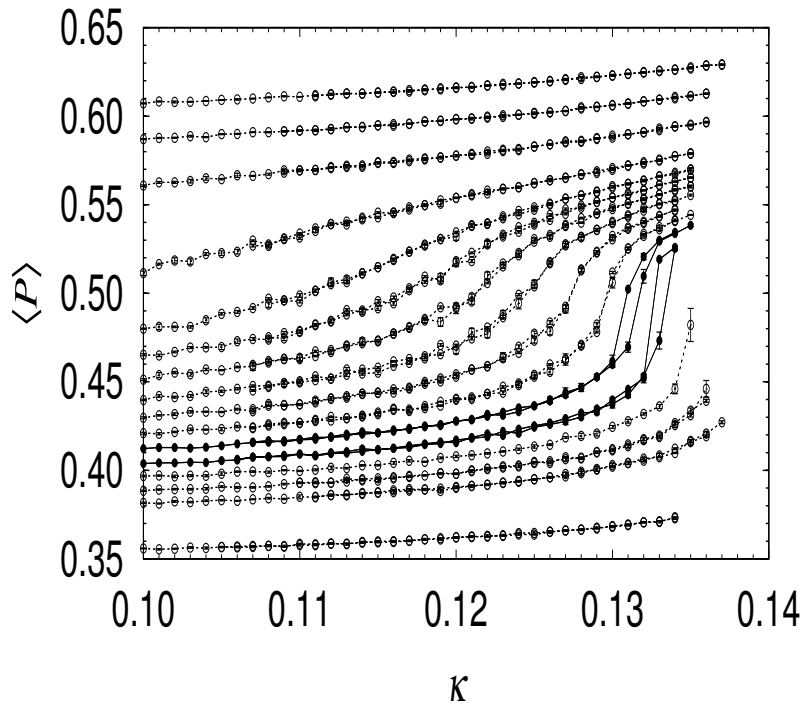


⇒ two action method?

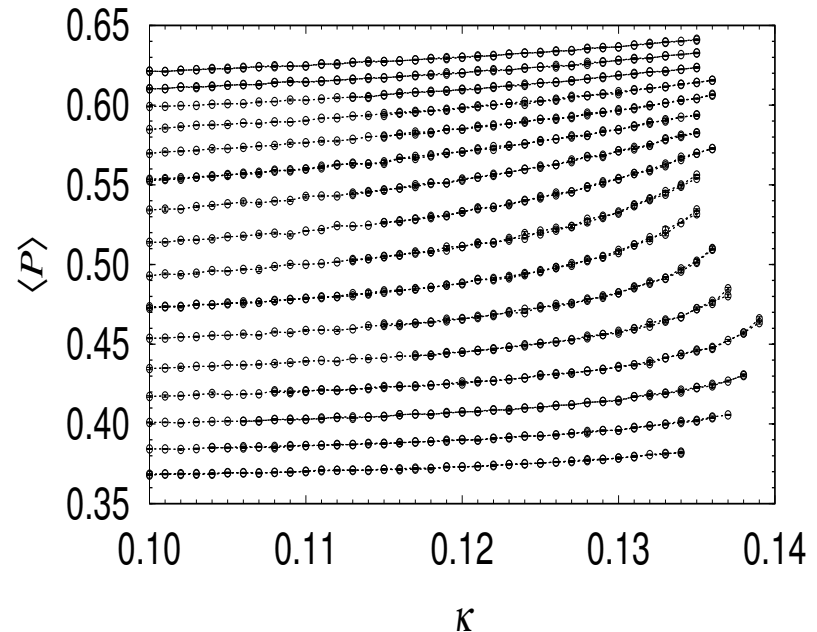
III: difficulty of sampling topological charge sectors

Problems in practical dynamical simulations: **Wilson**

Simulations with $N_f = 3$ improved fermions **CP-PACS**

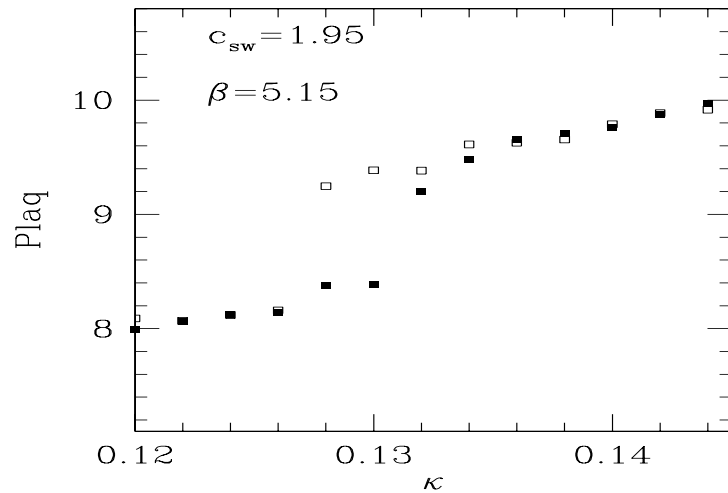


Wilson gauge



RG improved action

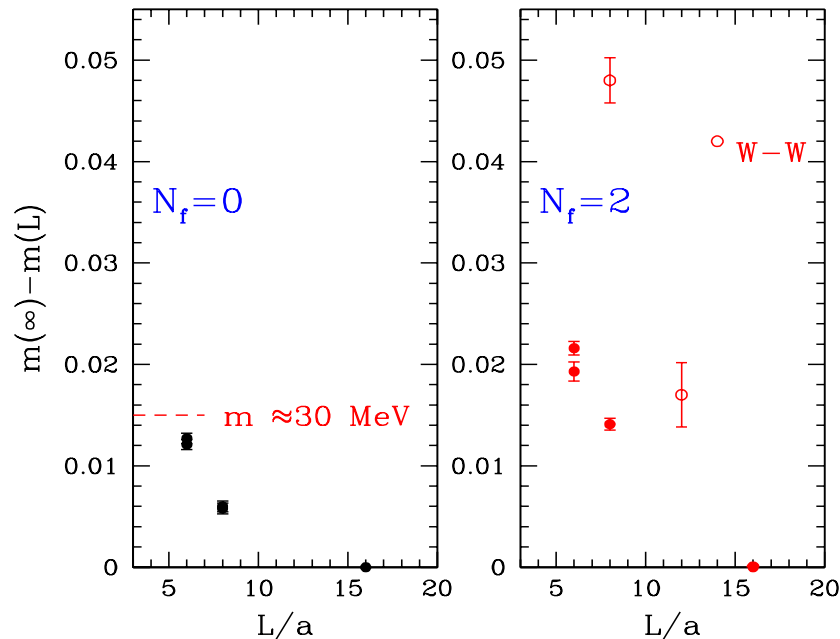
Simulations with $N_f = 2$ Wilson gauge non-perturbatively improved fermions (K.J.)



- hysteresis effect
- effects almost independent from values $1 < c_{sw} < 2$
- small lattice simulations

(unexpected) large lattice artefacts in quark mass

Wilson gauge, non-perturbatively improved Wilson fermions



Non-perturbatively improved Wilson
 $m(\infty) = m(16)$, W-W: Wilson action

My personal wishlist VI

Study and understand $T = 0$ phase diagram of QCD

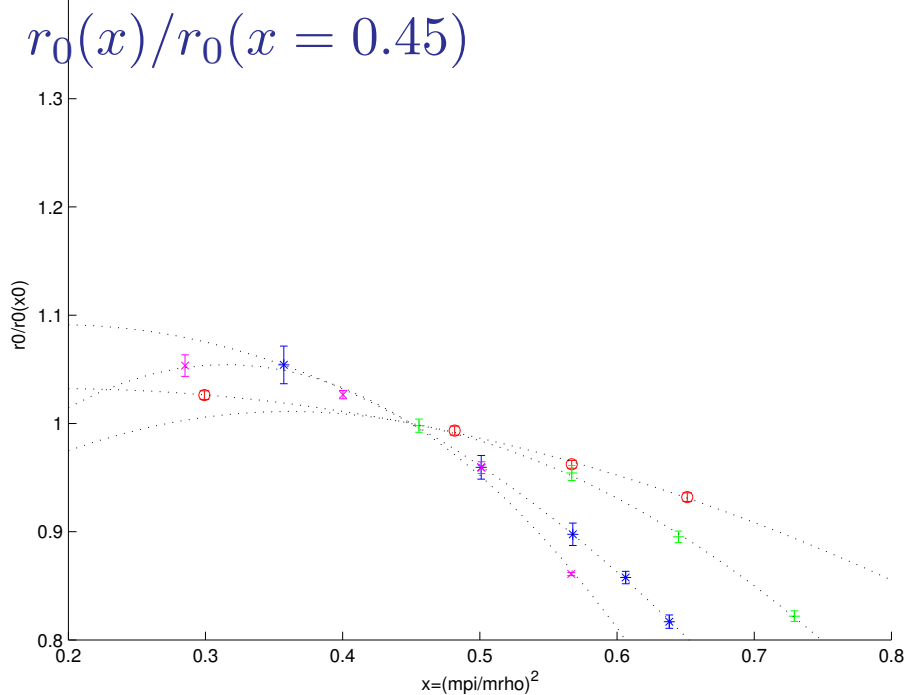
Investigate different gauge actions

understand nature of phase transition

W-KS: \times W-W: $+$ I-SW: \circ W-SWimpr: $*$

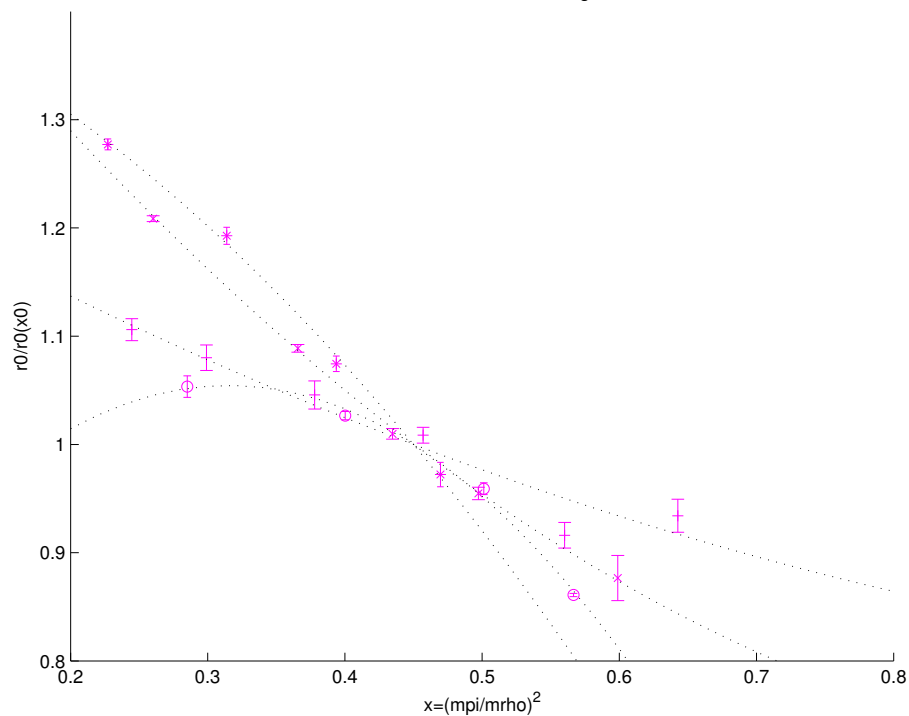
$R(y)$ for $y_{\text{ref}} = 0.45$, $a \approx 0.1$ fm

green +: SESAM, red o: CPPACS (o), blue *: JLQCD, KS-fermions: magenta x



KS-fermions various $a \geq 0.1$ fm

MILC, KS-fermions: a decreasing: $+ \times \circ$



$$x = (m_\pi/m_\rho)^2$$

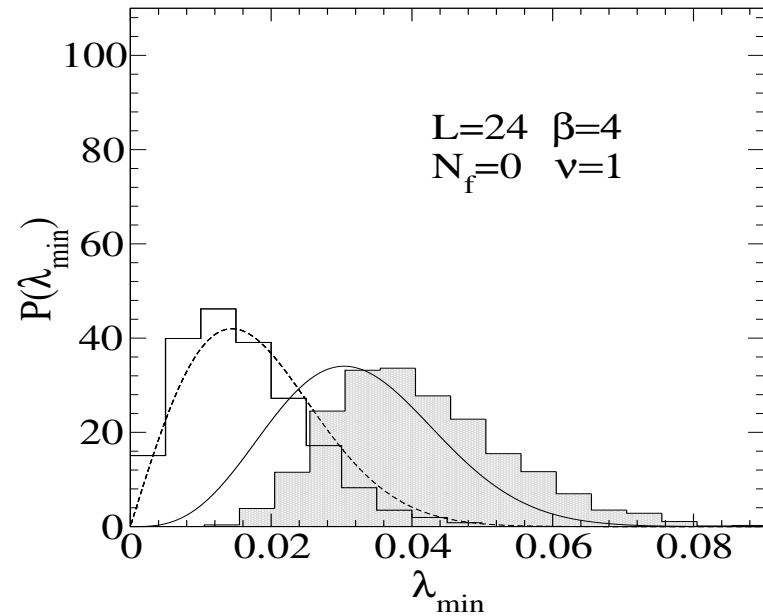
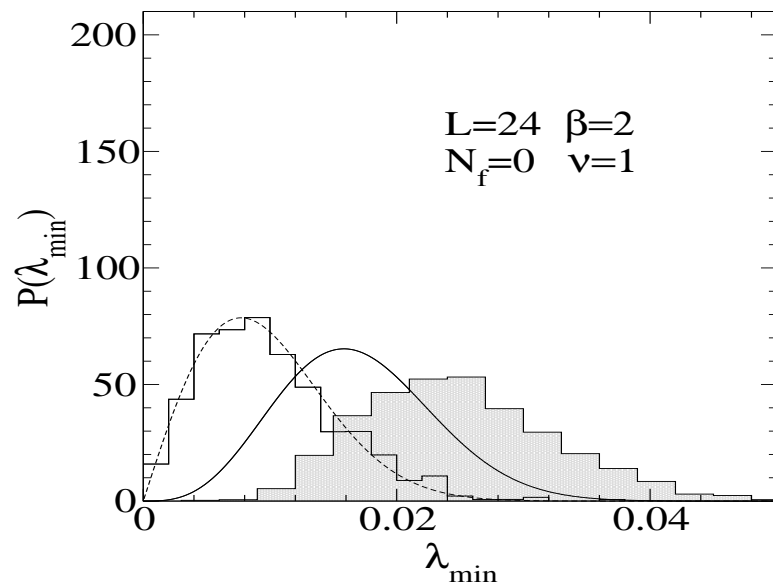
Large effects?

Large cutoff effects!

R. Sommer

Problems in practical dynamical simulations: Staggered

Eigenvalue distribution of staggered operator in comparison to Random matrix theory



Farchioni, Hip, Lang

see also: Damgaard, Heller, Niclasen, Rummukainen, Berg, Markum, Pullirsch, Wettig

does problem disappear for $a \ll 1$? How do improved actions behave?

Results (nevertheless) Wilson

- comparison to chiral perturbation theory
- finite size effects
- status of running quark mass
- towards $N_f = 3$
- meson spectrum

Chiral Perturbation theory

two strategies:

1.) *extrapolate to continuum limit and fit then to predictions of χPT*

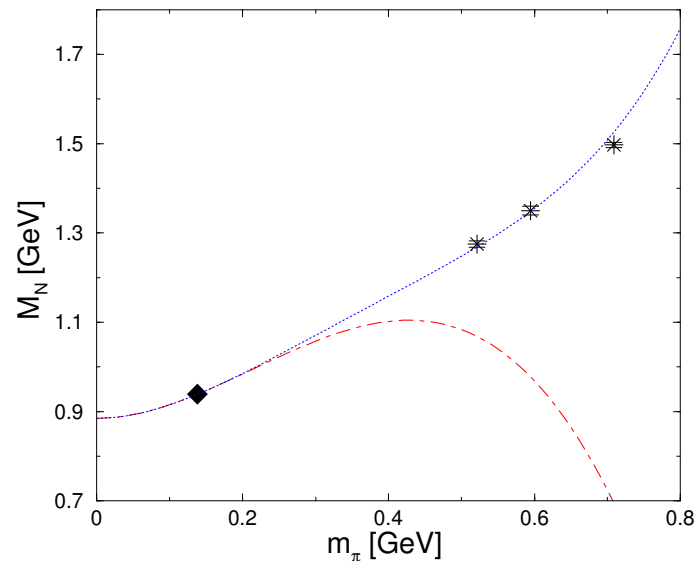
advantages/disadvantages

- chiral invariance ensured
- direct comparison possible
- computationally demanding

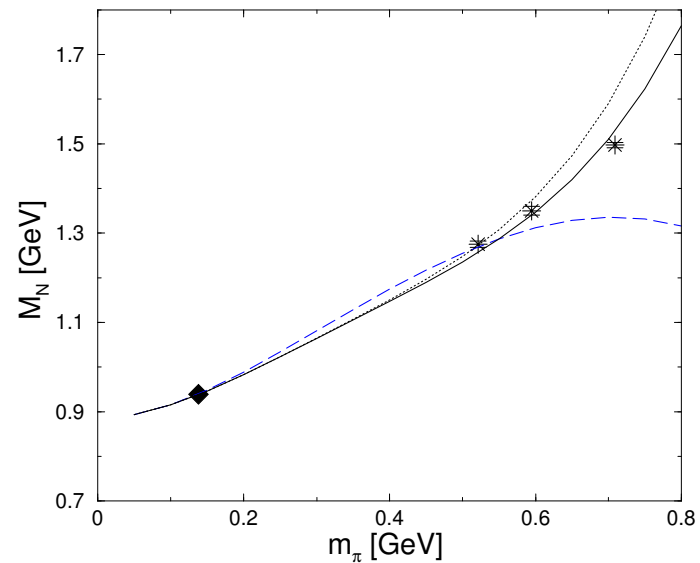
in practise (χPT practitioners): lattice data at non-vanishing lattice spacing are compared to continuum formulae

lattice data from chirally non-invariant lattice formulations

Example Bernard, Hemmert, Meissner

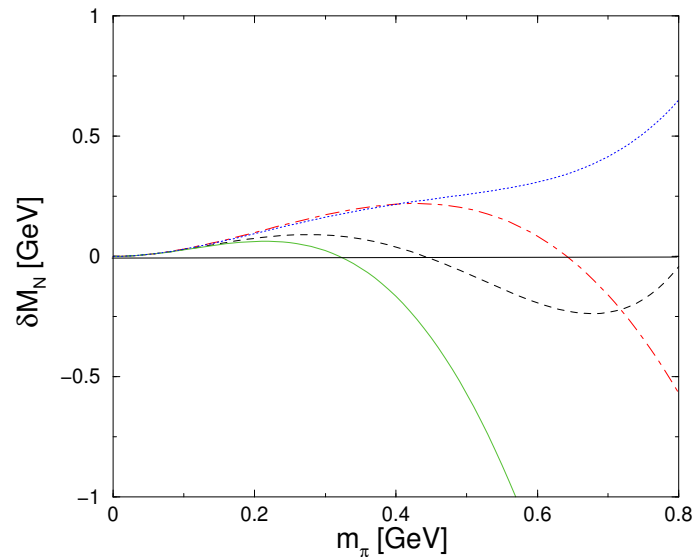


dimensional

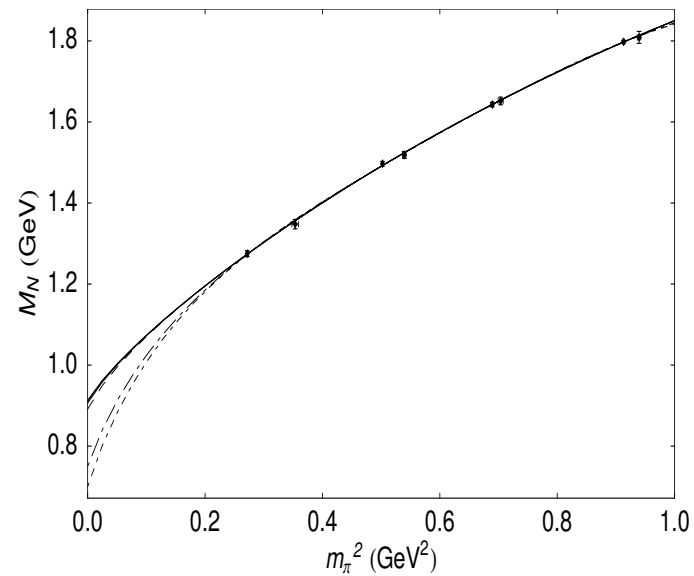


cut-off

- 4-loop results (quadratic in quark mass) keeping consistently chiral symmetry
- parameters of chiral lagrangian *fixed* at physical point
- “improvement term” added but only one!!
- data (CP-PACS) at smallest values of a available: close enough to the continuum? (see later)



size of 4-loop corrections



Leinweber, Thomas, Young (“pion cloud”)

BHM: *We stress again that applying the expressions to pion masses above 600 MeV is only done for illustrative purposes, for a realistic chiral extrapolation smaller pion masses are mandatory*

- it is possible to model lattice data
- clearly want, however, pure chiral perturbation theory

Chiral Perturbation theory

2.) *take discretization effects into account* Sharpe

Wilson fermions Baer, Rupak, Shoresh; Aoki

⇒ duplication of low energy constants

→ physical LEC $l_4, \dots, l_8 \leftrightarrow w_4, \dots, w_8$

Staggered fermions Aubin, Bernard, Goltermann, Lee, Sharpe + ...

start with Lee-Sharpe lagrangian

$$L = \frac{f^2}{8} \text{tr}(\partial_\mu \Sigma \partial_\mu \Sigma^\dagger) - \frac{1}{4} \mu m f^2 \text{tr}(\Sigma + \Sigma^\dagger) + \frac{2m_0^2}{3} (\Phi_I)^2 + a^2 V$$

$$\Sigma = \exp(i\phi/f), \quad \phi = \sum_a^{16} \phi_a T_a$$

Φ_I singlet field

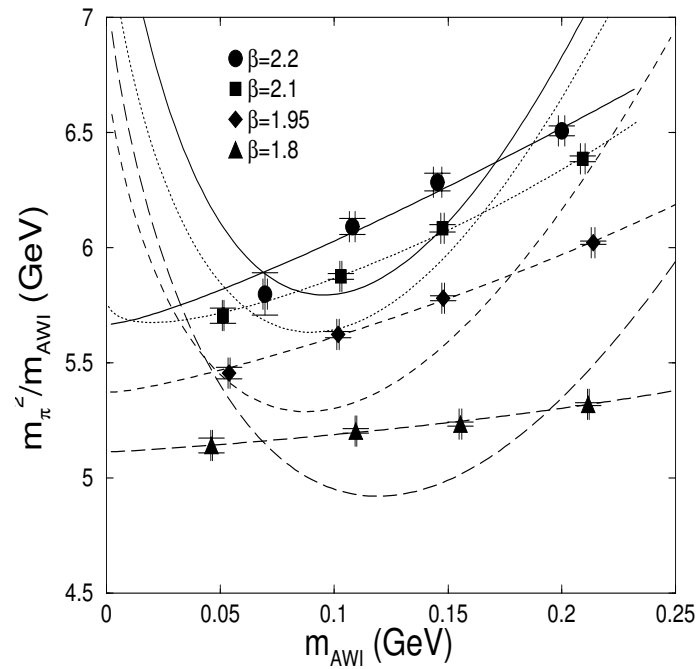
V staggered flavor breaking potential → six terms with coefficients C_1, \dots, C_6

Chiral Perturbation theory

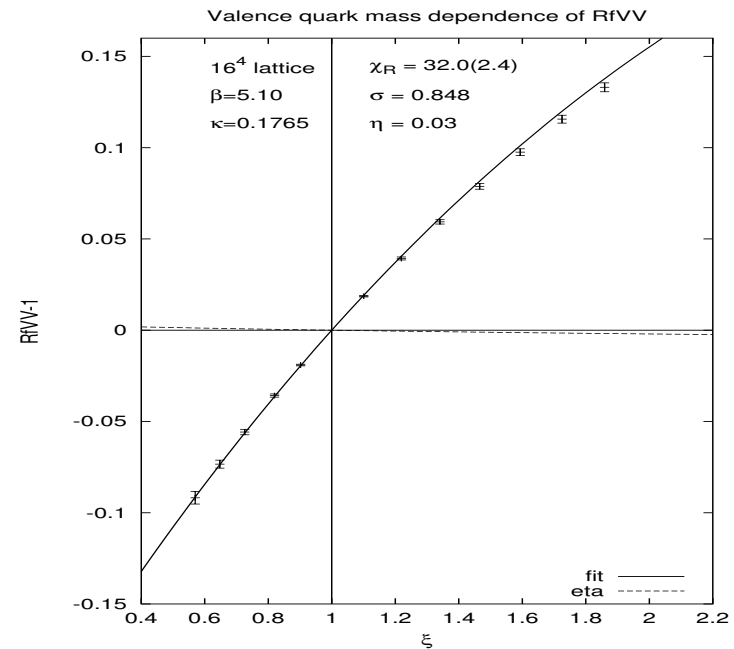
advantages/disadvantages

- lattice discretization effects can be (partly) absorbed
- allows for hybrid simulations such as improved fermions dynamical, overlap quenched
- (many) new parameters
- dependence of new fit parameters on g_0

Wilson examples



$$M_\pi^2/m_q \text{ (Aoki)}$$



double ratio (Farchioni, Gebert,
Montvay, Scholz, Scorzato)

$$Rf_{VV} = \frac{M_\pi^2(\text{sea})/m_q(\text{sea})}{M_\pi^2(\text{valence})/m_q(\text{valence})}$$

→ amazing cut-off cancelations in double ratios $a = 0.28\text{fm}$ (!)

Chiral Perturbation theory

remarks: it would be important to disentangle

- sea quark effects: plot only sea quark dependence
- a effects: de-double double ratios

a problem for universal LEC: assume

Aoki: shifting $\beta = 2.1 \rightarrow \beta = 2.2$, $\Lambda = 0.694(20) \rightarrow \Lambda = 0.128(88)$

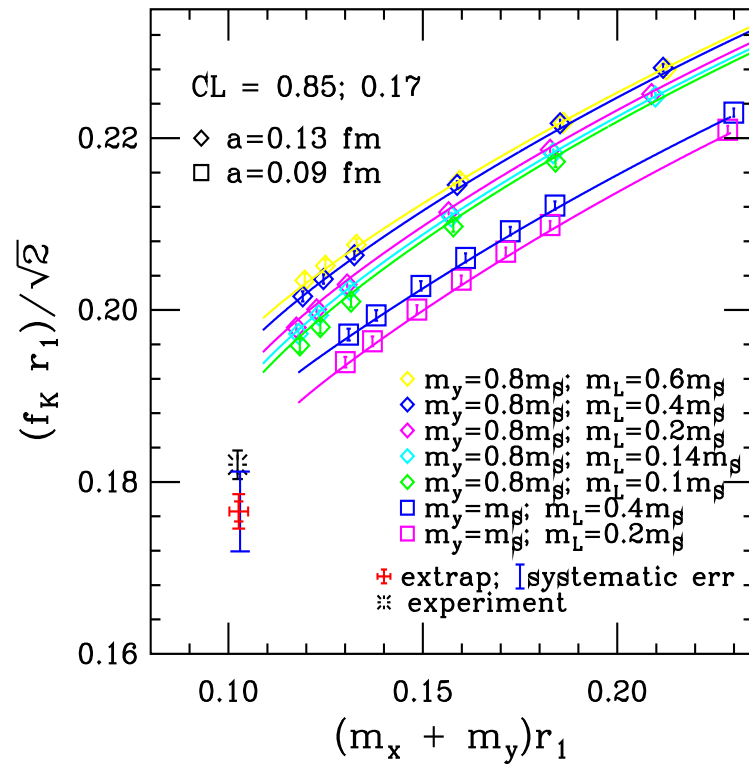
DESY group: different definitions of lattice spacing a :
 $\Lambda_3/f_0 = 30.4(2.9)$ or $\Lambda_3/f_0 = 6.51(57)$

results are a warning that more studies are needed

Chiral Perturbation theory

A staggered example (thanks to C. Bernard)

f_K ; taste viols; $N_f=2+1$, finite V effects removed



f_π, f_K agree with experiment

combinaton of GL coefficients seem to rule out $m_u = 0$ scenario

Two notes:

First: taking $\sqrt{\det}$ amounts in $S\chi PT$ to a partially quenched situation with 2 quenched fermions [Bernard, Golterman](#)

$\det^{1/4}$ with **u,d,s** quarks means in 1-loop of $S\chi PT$

- do $S\chi PT$ with $N_f = 4$ flavours; correct by hand: multiply loops by $1/4$

generalizable to all orders of $S\chi PT$? [Bernard](#)

- replica trick: computation with arbitrary N_u, N_d, N_s of **u,d,s** quarks
- correct/tune by hand: set $N_u = N_d = N_s = 1/4$

Second: nice example of application of χPT ([Chandrasekharan, Jiang](#))

→ very precise computation of condensate and susceptibility using meron cluster algorithm in strong coupling limit

My personal wishlist **V**

Need to discuss all these issues of chiral perturbation theory

→ workshop

Finite size effects

generally $M(L) - M = -\frac{3}{16\pi^2 ML} \int_{-\infty}^{\infty} F(iy) e^{\sqrt{M_\pi^2 + y^2} L}$

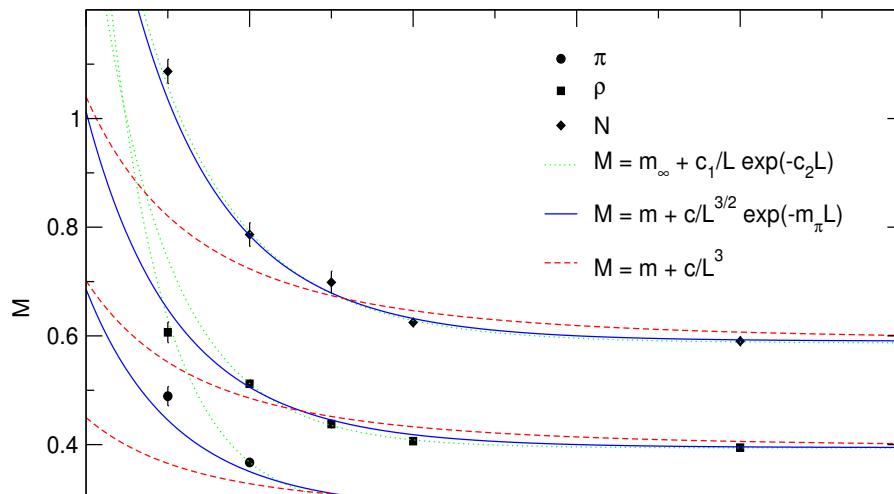
F : $\pi - \pi$ forward scattering amplitude in infinite volume

Lüscher's formula $L^{-3/2} e^{-m_\pi L}$: leading order of F

corrections: Colangelo, Dürr, Sommer

$$\frac{1}{2} L^{-3/2} e^{-m_\pi L} + \frac{1}{\sqrt{2} L^{-3/2}} e^{-\sqrt{2} m_\pi L} + \frac{1}{\sqrt{3} L^{-3/2}} e^{-\sqrt{3} m_\pi L}$$

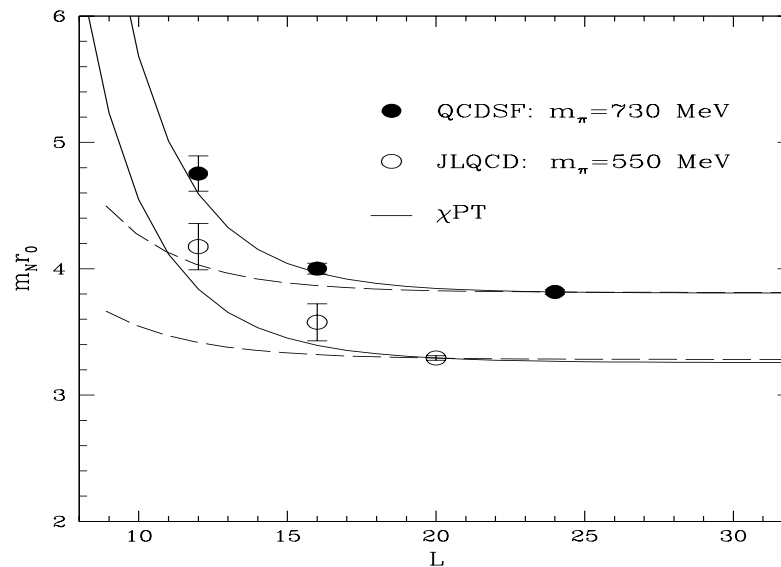
find (found) in practise $M = m_\infty + c/L^3$ (Fukugita, Mino, Okawa, Parisi, Ukawa)



Lippert, Orth, Schilling

→ claim: find expected exponential finite size effects
coeff. of exp. fitted

Finite size effects from chiral perturbation theory



QCDSF collaboration

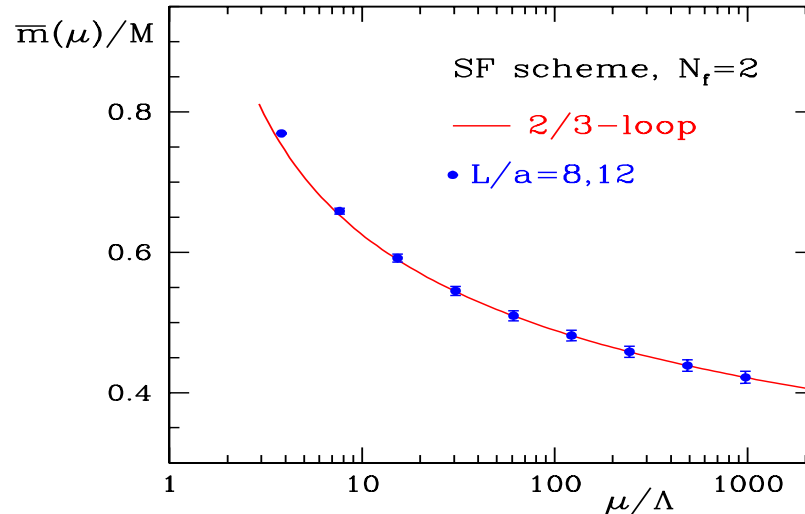
chiral perturbation theory result:

$$\delta = \frac{3g_A^2 m_\pi^2}{16\pi^2 F^2} \int dx \sum_n K_0(Ln \sqrt{m_{N_0}^2 x^2 + m_\pi^2 (1-x)})$$

m_{N_0} Nucleon mass in chiral limit

no free parameter! Leading order agrees with Lüscher formula

Running quark mass: status

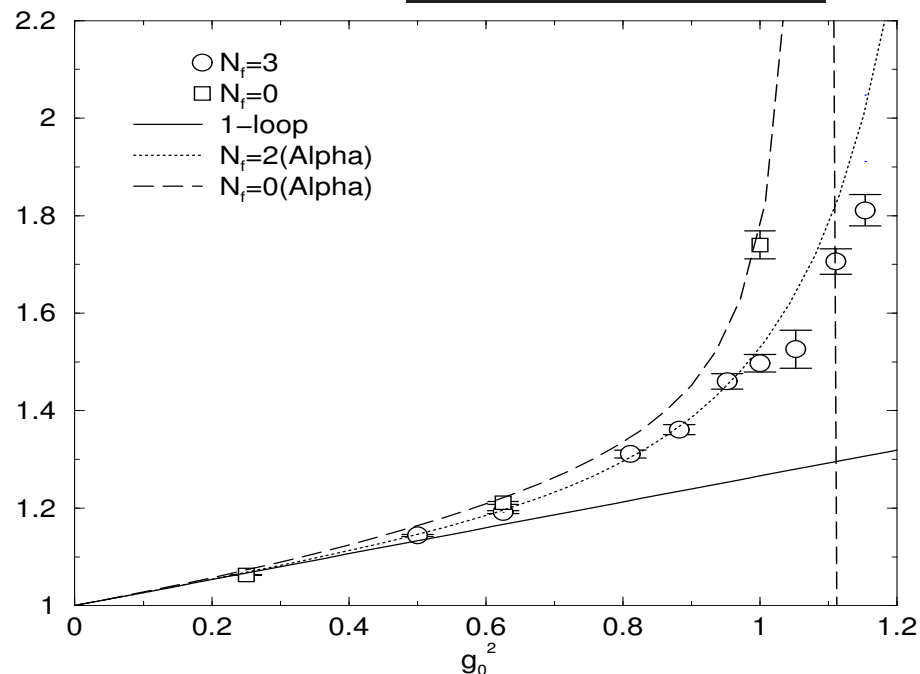


ALPHA
Collaboration

- \rightarrow averaged over $L = 8 \rightarrow L = 16$ and $L = 12 \rightarrow L = 24$
- perturbation theory works well (unfortunately!?)
- point for smallest μ/Λ corresponds to large value of coupling ($L = \max$)
- scale still missing

Towards $N_f = 3$ dynamical Wilson simulations

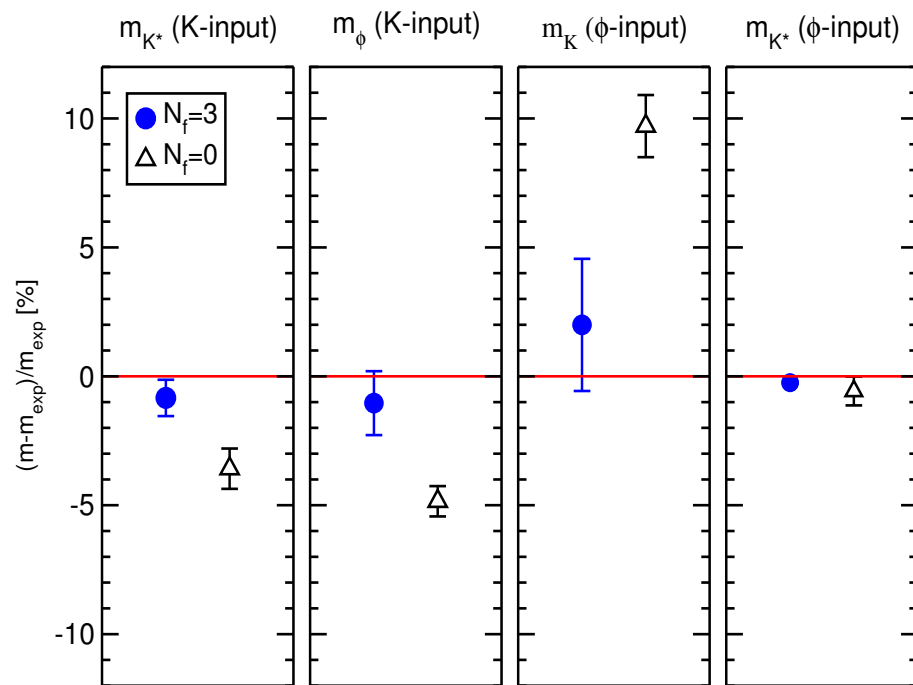
- joint japaneses forces of CP-PACS and JLQCD collaborations
- RG improved gauge and $O(a)$ improved Wilson fermion action
- ⇐ phase transition
- determination of c_{SW} non-perturbatively



- Schrödinger functional

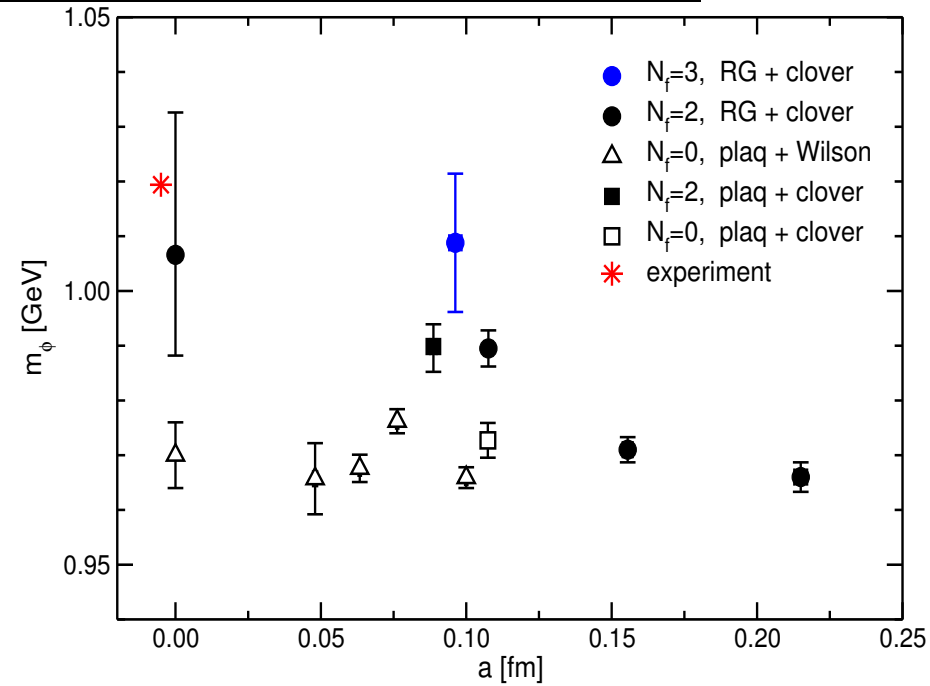
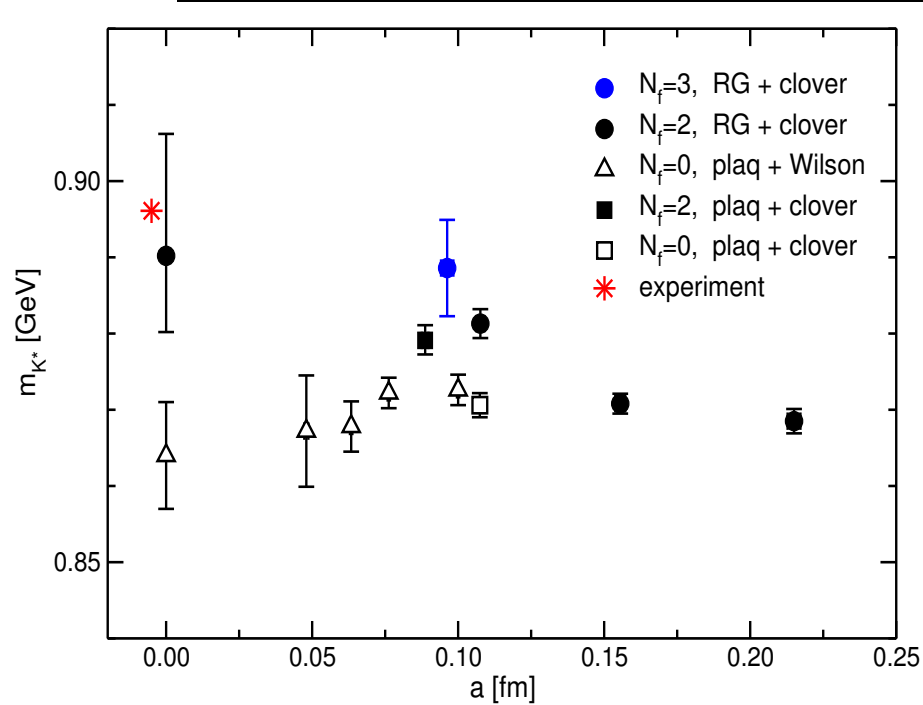
Towards $N_f = 3$ dynamical Wilson simulations:
physical input

$L \approx 1.6\text{fm}$



→ re-assuring: dynamical results can eliminate systematic uncertainty

Towards $N_f = 3$ dynamical Wilson simulations:
meson spectrum



My personal wishlist VII
Add improved staggered results

What was left out, with all my apologies

- domainwall fermions RBC
- localization in QCD Golterman, Shamir
- structure functions MIT, SESAM, QCDSF
- topological susceptibility Hart et.al.
- η' from low-lying eigenmodes SESAM, MIT

Conclusion

New powerful computers (apeNEXT, QCDOC, PC cluster, comm. supercomputers)

- ★ allow transition to serious dynamical fermion simulations
- ★ they are expensive machines that should be used wisely
 - check that your lattice formulation of continuum theory is okay
 - support and participate in ILDG to share configurations/propagators
 - work hard on algorithmic improvements

what do we answer somebody coming with a really big machine and asks

- what action to choose
- what algorithm to employ

Conclusion

there are dangerous animals on the lattice that lurk in the dark

← found surprises in dynamical simulations

⇒ try to use always two actions, depending on your question

baryon spectrum, decay constants etc. (heavier quarks):

improved staggered ↔ improved Wilson with (carefully selected) gauge action

very light quarks: chirally improved actions

(truncated fixed point, domain wall with $L_s \ll 1$, hypercube, FLIC)

↔ actions with exact chiral symmetry

(overlap, domain wall with $L_s \gg 1$)